

Evidence for *D^o-D^o* Mixing at BaBar

Milind V. Purohit University of South Carolina

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<u>Outline</u>

- Introduction
- $D^0 \overline{D}^0$ oscillations
- Search for Mixing / CP violation using $D^0 \rightarrow K^- \pi^+$ decays
- Other searches for mixing / CPV:
 - Lifetime Ratios: $\tau(D^0 \rightarrow K^+ K^-, \pi^+ \pi^-)$ vs $\tau(D^0 \rightarrow K^- \pi^+)$
 - CPV in time-integrated $D^0 \rightarrow K^+ K^-$ and $D^0 \rightarrow \pi^+ \pi^-$ rates.
 - Mixing study using $D^0 \rightarrow K^+ \pi^- \pi^0$ decays
- Comparison with other results, theory





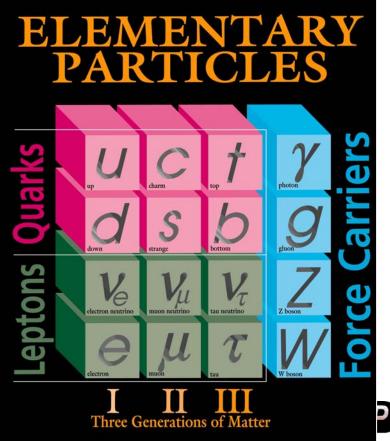
Introduction



Particle Physics

- In the last 100 years or so, starting with the discoveries of the electron, atoms, nuclei, and so on, we have discovered a lot about what the world is made of.
- After 50 years of intense effort, we now know that the physical world is
 - Composed of quarks and leptons
 - Interacting via force carriers called gauge bosons





Standard Model of FUNDAMENTAL PARTICLES AND INTERACTIONS

The Standard Model summarizes the current knowledge in Particle Physics. It is the quantum theory that includes the theory of strong interactions (quantum chromodynamics or QCD) and the unified theory of weak and electromagnetic interactions (electroweak). Gravity is included on this chart because it is one of the fundamental interactions even though not part of the "Standard Model."

FERMIONS matter constituents spin = 1/2, 3/2, 5/2, ...

Leptor		Quarks spin = 1/2					
Flavor	Mass GeV/c ²	Electric charge	Flav	vor	Approx. Mass GeV/c ²	Electric charge	
ve electron	<1×10 ⁻⁸	0	u u	ıp	0.003	2/3	
e electron	0.000511	-1	dd	lown	0.006	-1/3	
ν_{μ} muon neutrino	<0.0002	0	Ca	harm	1.3	2/3	
μ muon	0.106	-1	S s	trange	0.1	-1/3	
v_{τ} tau neutrino	<0.02	0	t t	ор	175	2/3	
au tau	1.7771	-1	b b	oottom	4.3	-1/3	

Spin is the intrinsic angular momentum of particles. Spin is given in units of h, which is the quantum unit of angular momentum, where $h = h/2\pi = 6.58 \times 10^{-25}$ GeV s = 1.05x10⁻³⁴ J s.

Electric charges are given in units of the proton's charge. In SI units the electric charge of the proton is 1.60×10⁻¹⁹ coulombs.

The energy unit of particle physics is the electronvolt (eV), the energy gained by one electron in crossing a potential difference of one volt. **Masses** are given in GeV/c² (remember $E = mc^2$), where 1 GeV = 10⁹ eV = 1.60×10⁻¹⁰ joule. The mass of the proton is 0.938 GeV/c² = 1.67×10⁻²⁷ kg.

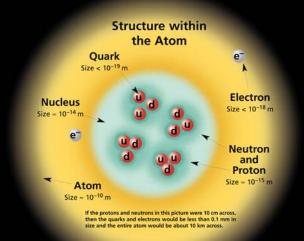
	Baryons qqq and Antibaryons qqq Baryons are fermionic hadrons. There are about 120 types of baryons.						
Symbol Name Quark Content Electric Mass GeV/c ² Spin							
р	proton	uud	1	0.938	1/2		
p	anti- proton	ūūd	~1	0.938	1/2		
n	neutron	udd	0	0.940	1/2		
Λ	lambda	uds	0	1.116	1/2		
Ω-	omega	555	-1	1.672	3/2		

Matter and Antimatter

For every particle type there is a corresponding antiparticle type, denoted by a bar over the particle symbol (unless + or - charge is shown) Particle and antiparticle have identical mass and spin but opposite charges. Some electrically neutral bosons (e.g., Z^0 , γ , and $\eta_c = c\overline{c}$, but not $K^0 = d\bar{s}$) are their own antiparticles.

Figures

These diagrams are an artist's conception of physical processes. They are not exact and have no meaningful scale. Green shaded areas represent the cloud of gluons or the gluon field, and red lines the quark paths.



PROPERTIES OF THE INTERACTIONS

BOSONS

Unified Electroweak spin = 1				
Name	Mass GeV/c ²	Electric charge		
γ photon	0	0		
W-	80.4	-1		
W+	80.4	+1		
Z ⁰	91.187	0		

force carriers spin = 0, 1, 2,

Unified Electroweak spin = 1					
Name	Electric charge				
γ photon	0	0			
W-	80.4	-1			
W+	80.4	+1			
70	01 197	0			

Strong (color) spin = 1					
Name	Mass GeV/c ²	Electric charge			
g gluon	0	0			

Color Charge

Each quark carries one of three types of "strong charge," also called "color charge." These charges have nothing to do with the colors of visible light. There are eight possible types of color charge for gluons. Just as electri-

cally-charged particles interact by exchanging photons, in strong interactions color-charged particles interact by exchanging gluons. Leptons, photons, and W and Z bosons have no strong interactions and hence no color charge.

Quarks Confined in Mesons and Baryons

One cannot isolate quarks and gluons; they are confined in color-neutral particles called hadrons. This confinement (binding) results from multiple exchanges of gluons among the color-charged constituents. As color-charged particles (quarks and gluons) move apart, the energy in the color-force field between them increases. This energy eventually is converted into additional quark-antiquark pairs (see figure below). The quarks and antiquarks then combine into hadrons; these are the particles seen to emerge. Two types of hadrons have been observed in nature: mesons qq and baryons qqq.

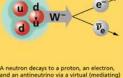
Residual Strong Interaction

The strong binding of color-neutral protons and neutrons to form nuclei is due to residual strong interactions between their color-charged constituents. It is similar to the residual elec-trical interaction that binds electrically neutral atoms to form molecules. It can also be viewed as the exchange of mesons between the hadrons.

qq	G									Meso	ns aā		
1.0	balls.	Property	Gravitational	Weak	Electromagnetic	The State of the S	ong			ons are bo about 14	sonic hadro		
-				(Electr	oweak)	Fundamental	Residual		1	-		1	-
5	pin	Acts on:	Mass – Energy	Flavor	Electric Charge	Color Charge	See Residual Strong Interaction Note	Symbol	Name	Quark content	Electric charge	Mass GeV/c ²	Spin
١,	1/2	Particles experiencing:	All	Quarks, Leptons	Electrically charged	Quarks, Gluons	Hadrons	π^+	pion	uđ	+1	0.140	0
		Particles mediating:	Graviton (not yet observed)	W+ W- Z ⁰	γ	Gluons	Mesons			sū			
1	/2	Strength relative to electromag 10 ⁻¹⁸ m	10-41	0.8	1	25	Not applicable	к-	kaon	su	-1	0.494	0
1	1/2	for two u quarks at: 3×10 ⁻¹⁷ m	Conservation of the second	10 ⁻⁴	1	60	to quarks	ρ^+	rho	ud	+1	0.770	1
1	1/2	for two protons in nucleus	10 ⁻³⁶	10 ⁻⁷	1	Not applicable to hadrons	20	B0	B-zero	db	0	5.279	0

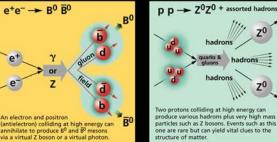
e⁺e⁻ → B⁰ B⁰

e



 $n \rightarrow p e^- \overline{\nu}_{o}$

and an antineutrino via a virtual (mediating) W boson. This is neutron B decay.



B⁰

Two protons colliding at high energy can produce various hadrons plus very high mass particles such as Z bosons. Events such as this one are rare but can yield vital clues to the structure of matter

hadrons

adrons

hadrons

Z⁰

The Particle Adventure

Visit the award-winning web feature The Particle Adventure at http://ParticleAdventure.org

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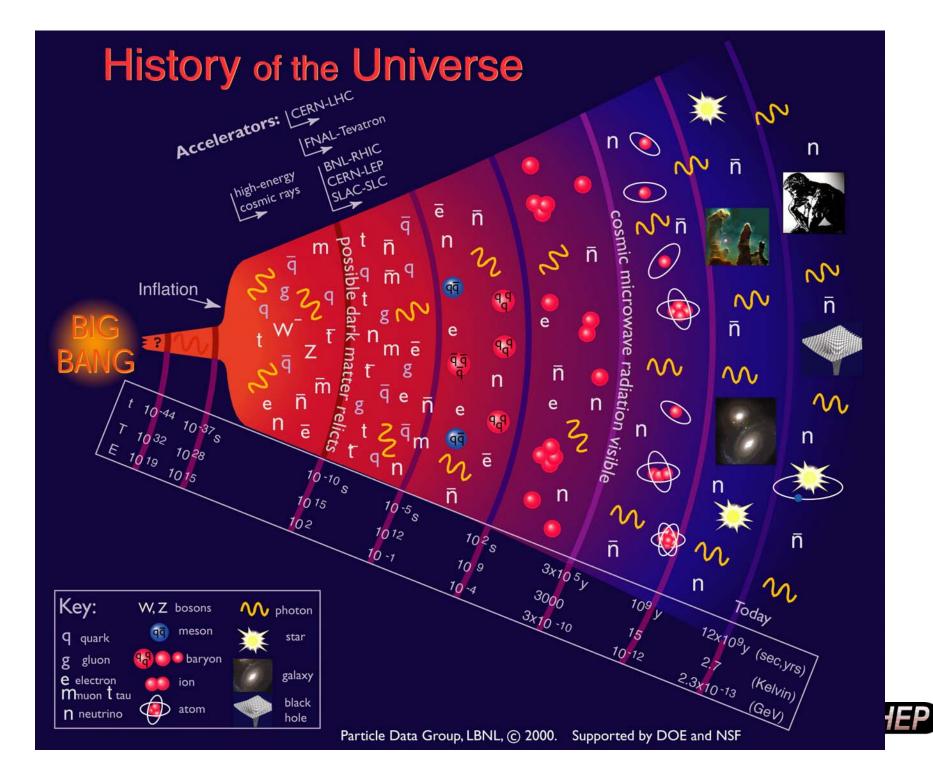
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Beyond the Standard Model

Is the Standard Model the final picture of particle physics?

Answer: Most certainly not!

How do we go beyond it?

- •Study existing problems with the model
- •Look for new particles / effects directly
- •Look for new particles / effects indirectly





Neutral Meson Oscillations



Neutral Meson Mixing

Mixing can occur in four neutral mesons:

K ⁰	Mass: $\sim 0.5 \text{ GeV/c}^2$
D ⁰	Mass: $\sim 1.9 \text{ GeV/c}^2$
B ⁰	Mass: $\sim 5.3 \text{ GeV/c}^2$
$B^0_{\ \mathrm{s}}$	Mass: $\sim 5.4 \text{ GeV/c}^2$

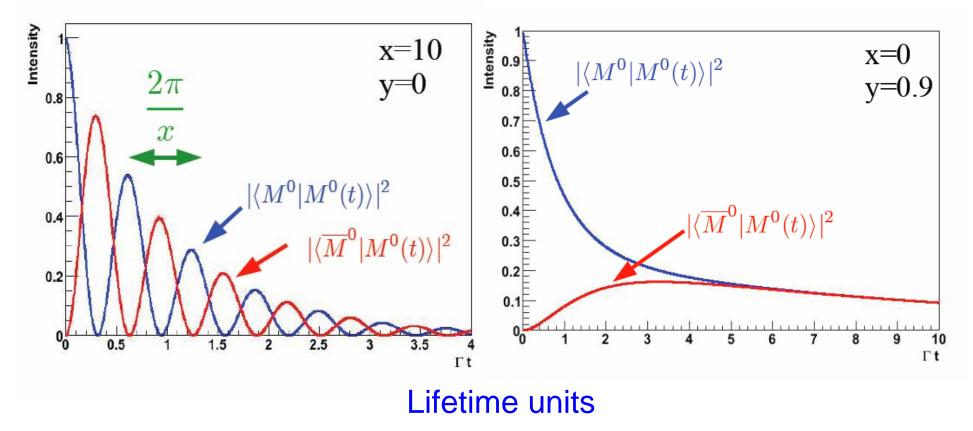
Will present mixing measurement for D^0 meson

Note: *D*⁰ meson first discovered at SLAC Mark-I, PRL 37, 255 (1976)



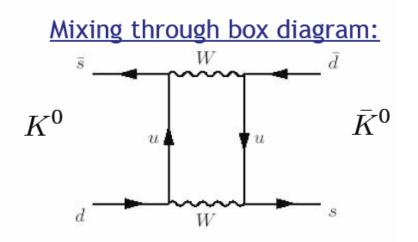
Some visual examples

Probability to find a $M^{0}(\overline{M}^{0})$ after a given time





The prediction of charm



No tree level Flavor Changing Neutral Currents (FCNC) in SM

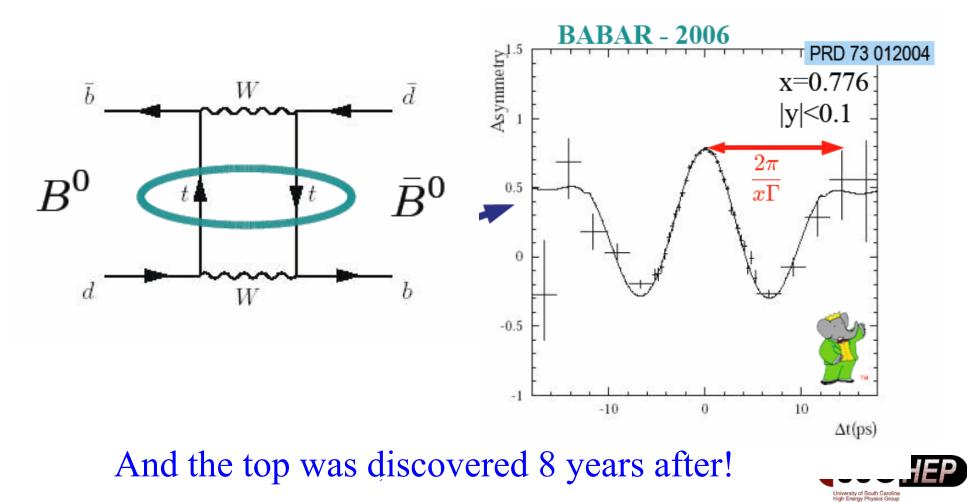
Glashow, Iliopoulus and Maiani (1970): FCNC calculated from single quark loop still too large Introduce additional loop with new c quark

GIM predicted charm quark 4 years before observation

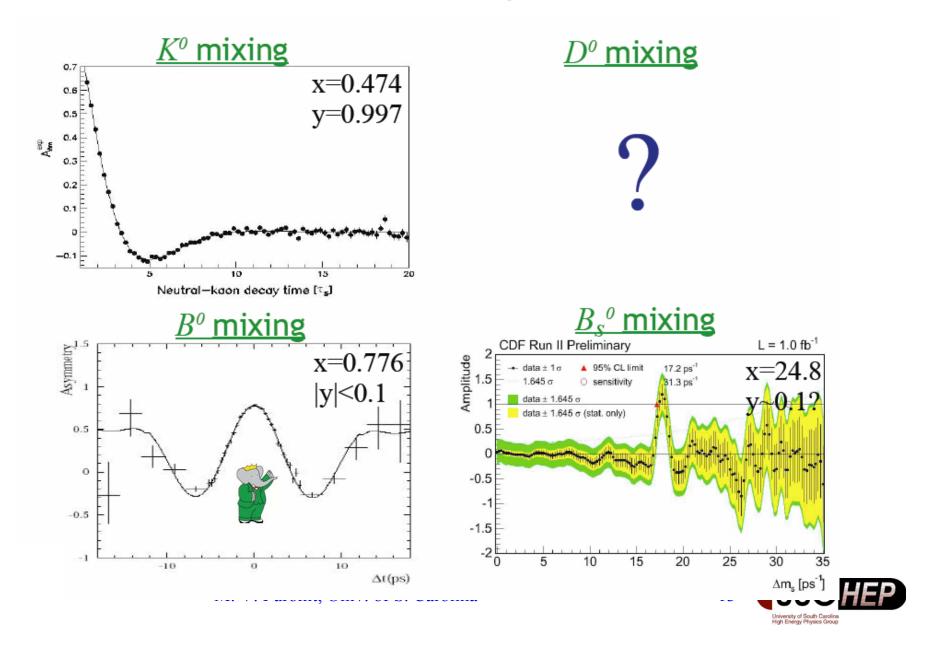


B⁰ mixing and the discovery of the t

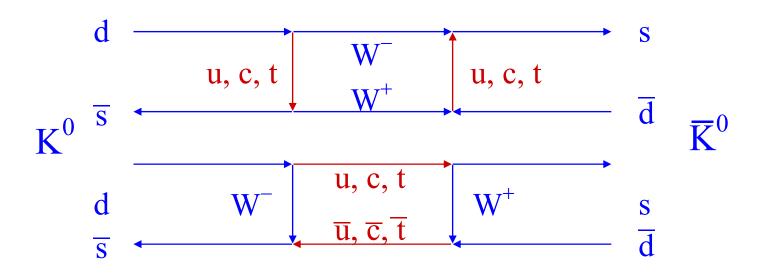
B⁰ mixing was argued by UA1 and directly observed by ARGUS in 1987 Large mixing frequency implied *t* quark was heavy (m_t > 50 GeV/c²)



The missing tile



Kaon oscillations



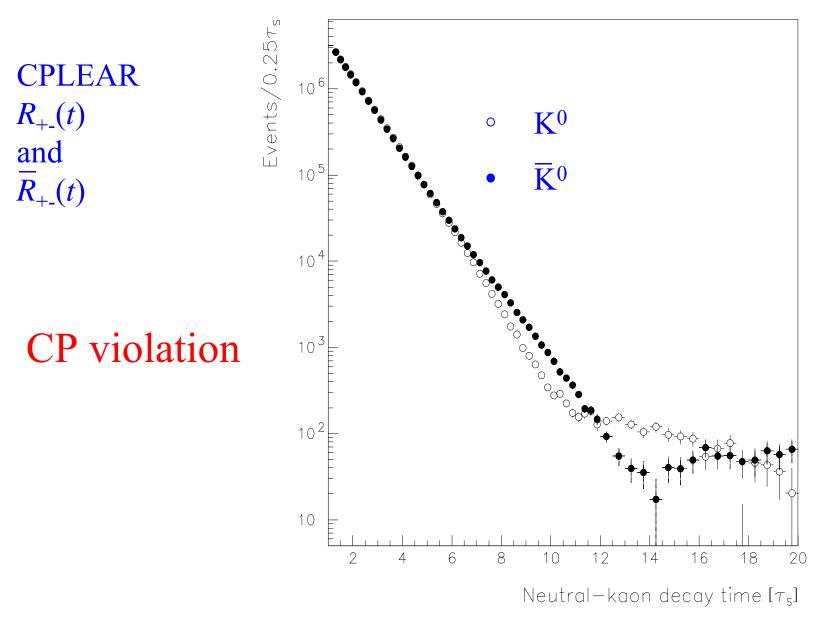
- So say at t=0, pure K^o,
 - later a superposition of states



Assume CP K⁰ Decay $|K_s\rangle \approx |K_1\rangle = \frac{1}{\sqrt{2}} \left(|K^0\rangle + |\overline{K^0}\rangle \right)$ CP=+1 mass eigenstates $\begin{cases} K_S \\ K_S \end{cases}$ $|K_L\rangle \approx |K_2\rangle = \frac{1}{\sqrt{2}} \left(|K^0\rangle - |\overline{K^0}\rangle \right) \quad \text{CP=-1} \quad p = q = \frac{1}{\sqrt{2}}$ • In that case

K_s branching fractions: $\pi^+\pi^- \approx 69\%$, $\pi^0\pi^0 \approx 31\%$ K_L branching fractions: $\pi^0 \pi^0 \pi^0 \approx 21\%$, $\pi^+ \pi^- \pi^0 \approx 13\%$, $\ell^\mp \pi^\pm \nu \approx 66\%$ $\tau_{\rm S} = \frac{1}{\Gamma_{\rm S}} = (0.8934 \pm 0.0008) \times 10^{-10} {\rm s}$ $\begin{array}{c} \mathsf{K}_{s} \to \pi^{o} \pi^{o} \\ \mathsf{K}_{s} \to \pi^{+} \pi^{-} \end{array} \right\} \qquad \mathcal{C}\mathsf{P}=+1 \qquad \tau_{L} = \frac{1}{\Gamma_{L}} = (5.17 \pm 0.04) \times 10^{-8} \text{ s} \\ \mathsf{K}_{L} \to \pi^{+} \pi^{-} \pi^{o} \\ \mathsf{K}_{L} \to \pi^{o} \pi^{o} \pi^{o} \end{array} \qquad \mathcal{C}\mathsf{P}=-1 \qquad \Delta m = m_{L} - m_{S} = (0.5301 \pm 0.0014) \times 10^{10} h \text{ s}^{-1} \\ \mathsf{K}_{L} \to \pi^{o} \pi^{o} \pi^{o} \end{array}$ $\therefore x = \frac{\Delta M}{\Delta M} \approx 0.95$





M. V. Purohit, Univ. of S. Carolina



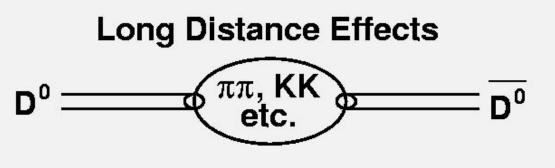


D^0 - $\overline{D}{}^0$ Oscillations



<u>Physics of $D^0\overline{D}^0$ – mixing</u>

• Common final states lead to mixing:

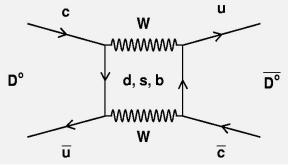


• One can naively estimate the mixing rate: Since BR ~ few x 10⁻³, one expects rate "r" ~ 10^{-5} .



Physics of $D^0\overline{D}^0$ – mixing, contd.

• Mixing at the quark level in the Standard Model:



• Predicted rate for mixing: $r_{mix} \sim 10^{-7}$



Possibility of CP violation in mixing:

- First noted by Pais and Treiman (Phys. Rev. D12, 2744 (1975)).
- It is possible that

 $r(D^0 \rightarrow D^0 bar) \neq r(D^0 bar \rightarrow D^0)$

• This is most likely to happen in the interference between the mixing and DCS amplitudes, since CP violation requires two amplitudes.



Neutral Meson systems

- Two-level system (M⁰,M⁰)
 - Weak interactions remove degeneracy, make them unstable

Time evolution by Schrödinger eq.:
$$i \frac{\partial}{\partial t} \begin{pmatrix} |M^{0}(t)\rangle \\ |\overline{M}^{0}(t)\rangle \end{pmatrix} = \begin{pmatrix} \mathsf{M} - \frac{i}{2} \Gamma \\ \mathsf{M} - \frac{i}{2} \Gamma \end{pmatrix} \begin{pmatrix} |M^{0}(t)\rangle \\ |\overline{M}^{0}(t)\rangle \end{pmatrix}$$

2x2 hermitian matrices Mesons decay!

Mass eigenstates:

$$|M_{1,2}\rangle = p|M^0\rangle \pm q|\overline{M}^0\rangle$$

Propagate with separate mass $m_{1,2}$ and width $\Gamma_{1,2}$: $|M_{1,2}(t)\rangle = e^{-i(m_{1,2}-i\Gamma_{1,2}/2)t}|M_{1,2}(t=0)\rangle$



Neutral meson oscillations

Time evolution for meson of *known flavor at t=0*

$$x = \frac{m_2 - m_1}{\Gamma}$$
$$y = \frac{\Gamma_2 - \Gamma_1}{2\Gamma}$$
$$\Gamma = \frac{\Gamma_2 + \Gamma_1}{2}$$

$$|M^{0}(t)\rangle = e^{-\bar{\gamma}t/2} \left(\cosh(\Delta\gamma t/2) |M^{0}\rangle - \frac{q}{p} \sinh(\Delta\gamma t/2) |\overline{M}^{0}\rangle \right)$$

Where $\Delta\gamma = (y + ix)\Gamma$ $\bar{\gamma} = (\Gamma_{1} + \Gamma_{2})/2 - i(m_{1} + m_{2})$

M⁰ "oscillates" into M⁰! (also dubbed "mixing")

An opposite flavor component appears after a while!



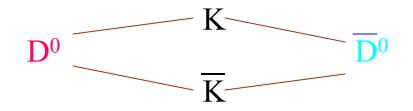
Short and Long distance

• Predictions for x and y:

$$\begin{pmatrix} \mathbf{M} - \frac{i}{2} \mathbf{\Gamma} \end{pmatrix}_{ij} = \frac{\langle D_i | H_{\text{eff}} | D_j \rangle}{2m_D} = m_D^{(0)} \delta_{ij} \\ + \frac{\langle D_i | H_w | D_j \rangle}{2m_D} + \frac{1}{2m_D} \sum_n \frac{\langle D_i | H_w | n \rangle \langle n | H_w | D_j \rangle}{m_D^{(0)} - E_n + i\epsilon}.$$

$$\mathbf{y} \quad \Gamma_{ij} = \frac{1}{2m_D} \sum_n \langle D_i | H_w | n \rangle \langle n | H_w | D_j \rangle \, \delta(E_n - m_D).$$

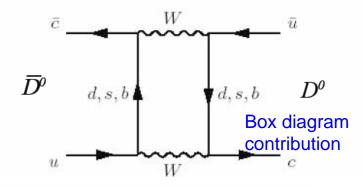
$$\mathbf{Sum of intermediate}_{\mathbf{REAL states}}$$





SM prediction for charm mixing

SM charm mixing box has down-type quarks in loop



Effective GIM suppression: $x \propto \frac{(m_s^2 - m_d^2)^2}{m_c^2}$ (bottom quark
ruled out by V_{CKM}) \checkmark $x \sim 10^{-5}$ Tiny!

Naively $x, y \sim \sin \theta_c^2 \times \left(\frac{m_s}{\Lambda_{\text{hadr.}}}\right)^2 \lesssim O(10^{-3})$

Always hard to evaluate SU(3) breaking !!! (*HQET*, propagation of common hadronic states,...)

x, $y \sim \sin \theta_c^2 \times [SU(3) \text{ breaking}].$

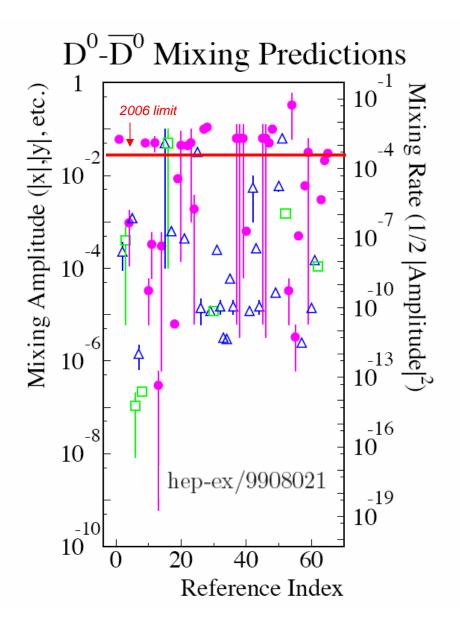
G. Burdman and I. Shipsey, Ann. Rev. Nucl. and Part. Sci. **53**, 431 (2003).

SU(3) breaking effect more important for y

$$x \lesssim 10^{-3}, \quad y \lesssim 10^{-2}.$$



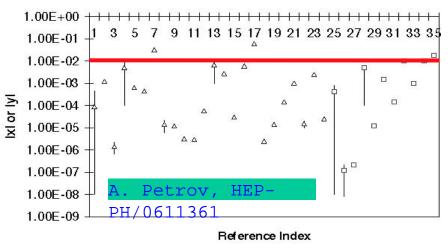
New Physics in Charm?



 Δ : Standard model predictions for x

□: Standard model predictions for y

- •: New physics predictions for x
 - Hard to see a clear prediction
 - Pushing the limit down excludes models



Standard Model mixing predictions

olina

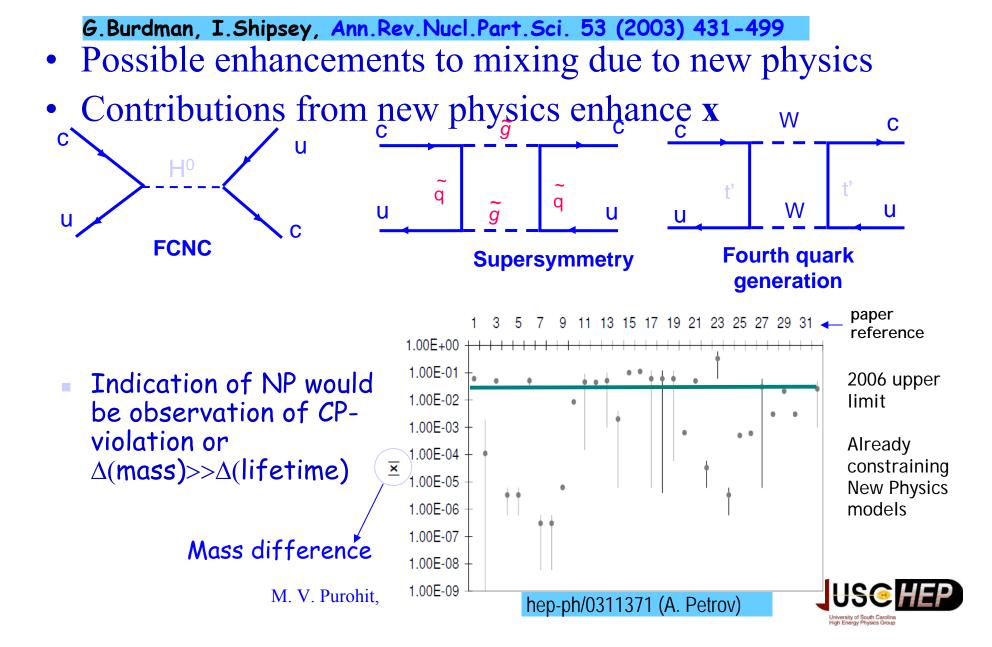


New Physics

- Small Standard Model mixing ⇒ Large window for discovery of new physics
- However, in 1995 Wolfenstein reminded us that long distance contributions can be large
- SM calculations of box diagram redone: r_{mix} in the SM could be as large as 10^{-3} according to Georgi (1992). Confirmed by Ohl, Ricciardi and Simmons (1993). More recently Falk *et al.* (2002) show that phase space effects alone can yield y ~ 1% via SU(3) breaking in the SM.



New Physics Contribution to Charm Mixing



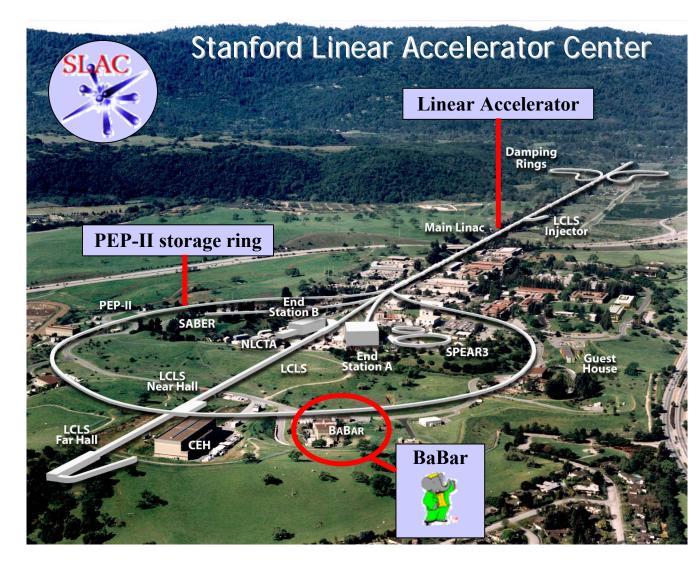


Charm Mixing in $D^{O} \rightarrow K\pi$ Decay at BaBar

(Phys. Rev. Lett. 98:211802, 2007)



PEP-II, a B-Factory (and Charm)



High-luminosity asymmetric energy e^+e^- collider at Υ (4S) resonance

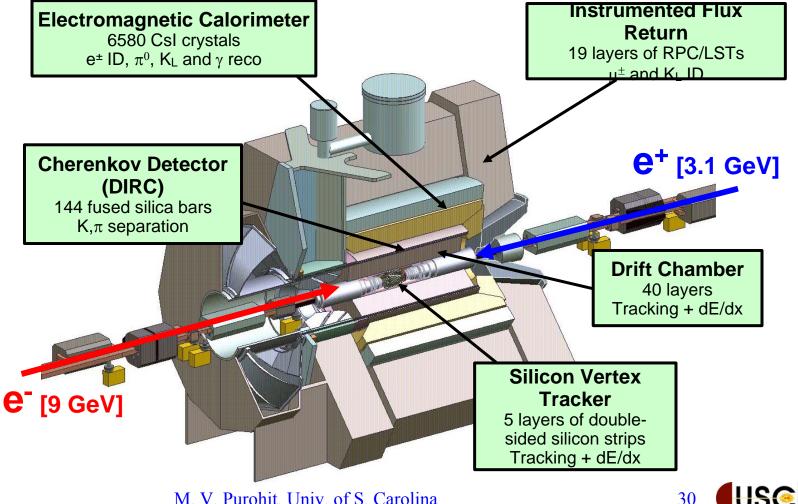
B-Factory built for study of CP-violation and other CKMphysics in *B* meson decays

~10 Hz of $B\overline{B}$



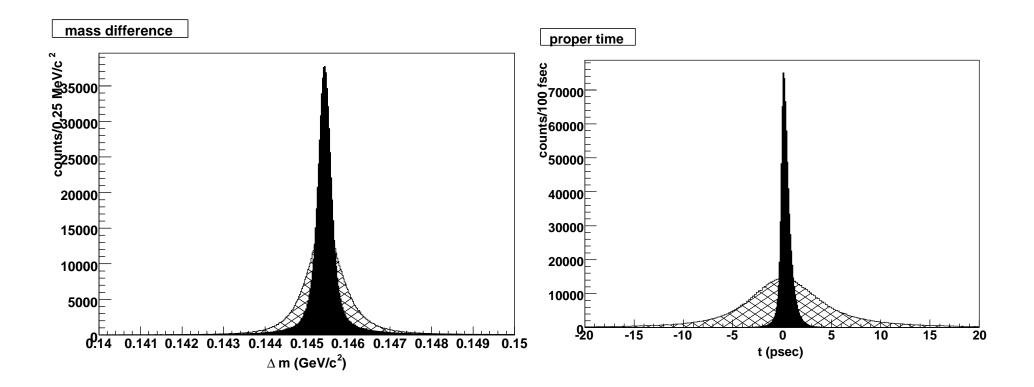
The BaBar Experiment

BaBar is a large acceptance experiment with excellent particle reconstruction and identification capability





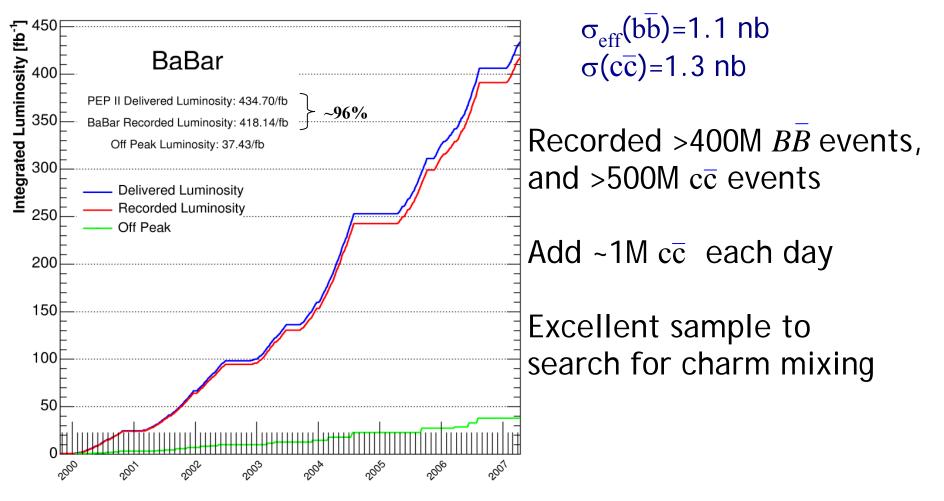
Effect of Beam Spot Constraint on ΔM and t





B-Factory: High Luminosity

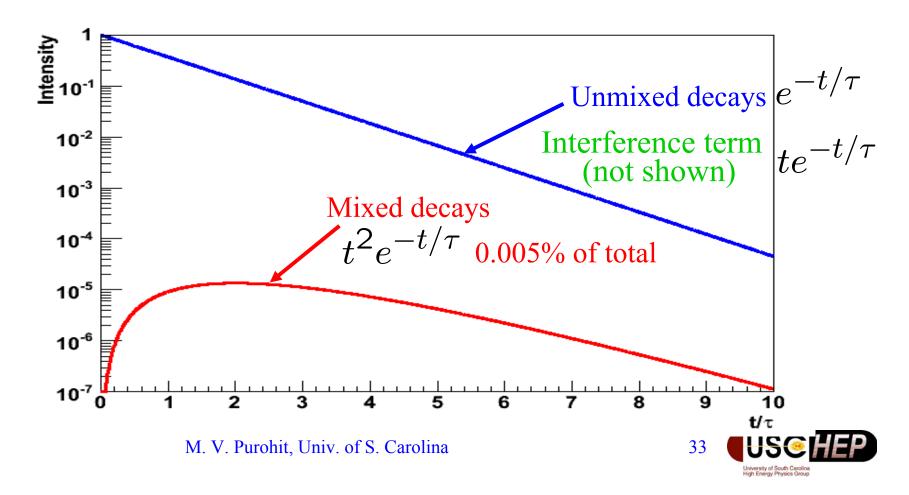
High luminosity recorded efficiently





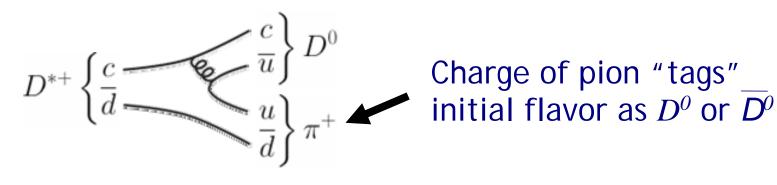
Principle of Mixing Measurement

- Produce clean sample of D^0 and \overline{D}^0
- *****Identify flavor (D^0 or \overline{D}^0 ?) at decay time
- Measure rate of mixed decays as function of time (Distributions shown without time smearing)



Production Flavor

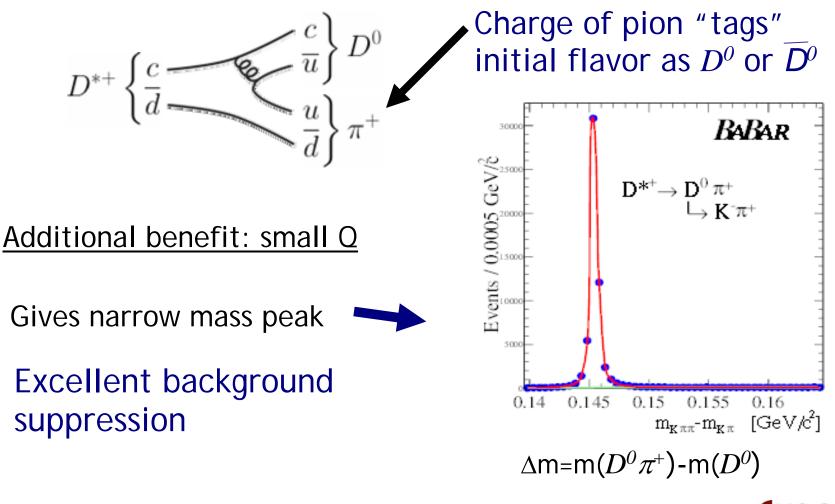
Use D^0 from $D^{*+} \rightarrow D^0 \pi^+$ decays:





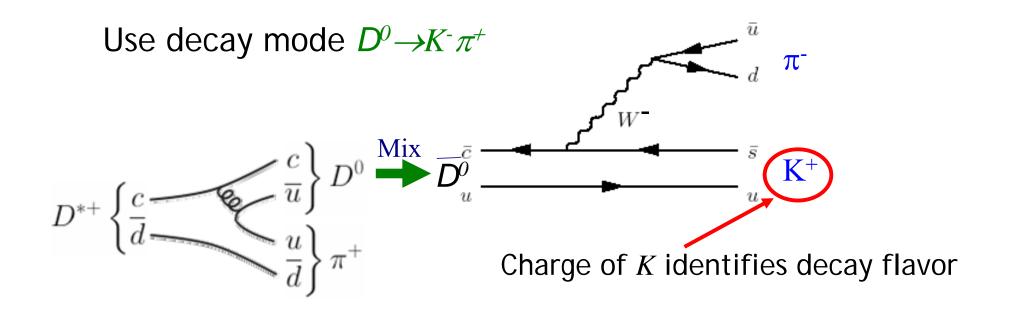
Production Flavor

Use D^0 from $D^{*+} \rightarrow D^0 \pi^+$ decays:





Flavor at Decay

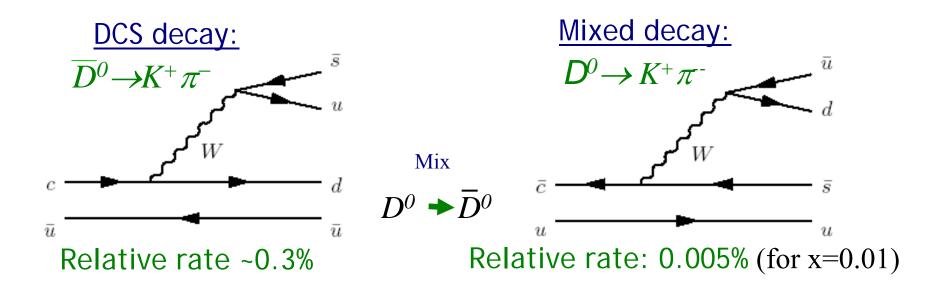


If opposite flavor: Wrong-sign (WS) event – mixing occurred If same flavor: Right-sign (RS) events – unmixed decay



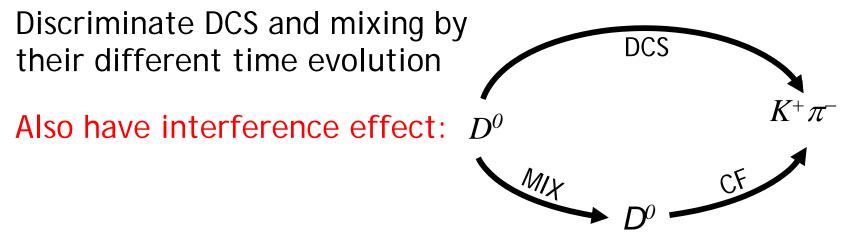
Doubly-Cabibbo Suppressed Decays

Hadronic decays do not uniquely identify decay flavor Get unmixed wrong-sign decays from DCS decays





Time-Evolution of $D^0 \rightarrow K\pi$ Decays



Time evolution:

$$\frac{\Gamma_{WS}(t)}{e^{-t/\tau}} \propto R_D + \sqrt{R_D} y' \left(\frac{t}{\tau}\right) + \left(\frac{x'^2 + y'^2}{4}\right) \left(\frac{t}{\tau}\right)^2$$

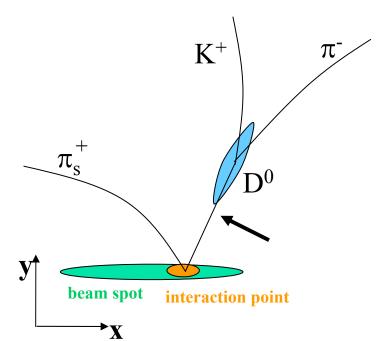
where $x' = x \cos \delta + y \sin \delta$ $y' = y \cos \delta - x \sin \delta$

and δ is the phase difference between DCS and CF decays.



Event Selection

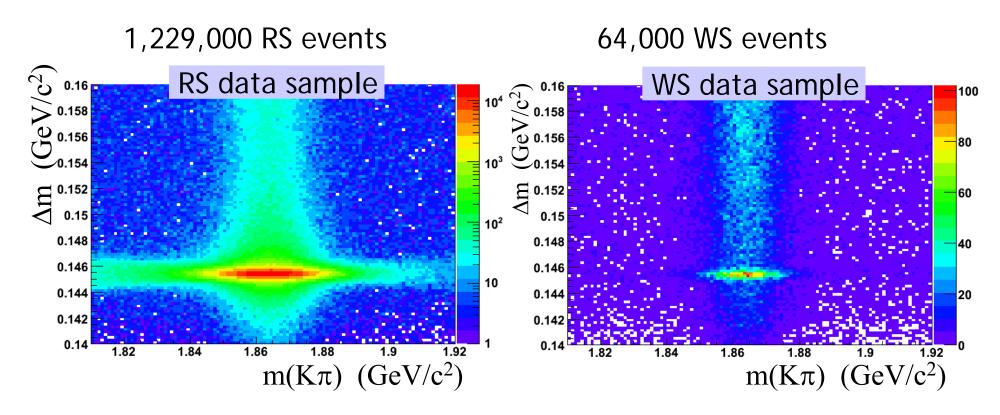
 $\frac{D^{0} \text{ selection}}{\diamond}$ \Rightarrow Identified K and π \Rightarrow p*(D⁰)> 2.5 GeV/c \Rightarrow 1.81<m(Kπ)<1.92 GeV/c² Slow π selection: \Rightarrow p*(π_s)< 0.45 GeV/c \Rightarrow p_{lab}(π_s)> 0.1 GeV/c \Rightarrow 0.14<Δm<0.16 GeV/c² Δ m=m(Kππ_s)-m(Kπ)



<u>Vertexing</u>: (Also greatly improves t resolution) D^0 and π_s constrained to luminous region Fit probability > 0.1% Reconstructed decay time, t: -2<t<4 ps Estimated decay time error, $\delta t < 0.5$ ps



Selected Events



Separate signal from background using m(K π) and Δ m



Fit Procedure

Unbinned maximum likelihood fit in several steps

(fitting 1+ million events takes a long time)

Fit to m(K π) and Δ m distribution:

RS and WS samples fit simultaneously
Signal and some background parameters shared

All parameters determined in fit to data, not MC

Fit RS decay time distribution:

*Determines D^0 lifetime and resolution function *Include event-by-event decay time error δt in resolution *Use m(K π) and Δm to separate signal/bkgd (fixed shapes)

Fit WS decay time distribution:

Use D⁰ lifetime and resolution function from RS fit
Compare fit with and without mixing (and CP violation)



Fit Procedure

Unbinned maximum likelihood fit in several steps (fitting 1+ million events takes a long time)

Fit to m(Kπ) and Δm distribution:
RS and WS samples fit simultaneously
Signal and some background parameters shared
All parameters determined in fit to data, not MC

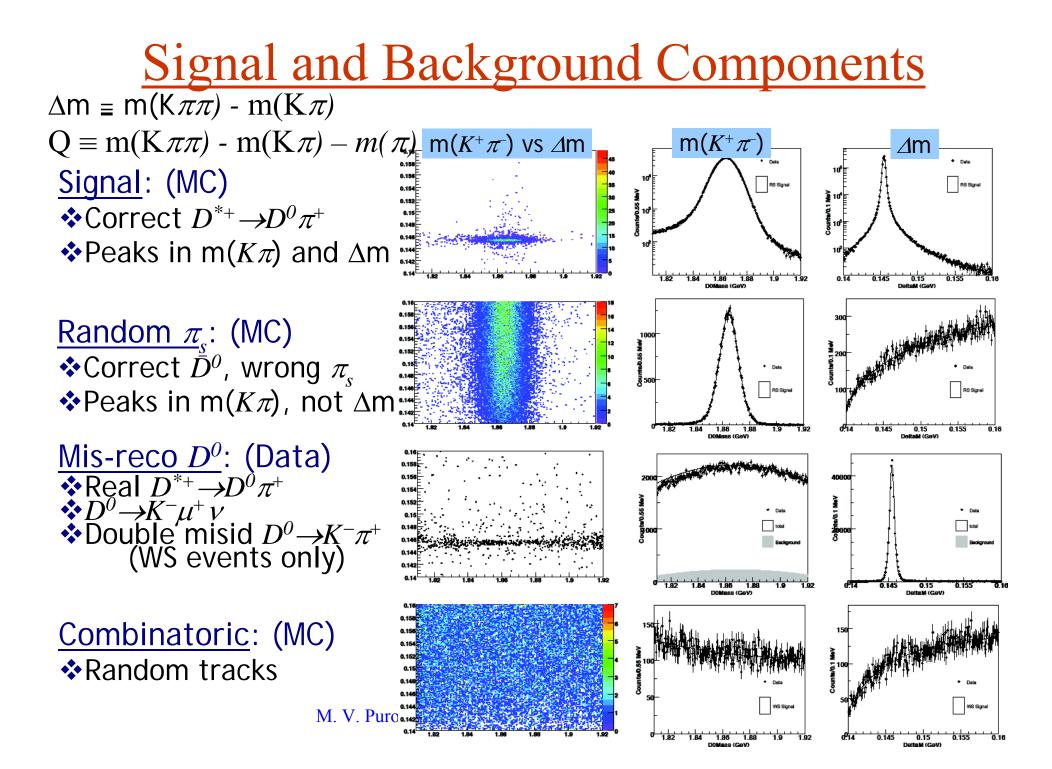
Fit RS decay time distribution:

*Determines D^0 lifetime and resolution function *Include event-by-event decay time error δt in resolution *Use m(K π) and Δm to separate signal/bkgd (fixed shapes)

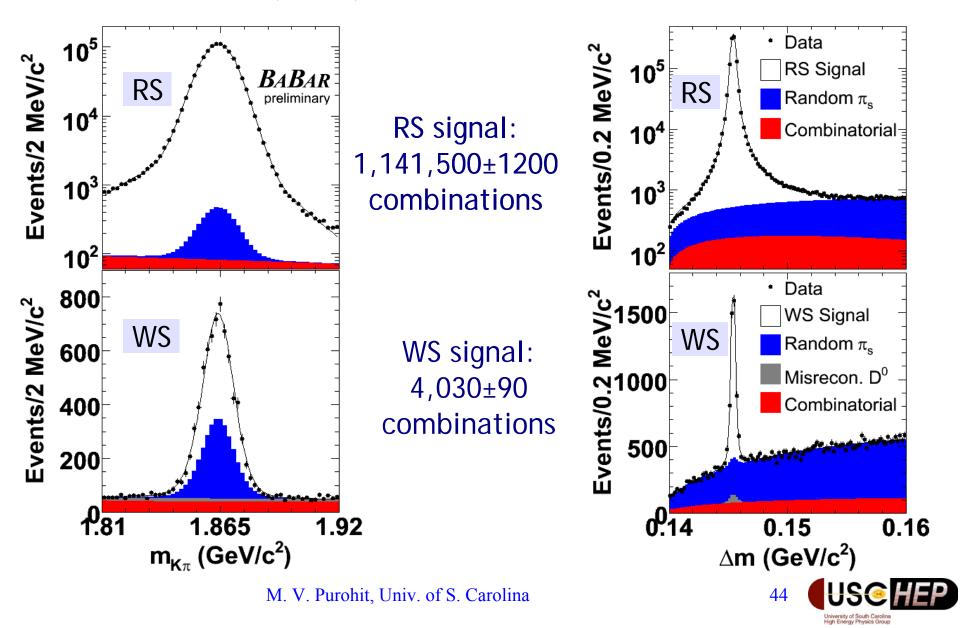
Fit WS decay time distribution:

Use D⁰ lifetime and resolution function from RS fit
Compare fit with and without mixing (and CP violation)





m(K π)- Δ m Fit Results



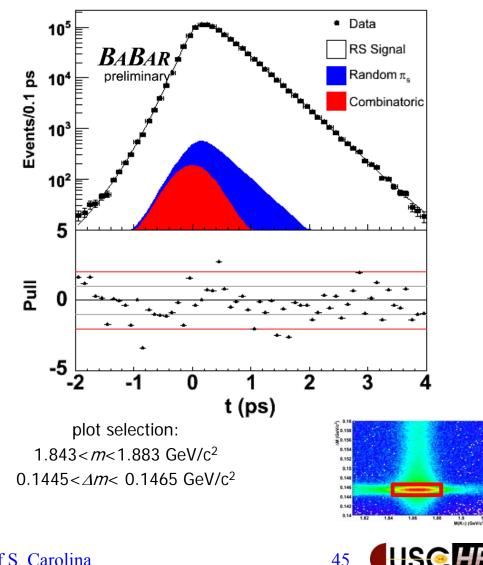
RS Decay Time Fit RS decay time, signal region

D⁰ lifetime and resolution function fitted in RS sample

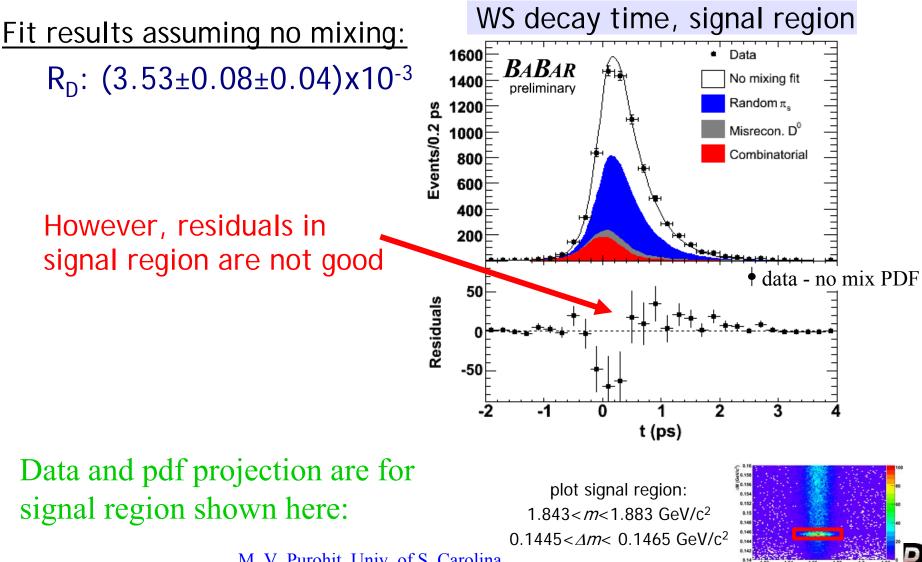
 $\tau = (410.3 \pm 0.6 (stat.))$ fs

Consistent with PDG (410.1±1.5 fs)

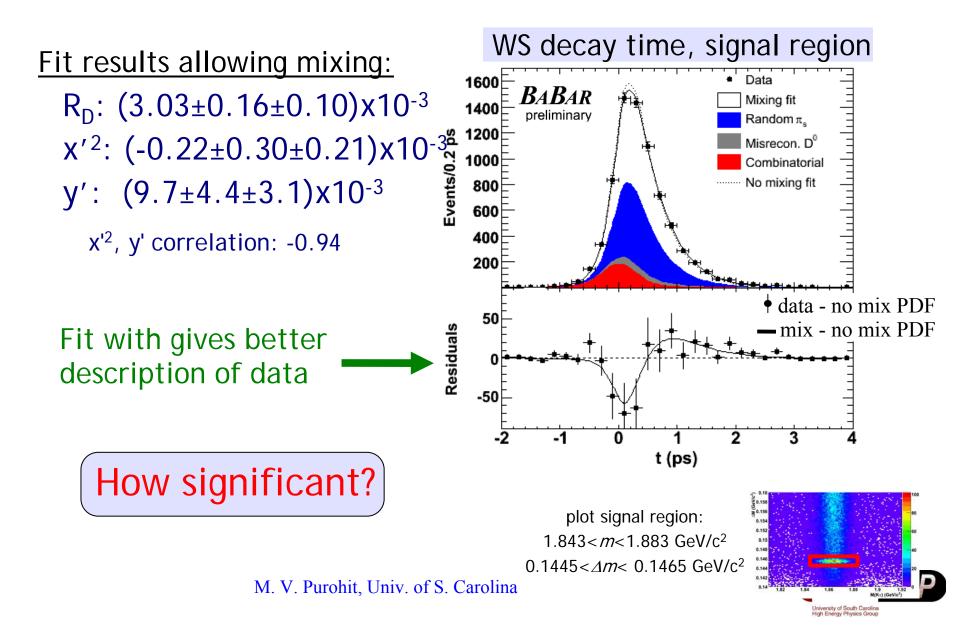
Systematics dominated by resolution function



WS Fit with no Mixing

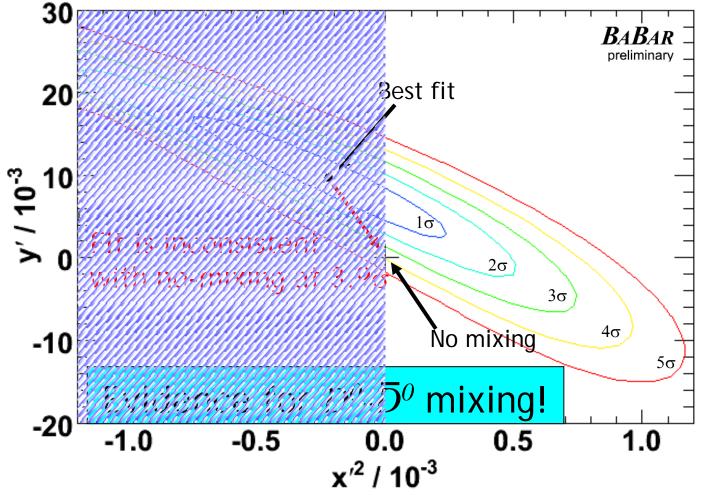


WS Fit with Mixing



Signal Significance with Systematics

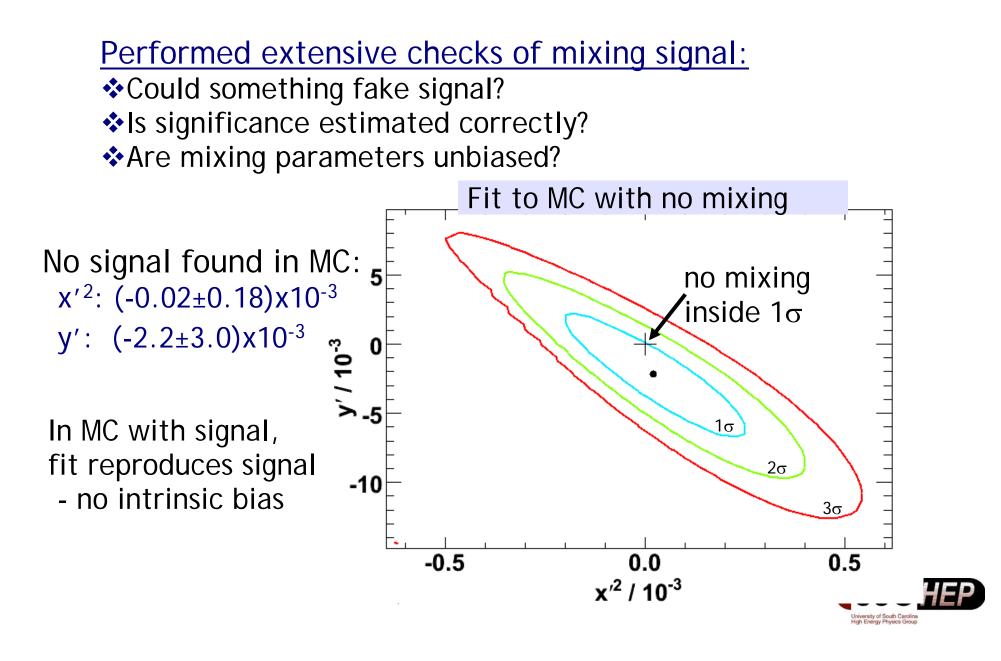
Including systematics decreases signal significance



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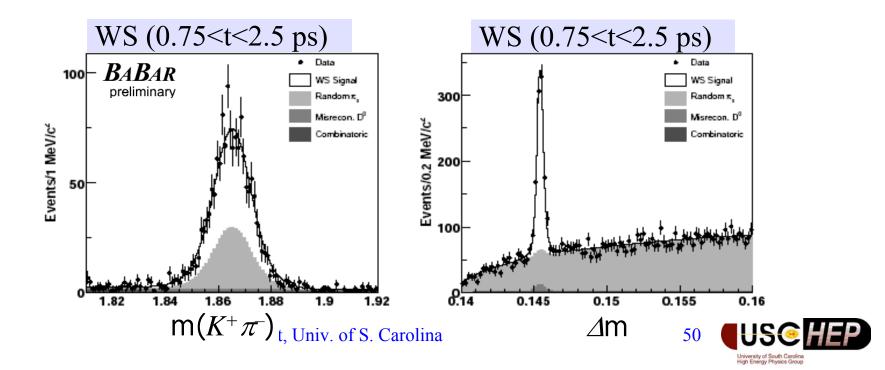
Validation Studies



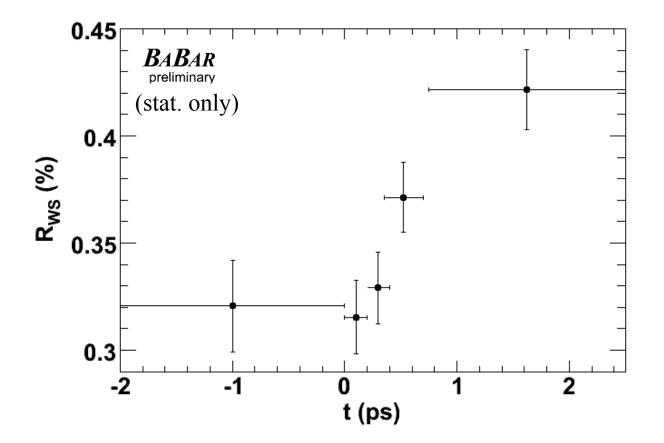
Fit m(Kπ) and Δm in bins of time:
◆If no mixing, ratio of WS to RS signal should be constant
◆No assumptions made on time-evolution of background
◆Each time bin is fit independently

Time bins:

 $\begin{array}{l} -2 < t < 0 \; \mathrm{psec} \\ 0 < t < 0.2 \; \mathrm{psec} \\ 0.2 < t < 0.4 \; \mathrm{psec} \\ 0.4 < t < 0.75 \; \mathrm{psec} \\ 0.75 < t < 2.5 \; \mathrm{psec} \end{array}$

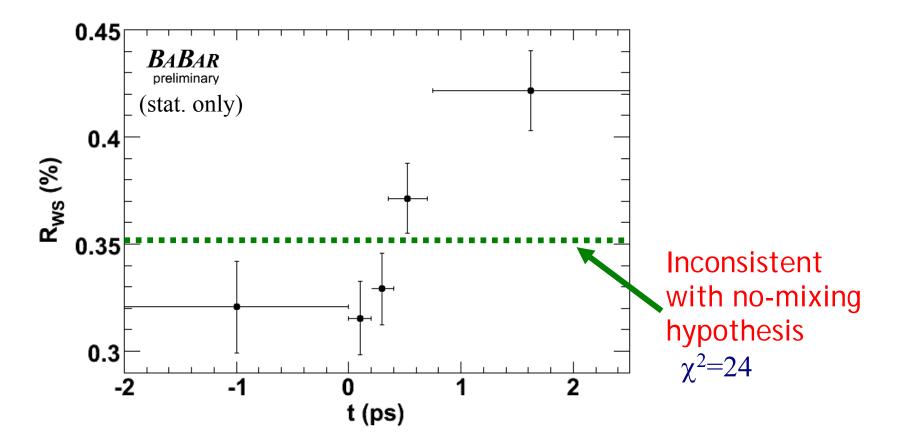


Rate of WS events clearly increase with time:



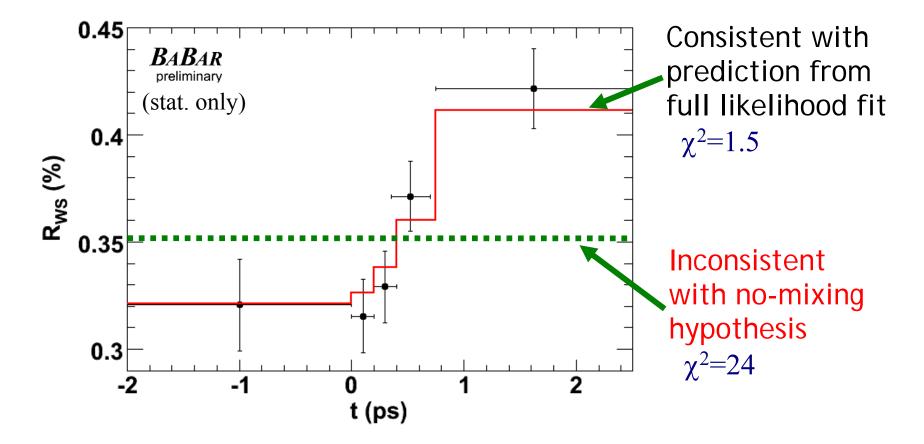


Rate of WS events clearly increase with time:





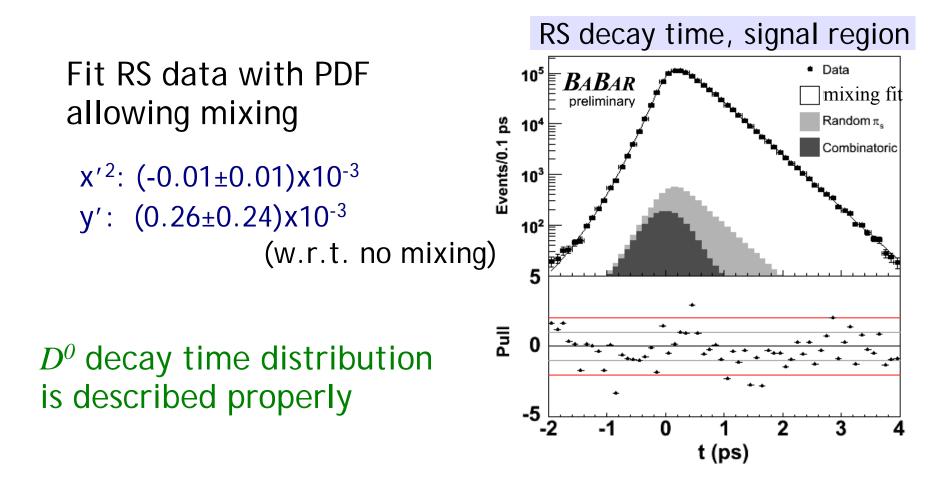
Rate of WS events clearly increase with time:



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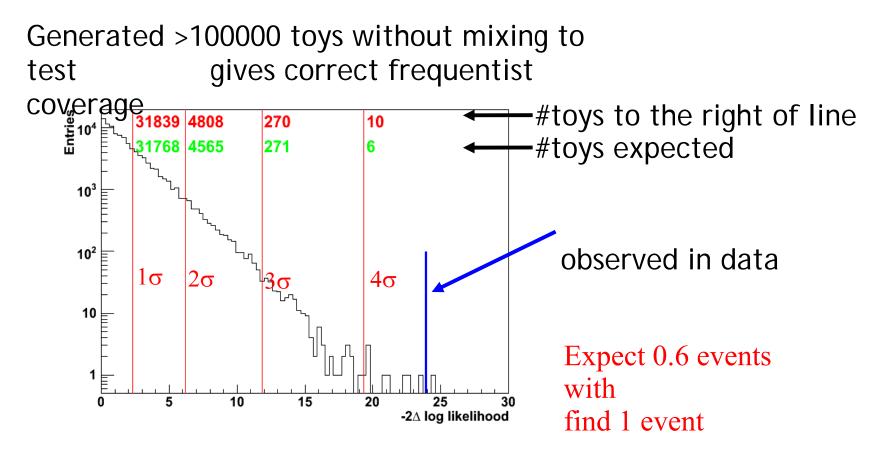
Validation: Fit RS Data for Mixing





Validation: Coverage of $-2\Delta Log \mathscr{Q}$

Significance of signal is calculated as change in log likelihood with respect to no-mixing hypothesis





Systematic Uncertainties

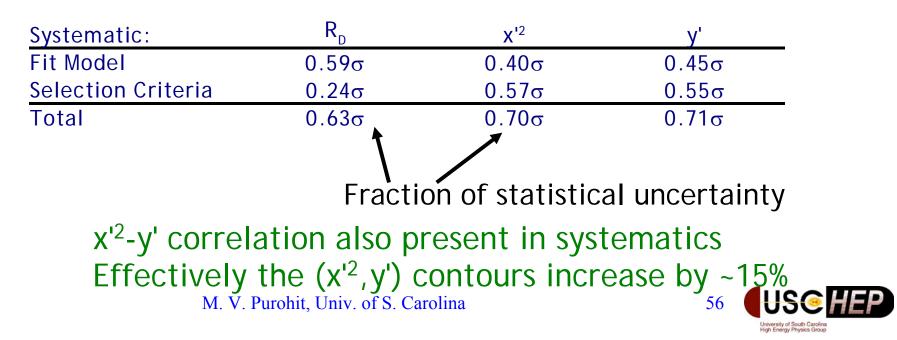
Two types of systematic uncertainties considered:

Fit model variations:

Change signal and background models used in fit, to test assumptions made

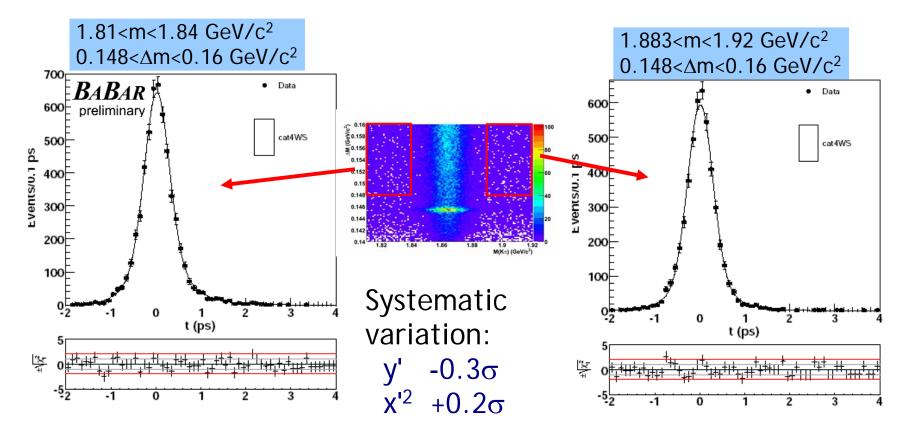
Selection criteria:

Mainly decay time (error) ranges used in fit



Systematic: Combinatorial Decay Time

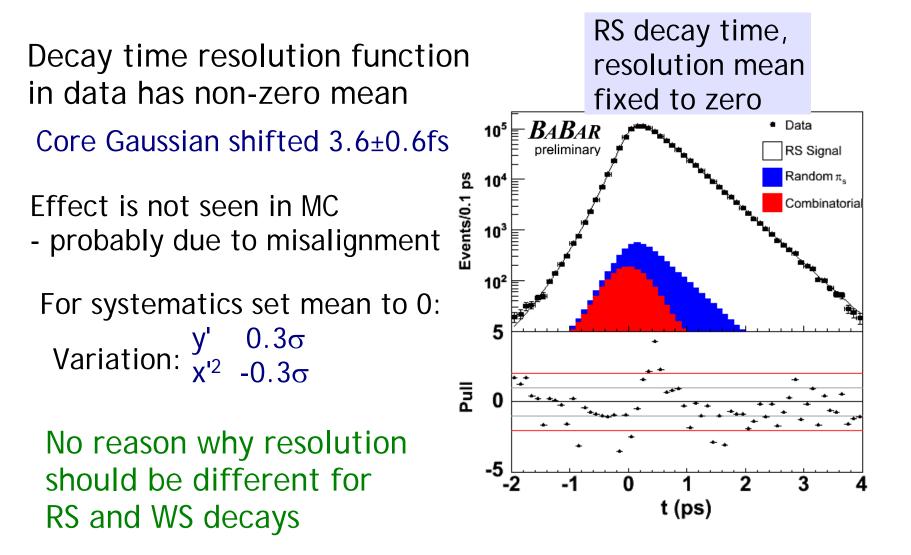
Decay time in combinatorial bkgd not independent of $m(K\pi)$ Fix PDF parameters to fits in different background sidebands:







Systematic: Decay Time Resolution





Allowing for CP Violation

CP violation could introduce different time dependence for D^0 (+) and D^0 (-):

$$\frac{T_{\rm WS}^{\pm}(t)}{e^{-\Gamma t}} = \sqrt{\frac{1 \pm A_{\rm D}}{1 \mp A_{\rm D}}} R_{\rm D} + \sqrt{R_{\rm D}} \sqrt[4]{\frac{(1 \pm A_{\rm D})(1 \pm A_{\rm M})}{(1 \mp A_{\rm D})(1 \mp A_{\rm M})}} (y' \cos \varphi \mp x' \sin \varphi) \Gamma t + \sqrt{\frac{1 \pm A_{\rm M}}{1 \mp A_{\rm M}}} \frac{{x'}^2 + {y'}^2}{4} (\Gamma t)^2$$

Three possible types of CP violation:

- Direct CP violation in DCS decay
- ♦ CP violation in mixing
- CP violation in interference between mixing and decay

Simpler to fit D^0 (+) and D^0 (-) separately:

$$\frac{\Gamma_{WS}^{\pm}(t)}{e^{-t/\tau}} \propto R_D^{\pm} + \sqrt{R_D^{\pm}} y'^{\pm} \left(\frac{t}{\tau}\right) + \left(\frac{x'^{\pm 2} + y'^{\pm 2}}{4}\right) \left(\frac{t}{\tau}\right)^2$$

CP violation if one or more "±" parameters are different



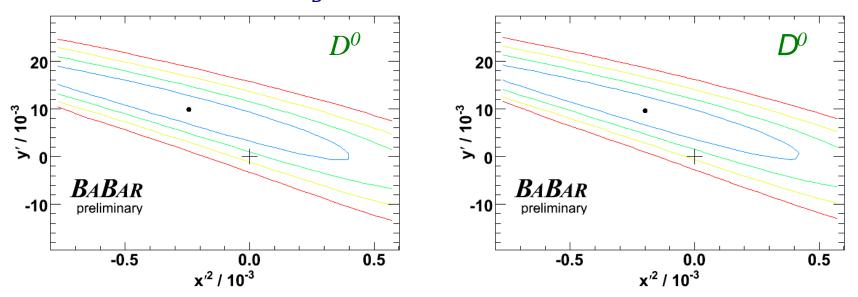
CPV Allowed Contours

<u>Results of fitting D^0 and \overline{D}^0 separately:</u>

 y'^+ : (9.8±6.4±4.5)x10⁻³

 x'^{+2} : (-0.24±0.43±0.30)x10⁻³ x'^{-2} : (-0.20±0.41±0.29)x10⁻³ y'-: (9.6±6.1±4.3)x10-3

 $A_{D} = (-2.1 \pm 5.2 \pm 1.5)\%$



No evidence for CP violation found





Other searches for D^{0} mixing and for CP violation in D^{0} decays



$D^0 - \overline{D}^0$ Mixing in Lifetime Ratio of $D^0 \rightarrow K^+ K^-$, $\pi^+ \pi^- \text{vs } D^0 \rightarrow K^- \pi^+$

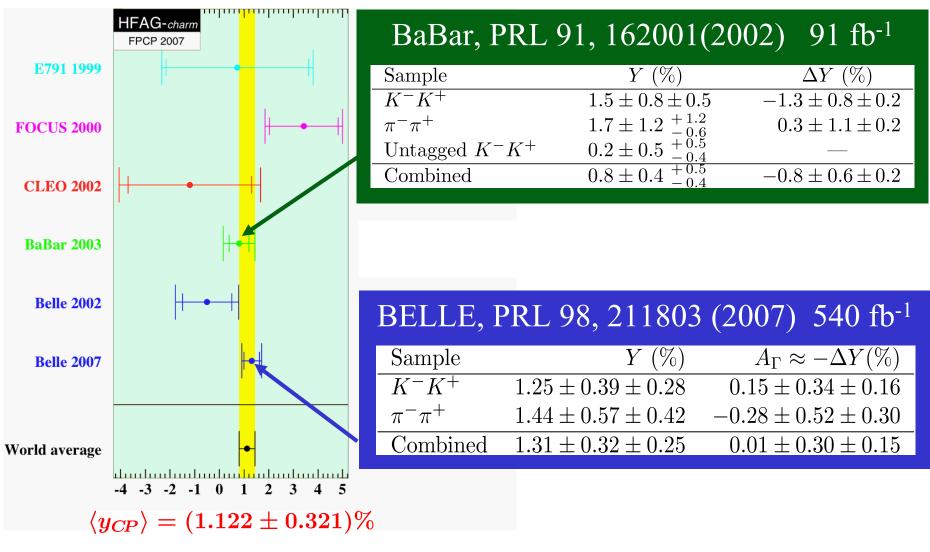
 $D^{0} \rightarrow K^{-} \pi^{+}: \text{CP-mixed} \quad D^{0}(t) \rightarrow K^{+} K^{-}, \ \pi^{+} \pi^{-}: \text{CP-even}$ $Determine the quantities \qquad h = K \text{ or } \pi$ $\langle \tau_{hh} \rangle = (\tau_{hh}^{+} + \tau_{hh}^{-})/2$ $A_{\tau} = (\tau_{hh}^{+} - \tau_{hh}^{-})/(\tau_{hh}^{+} + \tau_{hh}^{-})$ $x \equiv 2 \frac{m_{1} - m_{2}}{\Gamma_{1} + \Gamma_{2}} \qquad y \equiv \frac{\Gamma_{1} - \Gamma_{2}}{\Gamma_{1} + \Gamma_{2}}$ $CPV \text{ in interference} \text{ of mixing and decay:} \qquad \varphi_{f} \equiv \arg\left(\frac{q}{p} \frac{\langle f | \mathcal{H}_{D} | \overline{D}^{0} \rangle}{\langle f | \mathcal{H}_{D} | D^{0} \rangle}\right) \neq 0$

If CP is conserved $y_{CP} = y, \Delta Y = 0$

$$egin{array}{rcl} y_{CP} &=& y\cosarphi_f\ \Delta Y &=& x\sinarphi_f \end{array}$$

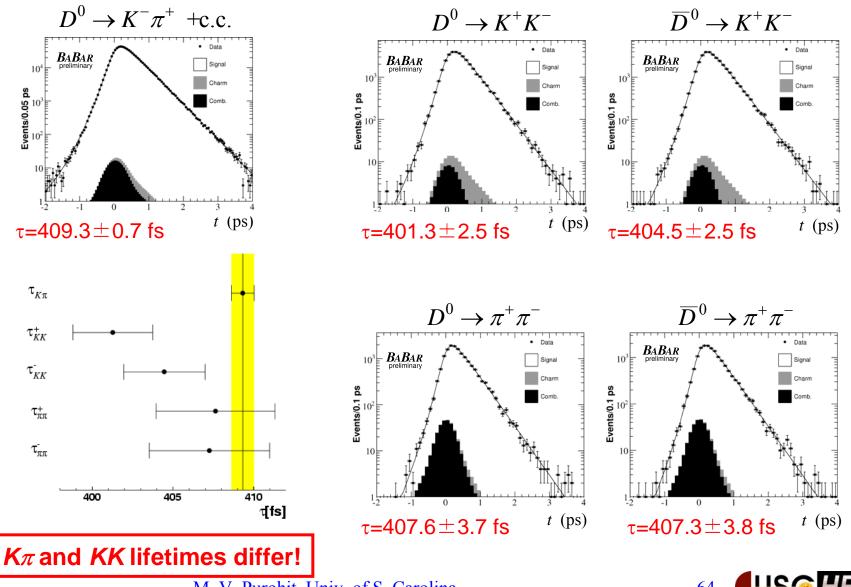


Previous lifetime ratio results





Decay time fits to determine $(y_{CP}, \Delta Y)$

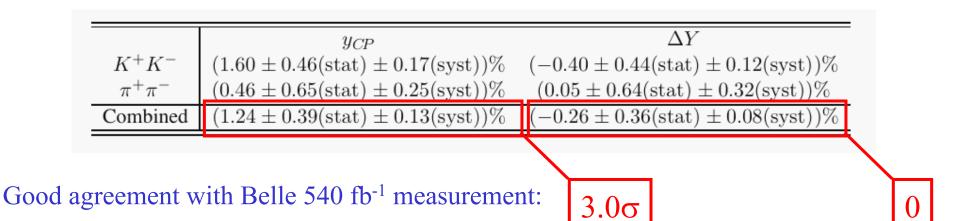


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BaBar (y_{CP} , ΔY) results

Tagged results from 384 fb⁻¹:



 $y_{CP} = (1.31 \pm 0.32 \pm 0.25)\%$

 $A_{\Gamma} = (0.01 \pm 0.30 \pm 0.15)\%$

M. Staric et al. (Belle Collab.), Phys. Rev. Lett. 98, 211803 (2007).



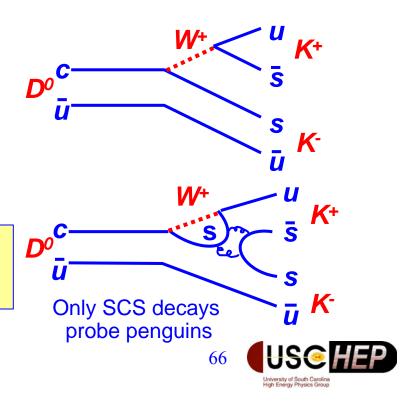
Search for direct CPV in time-integrated $D^0 \rightarrow K^+ K^-$, $\pi^+ \pi^-$ rates $A_{CP} = \frac{\Gamma(f) - \Gamma(\overline{f})}{\Gamma(f) + \Gamma(\overline{f})} = \frac{2 \operatorname{Im} A_1 A_2^* \sin(\delta_1 - \delta_2)}{|A_1|^2 + |A_2|^2 + 2 \operatorname{Re} A_1 A_2^* \cos(\delta_1 - \delta_2)}$ 2 weak amplitudes with phase difference strong phase difference

Two amplitudes with different strong & weak phases needed to observe CPV (in SM from tree and penguins)

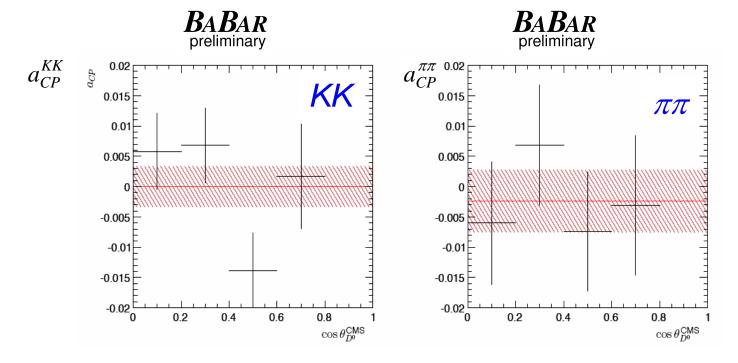
Standard model predictions for direct CPV asymmetries in these modes: O(0.001% - 0.01%)

F. Bucella et al., Phys. Rev. **D51**, 3478 (1995).S. Bianco et al., Riv. Nuovo Cim. 26N7, 1(2003).

e.g., $D^0 \rightarrow K^+K^-$:



Search for CPV in $D^0 \rightarrow K^+ K^-$, $\pi^+ \pi^-$



$$a_{CP}^{KK} = (0.00 \pm 0.34 \text{ (stat.)} \pm 0.13 \text{ (syst.)})\%$$

 $a_{CP}^{\pi\pi} = (-0.24 \pm 0.52 \text{ (stat.)} \pm 0.22 \text{ (syst.)})\%$

No evidence for CPV in either mode



Mixing in $D^0 \rightarrow K^+ \pi^- \pi^0$

Two types of WS Decays:

- Doubly Cabbibo-supressed (DCS)
- Mixing followed by Cabibbo-Favored (CF) decay $D^0 \rightarrow \overline{D}^0 \rightarrow K^+ \pi^- \pi^0$

 $D^{0} \to K^{+} \pi^{-} \pi^{0}$ $D^{0} \xrightarrow{}_{\text{mix}} \overline{D}^{0} \to K^{+} \pi^{-} \pi^{0}$

Two ways to reach same final state \Rightarrow interference!

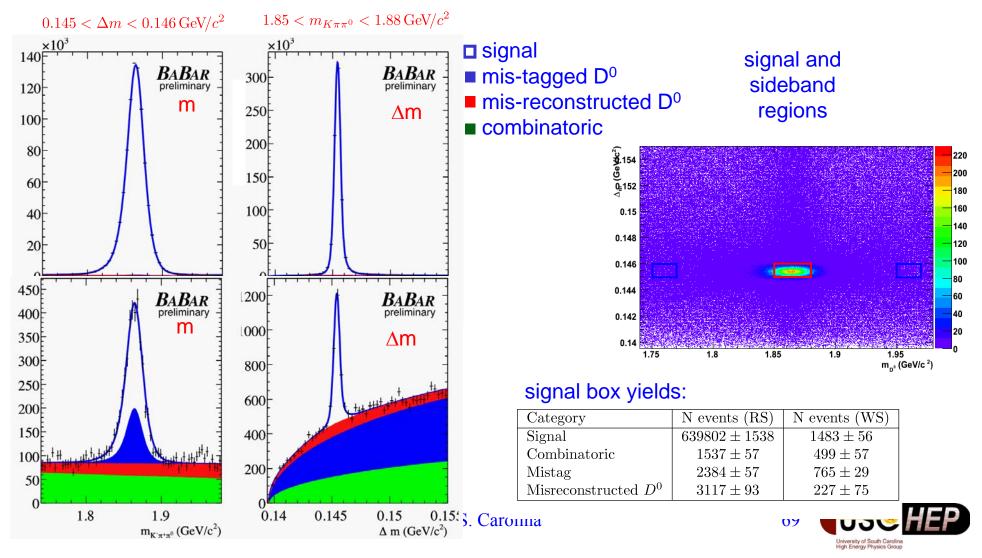
 $\begin{array}{rcl} \text{Time dependent WS rate :} \\ & \Gamma_{\bar{f}}(s_{12}, s_{13}, t) &= e^{-\Gamma t} \{ |A_{\bar{f}}|^2 \leftarrow \text{DCS} \\ \hline & \text{Interference} \rightarrow + & |A_{\bar{f}}| |\bar{A}_{\bar{f}}| \left[y'' \cos \delta_{\bar{f}} - x'' \sin \delta_{\bar{f}} \right] (\Gamma t) \\ & \hline & \text{Mixing} \rightarrow + & \frac{x''^2 + y''^2}{4} |\bar{A}_{\bar{f}}|^2 (\Gamma t)^2 \} \\ & \bar{f} = K^+ \pi^- \pi^0 \qquad A_{\bar{f}} = \langle \bar{f} |\mathcal{H}| D^0 \rangle, \ \bar{A}_{\bar{f}} = \langle \bar{f} |\mathcal{H}| \overline{D}^0 \rangle \end{array}$

 $egin{aligned} y'' &= y\cos\delta_{K\pi\pi^0} - x\sin\delta_{K\pi\pi^0} \ x'' &= x\cos\delta_{K\pi\pi^0} + y\sin\delta_{K\pi\pi^0} \end{aligned}$

 $\delta_{K\pi\pi^0}$: strong phase difference between CF and DCS decay amplitudes

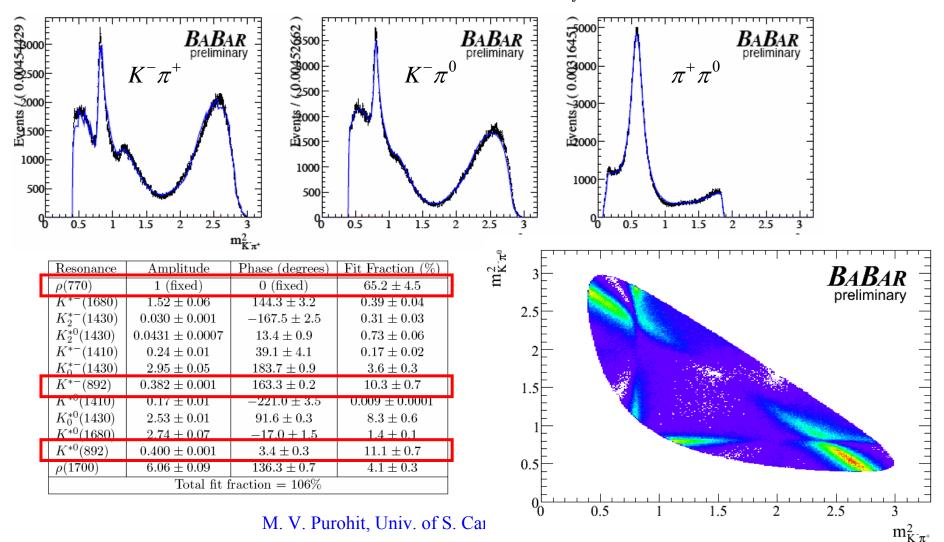


RS and WS $(m_{K\pi\pi} \Delta m)$ fits Determine signal and background yields in subsequent Dalitz analyses.



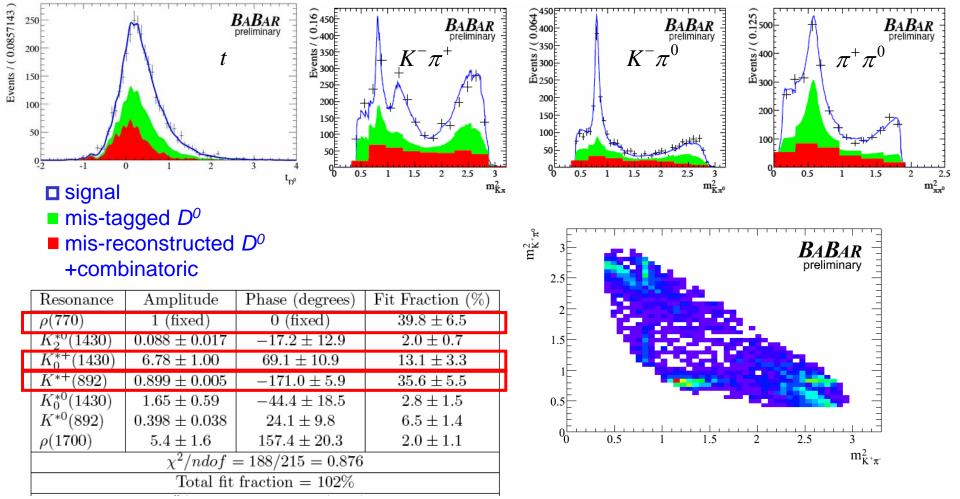
$D^0 \rightarrow K^- \pi^+ \pi^0 \text{RS}$ Dalitz fit

Time-integrated analysis to determine CF amplitudes, $\bar{A}_{\bar{f}}$

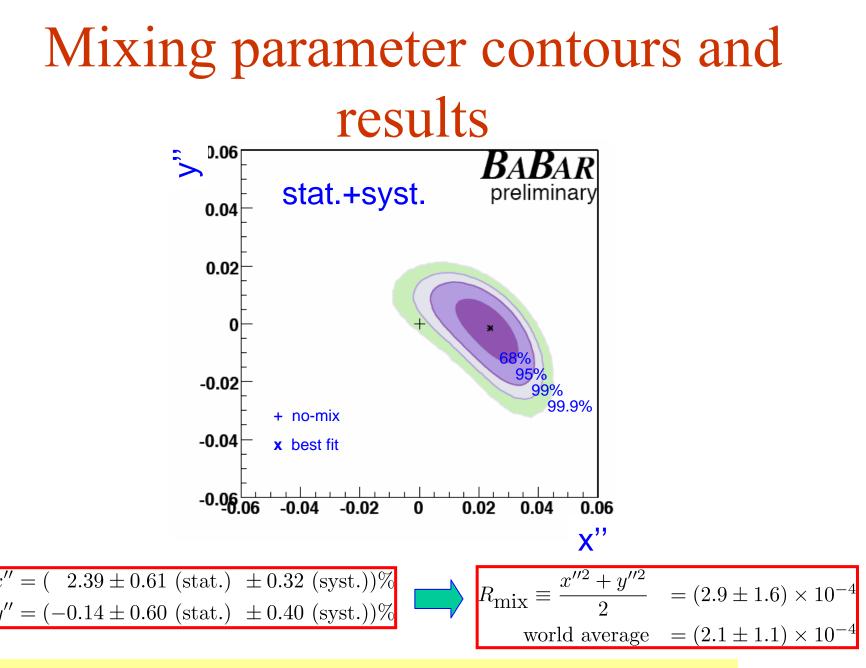


$D^{0}(t) \rightarrow K^{+} \pi^{-} \pi^{0}$ WS Dalitz fit results

Through t-dependence, distinguish DCS amplitudes from the CF amplitudes arising from mixing.







Results are consistent with no mixing at 0.8%, including systematics



BaBar D^0 - \overline{D}^0 Mixing Summary

From $K^{\pm}\pi^{\mp}$ decays:

x'²: (-0.22±0.30±0.21) x 10⁻³, y': (9.7±4.4±3.1) x 10⁻³ Further evidence for $D^{0}-\overline{D}^{0}$ mixing from the *BaBar* experiment:

- $D^0 \rightarrow K^- \pi^+$ to $D^0 \rightarrow K^+ K^-$, $\pi^+ \pi^-$ lifetimes:

 $y_{CP} = (1.24 \pm 0.39 \text{ (stat.)} \pm 0.13 \text{ (syst.)})\%$

 $- D^{0} \rightarrow K^{+} \pi^{-} \pi^{0} \text{ time-dependent Dalitz analysis:}$ $x'' = (2.39 \pm 0.61 \text{ (stat.)} \pm 0.32 \text{ (syst.)})\%$ $y'' = (-0.14 \pm 0.60 \text{ (stat.)} \pm 0.40 \text{ (syst.)})\%$

In $D^0 \rightarrow K^+ K^-$, $\pi^+ \pi^-$ decays,

- no evidence for direct CP violation

 $a_{CP}^{KK} = (0.00 \pm 0.34 \text{ (stat.)} \pm 0.13 \text{ (syst.)})\%$ $a_{CP}^{\pi\pi} = (-0.24 \pm 0.52 \text{ (stat.)} \pm 0.22 \text{ (syst.)})\%$

- no evidence for CP violation in mixing: $\Delta Y = (-0.26 \pm 0.36 \text{ (stat.)} \pm 0.08 \text{ (syst.)})\%$



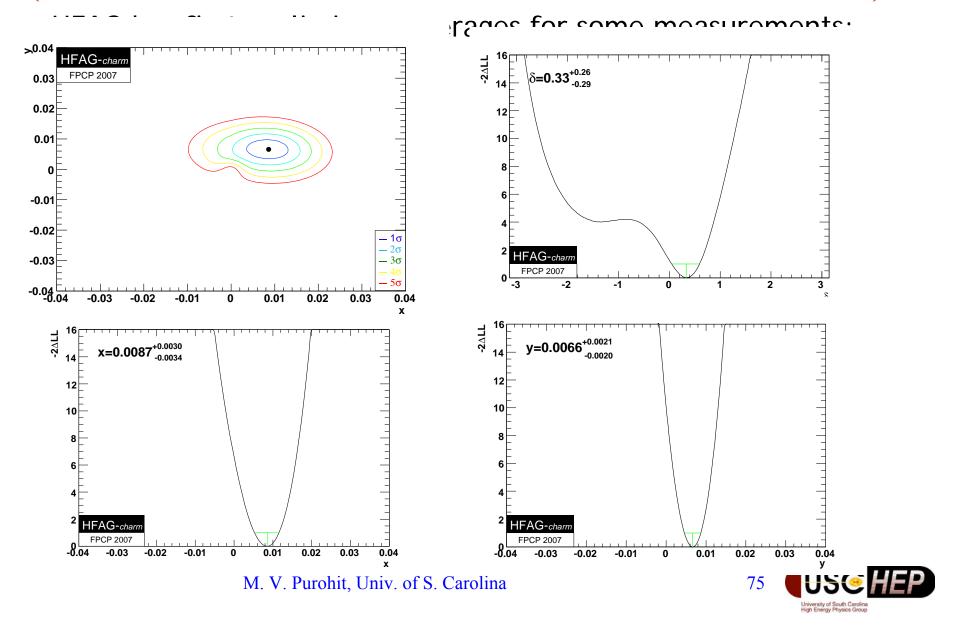


Combining with other results, a comparison with Theory, and Conclusions



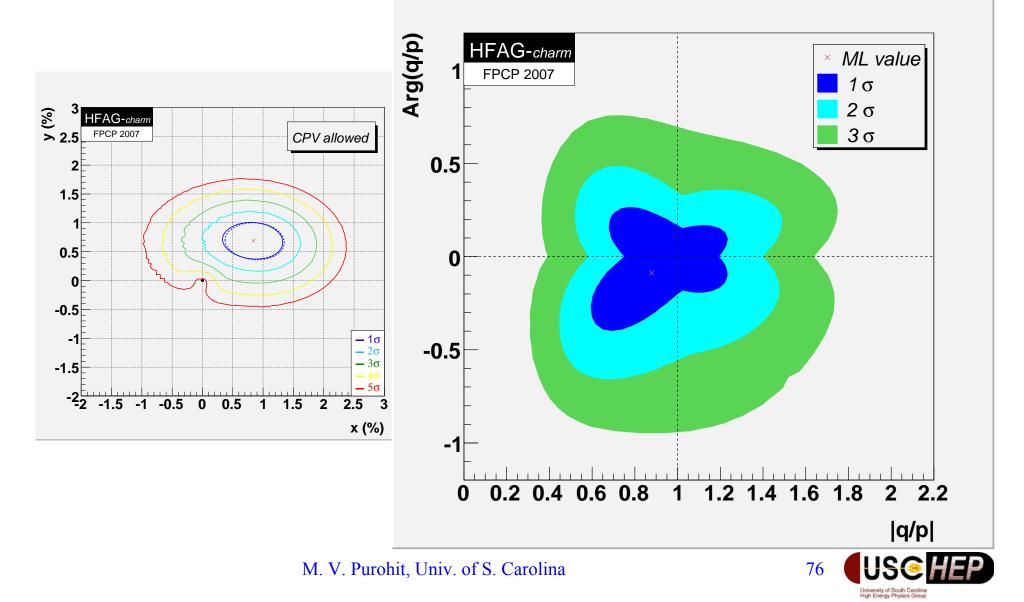
HFAG Results assuming no CPV

(Visit http://www.slac.stanford.edu/xorg/hfag/charm/index.html)



HFAG results allowing for CPV

(Visit <u>http://www.slac.stanford.edu/xorg/hfag/charm/index.html</u>)



Implications of Charm Mixing

BaBar and Belle mixing results first presented at Moriond electroweak conference on March 17

Several new hep-ph preprints on charm mixing since then, e.g.,

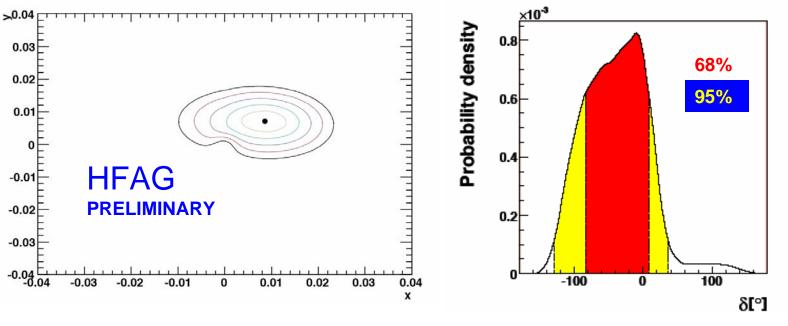
Five use D⁰ mixing results to evaluate limits on:
 ◆Certain SUSY models (flavor suppresion by "alignment") hep-ph/0703204 hep-ph/0703254, arXiv:0704.0601
 ◆Non-universal Z' model
 hep-ph/0703270

"Models are further constrained, "Light non-degenerate but constraints are limited by lack of precise SM value" "Light non-degenerate squarks unlikely to be observed at LHC"

Currently, only an observation of CP violation in mixing would be a clear sign of New Physics



Interpreting the results

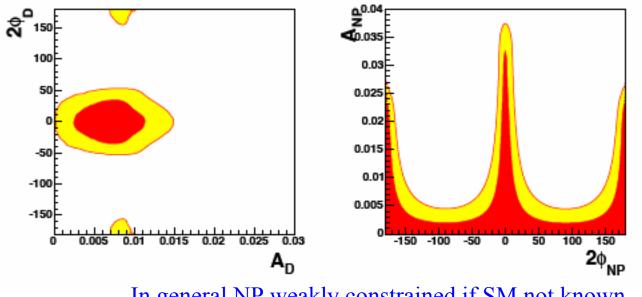


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And CP violation?

In the standard model, $\phi \sim 2\,A^2\lambda^4\eta \lesssim 10^{-3}$



Ciuchini et al. hep-ph/0703294

In general NP weakly constrained if SM not known Nevertheless SUSY coupling can be constrained hints on squark and gluino masses!

Neutral meson mixing always a window into unknown (virtual) states!



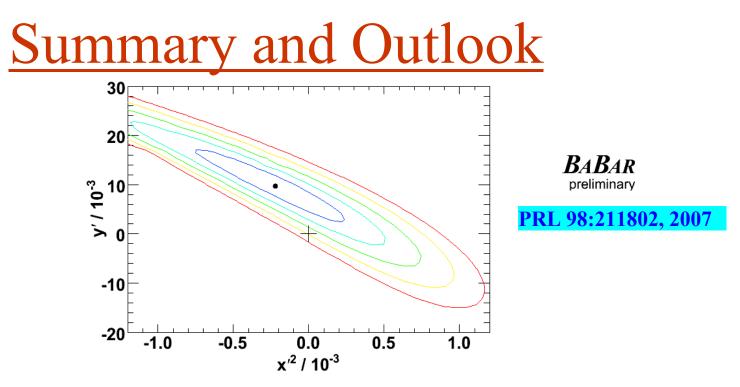
Model	Approximate Constraint		
Fourth Generation (Fig. 2)	$ V_{ub'}V_{cb'} \cdot m_{b'} < 0.5 \; (GeV)$		
Q = -1/3 Singlet Quark (Fig. 4)	$s_2 \cdot m_S < 0.27 \; (\text{GeV})$		
Q = +2/3 Singlet Quark (Fig. 6)	$ \lambda_{uc} < 2.4 \cdot 10^{-4}$		
Little Higgs	Tree: See entry for $Q = -1/3$ Singlet Quark		
	Box: Region of parameter space can reach observed $x_{\rm I}$		
Generic Z' (Fig. 7)	$M_{Z'}/C > 2.2 \cdot 10^3 \text{ TeV}$		
Family Symmetries (Fig. 8)	$m_1/f > 1.2 \cdot 10^3 \text{ TeV}$		
	(with $m_1/m_2 = 0.5$)		
Left-Right Symmetric (Fig. 9)	No constraint		
Alternate Left-Right Symmetric (Fig. 10)	$M_R > 1.2 \text{ TeV} (m_{D_1} = 0.5 \text{ TeV})$		
	$(\Delta m/m_{D_1})/M_R > 0.4 \text{ TeV}^{-1}$		
Vector Leptoquark Bosons (Fig. 11)	$M_{VLQ} > 55(\lambda_{PP}/0.1)$ TeV		
Flavor Conserving Two-Higgs-Doublet (Fig. 13)	No constraint		
Flavor Changing Neutral Higgs (Fig. 15)	$m_H/C > 2.4 \cdot 10^3 \text{ TeV}$		
FC Neutral Higgs (Cheng-Sher ansatz) (Fig. 16)	$m_H/ \Delta_{uc} > 600 \text{ GeV}$		
Scalar Leptoquark Bosons	See entry for RPV SUSY		
Higgsless (Fig. 17)	$M > 100 { m ~TeV}$		
Universal Extra Dimensions	No constraint		
Split Fermion (Fig. 19)	$M/ \Delta y > (6\cdot 10^2 { m ~GeV})$		
Warped Geometries (Fig. 21)	$M_1 > 3.5 { m ~TeV}$		
Minimal Supersymmetric Standard (Fig. 23)	$ (\delta^u_{12})_{\rm LR,RL} < 3.5 \cdot 10^{-2} \mbox{ for } \tilde{m} \sim 1 \mbox{ TeV}$		
	$ (\delta^u_{12})_{\rm LL,RR} < .25$ for $\tilde{m} \sim 1$ TeV		
Supersymmetric Alignment	$\tilde{m} > 2$ TeV		
Supersymmetry with RPV (Fig. 27)	$\lambda_{12k}'\lambda_{11k}'/m_{\bar{d}_{R,k}} < 1.8 \cdot 10^{-3}/100~{\rm GeV}$		
Split Supersymmetry 75			

Table from Golowich, Hewett, Pakvasa and Petrov: **arXiv:0705.3650** [hep-ph]

"... for some models (Split Fermions, Flavor Changing Neutral Higgs) the constraints can be strong."

"Such a list is by nature approximate, and we refer the reader to the body of the paper for a more precise presentation of our results."





BaBar studied $D^0 \rightarrow K\pi$ and other D^0 decays for mixing, CPV Evidence for mixing in $K\pi$ decays (3.9 σ)

 \bullet Evidence for mixing in lifetime differences (3.0 σ)

✤No sign of CP violation at the ~½% level

Consistent with other measurements and SM

More BaBar data and analyses coming up





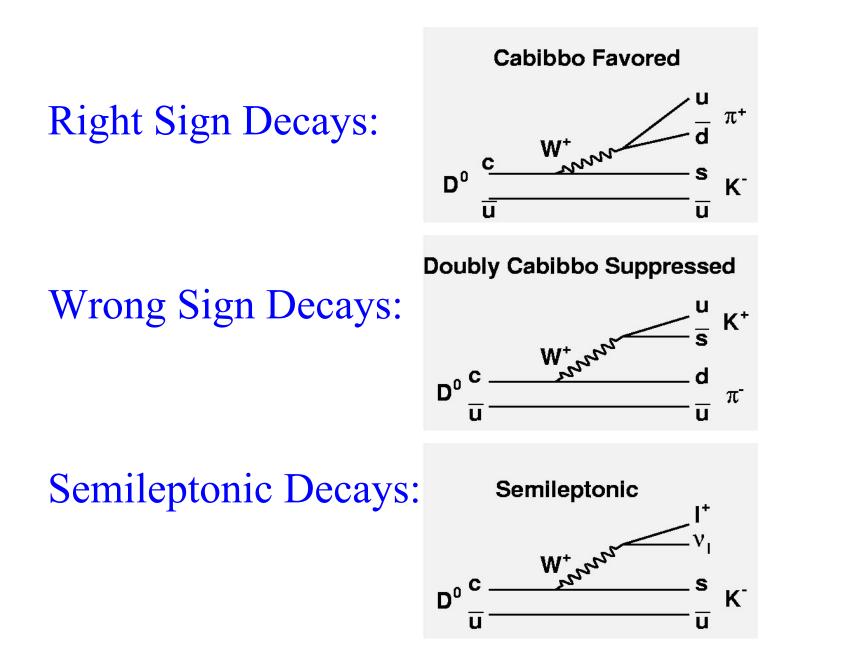
Backup Slides



"Right-sign" and "Wrong-sign" decays

- Most decays of the D⁰ are "Cabibbofavored", e.g., $D^0 \rightarrow K^-\pi^+$.
- Hadronic "wrong-sign" decays (D⁰ → K⁺π⁻ in this case) can occur either via double Cabibbo-suppression (DCS) or due to mixing.
- Semileptonic "wrong-sign" decays only occur due to mixing.







Definitions of x, y, etc.

- The off-diagonal elements of the DDbar mass and decay matrices give rise to mass
 - and lifetime differences:
- $\Delta m \equiv m_1 m_2$, $\Delta \Gamma \equiv \Gamma_1 \Gamma_2$ Γ_1 corresponds to CP even final states as in the decay $D^0 \rightarrow K^-K^+$.
- It is convenient to define

 $x \equiv \Delta m / \Gamma$ and $y \equiv \Delta \Gamma / 2\Gamma$

• When there is a possible strong phase δ between the RS and WS amplitudes, we use instead

 $x' = x \cos \delta + y \sin \delta$ and

 $y' = y \cos \delta - x \sin \delta$



Time dependence of mixing:

- Ordinarily, decays proceed according to the exponential law:
 - $r_{CF}(t) \sim e^{-\Gamma t}$
- However, mixed decays have a modified time dependence:

 $r_{WS}(t) \sim [R_{DCS} + \sqrt{R_{DCS}} y't + (x^2 + y^2)t^2] e^{-\Gamma t}$

• This different time dependence is crucial in separating mixed and DCS contributions to the wrong-sign (WS) rate.



Semileptonic Modes

 In semileptonic modes such as D⁰ → K⁻e⁺v_e we need not worry about DCS backgrounds and simply observe a "wrong-sign" signal, correct for any time-dependent acceptance effects and thereby measure (x² + y²).



Hadronic modes

- If a $D^0 D^0$ bar pair is produced at the ψ'' then the pair must remain coherent until the first decay. Thus, one decay can be used for tagging, while the other (WS) decay is used to measure the rate. This technique was used by Mark III and will be used by CLEO-c.
- In a technique first suggested by Val Fitch, one can use the decay chain D^{*+} → D⁰π⁺ followed by D⁰ → K⁻π⁺. The Q in the D^{*+} decay is so small that combinatoric backgrounds are kinematically suppressed. The "slow pion" from the D^{*+} tags the flavor of the D⁰ at birth.



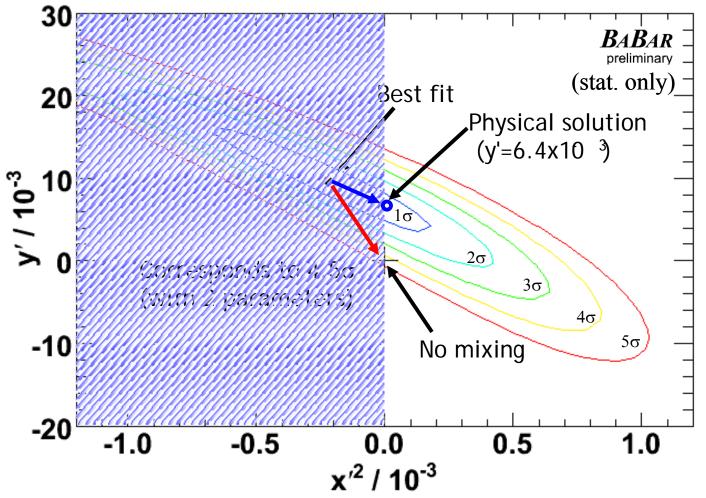
The lifetime difference technique

- First suggested by Ted Liu, in this technique one simply measures the lifetime of the D⁰ in modes such as K⁻K⁺, which are CP eigenstates and measure Γ₁, and in copious modes such as K⁻π⁺ which yield an average of Γ₁ and Γ₂.
- Then $y \equiv \Delta \Gamma / 2\Gamma \cong (\tau_{K\pi} / \tau_{KK}) 1$



Signal Significance

Best fit is in unphysical region (x²<0)



M. V. Purohit, Univ. of S. Carolina



Systematic Uncertainties

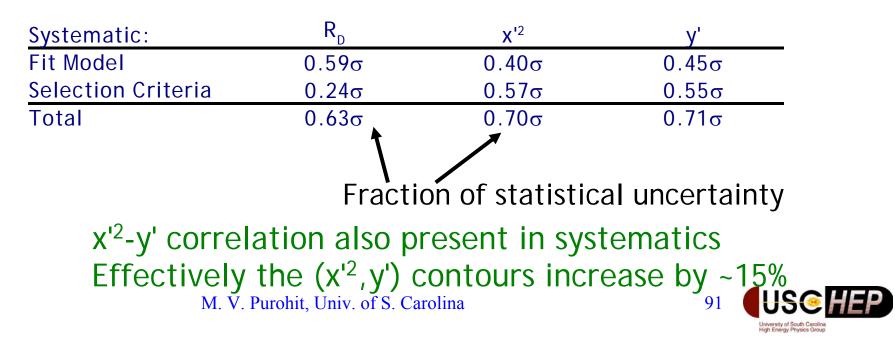
Two types of systematic uncertainties considered:

Fit model variations:

Change signal and background models used in fit, to test assumptions made

Selection criteria:

✤Mainly decay time (error) ranges used in fit



Double tag at v(3770) [CLEO-c]

- Reconstruct Double Tags: CP vs Kπ
- Asymmetry in CP+ vs CP- related to cosδ

D_{CP±} neutral D CP eigenstate

$$A = \frac{B(D_{CP+} \rightarrow K^{-}\pi^{+}) - B(D_{CP-} \rightarrow K^{-}\pi^{+})}{B(D_{CP+} \rightarrow K^{-}\pi^{+}) + B(D_{CP-} \rightarrow K^{-}\pi^{+})}$$

R_D is ratio of DCS to Cabibbo favored rates

$$\cos \delta = \frac{A}{2\sqrt{R_D}}$$

• Input $R_D = (3.60 \pm 0.08)\%$ from PDG2006+CDF ~±2%,

 $\psi(3770)$ decay conserves CP

Need to run On threshold

- Updated results with 281 pb⁻¹ at Winter Conferences
 - Expect $\sigma(y)$ ~ ±1.5% and $\sigma(\cos \delta_{K\pi})$ ~ ±0.3
 - Including systematic uncertainties
- Full CLEO-c dataset ~750 pb⁻¹
 - Expect $\sigma(y)$ ~ ±1.0% and $\sigma(\cos \delta_{K\pi})$ ~ ±0.1-0.2



BaBar (y_{CP} , ΔY) systematics

Systematic uncertainties (%):

Systematic	Δy_{CP}^{KK}	$\Delta y_{CP}^{\pi\pi}$	Δy_{CP}	$\Delta(\Delta Y^{KK})$	$\Delta(\Delta Y^{\pi\pi})$	$\Delta(\Delta Y)$
Signal Model	0.130	0.059	0.085	0.072	0.265	0.062
Charm Bkgd	0.062	0.037	0.043	0.001	0.002	0.001
Combinatorial Bkgd	0.019	0.142	0.045	0.001	0.005	0.002
Selection	0.068	0.178	0.046	0.083	0.172	0.011
Detector Model	0.064	0.080	0.064	0.054	0.040	0.054
Quadrature sum	0.172	0.251	0.132	0.122	0.318	0.083

Variations:

- Signal: PDF shape, polar angle dependent resolution offset, signal interval
- Charm backgrounds: yields and charm lifetime
- Combinatorial backgrounds: yields, shape and sideband region
- Selection: σ_t criterion, treatment of multiple candidates
- Detector: Alignment and energy loss



Search for CPV in $D^0 \rightarrow K^+ K^-$, $\pi^+ \pi^-$

Measure the time integrated CP asymmetries

$$a_{CP}^{KK} = rac{\Gamma(D^0
ightarrow K^- K^+) - \Gamma(\overline{D}{}^0
ightarrow K^+ K^-)}{\Gamma(D^0
ightarrow K^- K^+) + \Gamma(\overline{D}{}^0
ightarrow K^+ K^-)} \ a_{CP}^{\pi\pi} = rac{\Gamma(D^0
ightarrow \pi^- \pi^+) - \Gamma(\overline{D}{}^0
ightarrow \pi^+ \pi^-)}{\Gamma(D^0
ightarrow \pi^- \pi^+) + \Gamma(\overline{D}{}^0
ightarrow \pi^+ \pi^-)}.$$

Experimental procedure:

- fit $m, \Delta m$ distributions to determine raw signal weights
- Determine relative D^0/\overline{D}^0 soft pion tagging efficiency using $\overline{D}^0 \rightarrow K\pi$ data

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\Rightarrow greatly	reduces	systematic	uncertainties-	~
		✓		- 2.

Category	Δa_{CP}^{KK}	$\Delta a_{CP}^{\pi\pi}$
2-Dim. PDF shapes	$\pm 0.04\%$	$\pm 0.05\%$
π_s correction		$\pm 0.08\%$
a_{CP} extraction	$\pm 0.09\%$	$\pm 0.20\%$
Quadrature sum	$\pm 0.13\%$	$\pm 0.22\%$



