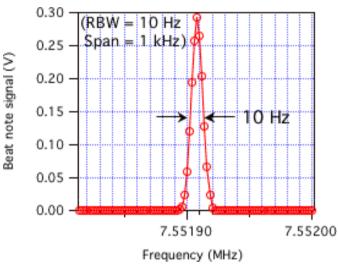
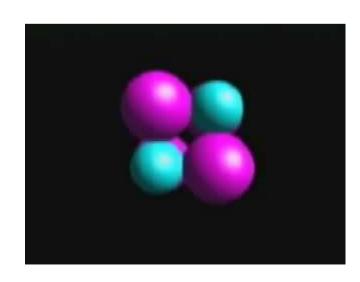
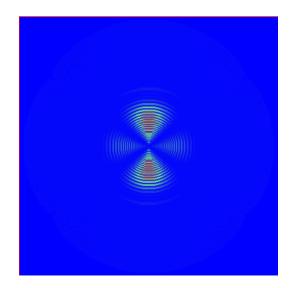
An Introduction to AMO Physics Output Output Displayed to the AMO Output O

Cass Sackett UVA UVA







What is AMO?

= Atomic, molecular, and optical physics

Linked because share basic approach:

- Study "simple" quantum systems with high precision
- Learn to control systems for applications
- Central tool: laser

Purpose of Talk

Introduce key background ideas of AMO Helpful for subsequent talks

Aimed at new students

Outline

- I. Atoms and molecules
 Hierarchy of structures, notation
- II. Driving transitions

 Transition rate, saturation, coherent evolution
- III. Lasers
 Basic theory, types of laser
- IV. Frequency conversionNonlinear response, phase matching
- V. AMO at UVA Groups, courses

Atoms

Simplest atom: hydrogen

Energy levels have hierarchy

Highest level: Coulomb $E_n = -\frac{mc^2}{2n^2}\alpha^2$

$$E_n = -\frac{mc^2}{2n^2}\alpha^2$$

fine structure constant
$$\alpha = \frac{e^2}{4\pi\epsilon_0 \hbar c} = \frac{1}{137}$$

Often write $E_n = -\frac{1}{2n^2}$ using atomic units unit of energy = $mc^2\alpha^2$ = Hartree = 27.2 eV

Label states with n, l: 3d state has n = 3, l = 2

Fine structure

Spin orbit coupling correlates L and S
Use J = L +S

Relativistic effects shift energy levels $\Delta E/E \sim (Z\alpha)^2$ for atomic number Z

More important for heavier atoms

Energy levels split by J

3d
$$\longrightarrow$$
 $J = 5/2$ (6 states)
(10 states) $J = 3/2$ (4 states)

Notation: n ^{2S+1}L_J

So " $3^{2}D_{3/2}$ " means n = 3, S = $\frac{1}{2}$, L = 2, J = $\frac{3}{2}$

Hyperfine structure

Electronic angular momentum **J** coupled to nuclear angular momentum **I**

 \Rightarrow Total momentum $\mathbf{F} = \mathbf{I} + \mathbf{J}$

Hydrogen has $I = \frac{1}{2}$:

$$3^{2}D_{3/2} \longrightarrow F = 2$$

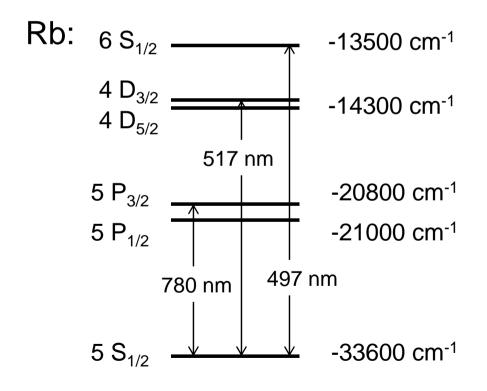
$$F = 1$$

Energy scale
$$\approx \frac{\text{nuclear moment}}{\text{electron moment}} \times \text{FS scale}$$

Typically about 10⁻³

Other atoms

Alkali atoms (Li, Na, K, Rb, Cs)
have one valence electron + core electrons
Act like H, but levels shifted due to core



units cm⁻¹ = wave numbers 1 eV = 8050 cm⁻¹ = inverse of corresponding wavelength

Multielectron Atoms

More than one valence electron = more complex Multiple s_i 's and ℓ_i 's, couple in different ways

Typically
$$L = \sum \ell_i$$
 $S = \sum s_i$ $J = L + S$

$$S = \sum S_i$$

$$J = L + S$$

Label state with L, S and J:

Calcium: $^{1}S_{0}$ (4s5s) $\overline{\text{A}}$ 33300 cm⁻¹ 31500 $^{3}S_{0}$ (4s5s) 300 nm 20300 \longrightarrow $^{3}D_{1,2,3}$ (4s3d) 15300 \longrightarrow $^{3}P_{0,1,2}$ (4s4p) 422 nm 653 nm

Molecules

Put two atoms together, even more fun!

Interact via molecular potential V(R)

= energy as function of nuclear separation R

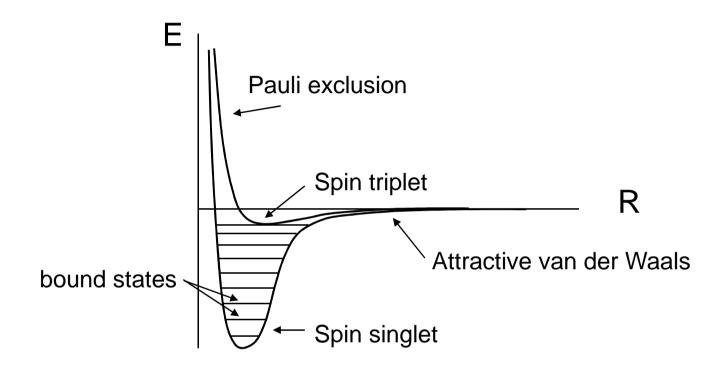
Born-Oppenheimer approximation:

- Fix nuclei at spacing R
- Find electronic energies $E_n(R)$
- Use $E_n(R)$ as potential for nuclear motion

Usually valid $(m_e << m_N)$

At large R, $E_n(R) \rightarrow E_{n1} + E_{n2}$ atomic energies lots of quantum numbers!

Alkali diatom ground state potential:



Nuclei move in V(r): have own eigenstates and energies vibrational and rotational quantum numbers

Transitions

What do we do with atoms and molecules?

Typically, drive transitions using perturbation

$$H' = \hat{V} \cos \omega t$$

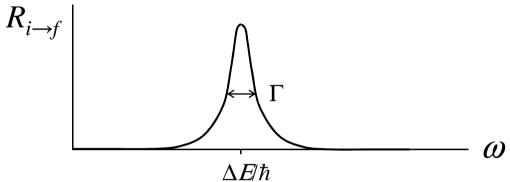
Basic result: Fermi's Golden Rule:

$$R_{i\rightarrow f} = \frac{2\pi}{\hbar} \left| \left\langle f \left| \hat{V} \right| i \right\rangle \right|^2 \delta(\hbar\omega - \Delta E)$$
 where
$$\left| i \right\rangle = \text{initial state}$$

$$\left| f \right\rangle = \text{final state}$$

$$\Delta E = \left| E_f - E_i \right|$$

In real life, delta function \rightarrow peak with finite line width $\hbar\Gamma$



Various reasons:

- Finite lifetime τ : $\Gamma \approx \tau^{-1}$
- Thermal atoms at temperature T: $\Gamma \approx \omega \sqrt{\frac{k_B T}{mc^2}}$ (Doppler broadening)
- Collision time t_c : $\Gamma \approx t_c^{-1}$
- Finite duration of drive pulse t_p : $\Gamma \approx t_p^{-1}$

On resonance $\hbar\omega = \Delta E$,

$$\delta(\hbar\omega - \Delta E) \rightarrow \frac{1}{\hbar\Gamma}$$
 and $R = \frac{2\pi}{\Gamma} \frac{\left|V_{if}\right|^2}{\hbar^2}$

Typically V = interaction with laser Strongest interaction: dipole coupling $V = e\mathbf{E} \cdot \mathbf{r}$

Dipole coupling

Only couples certain states: selection rules

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\Delta F, \Delta m_F = 0, \pm 1 (photon has \ell = 1)
\Delta S = 0 (no coupling to spin)
\Delta L = \pm 1 (require opposite parity)
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Can violate with higher order interactions (magnetic dipole, electric quadrupole) or higher order perturbations (two-photon transition)

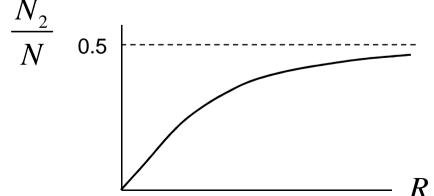
Saturation

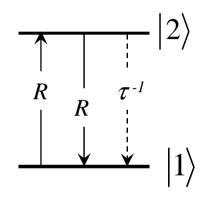
Easy to see that $R_{i op f} = R_{f op i}$ Transition goes both ways This limits population transfer

Simple model:
$$N_1 + N_2 = N$$

$$\frac{dN_2}{dt} = +RN_1 - RN_2 - \frac{N_1}{\tau}$$

In steady state, get $N_2 = N \frac{R\tau}{2R\tau + 1}$





Population saturates

Coherent Excitation

If $R_{i\rightarrow f} \ge \Gamma$ and drive field is \approx monochromatic, Fermi's golden rule is inadequate

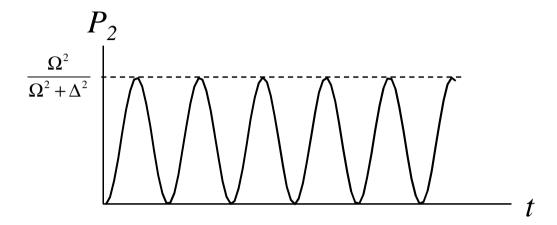
Need full solution to Schrodinger Eqn

Solution simple if just two levels coupled and if

$$\begin{split} \Omega \equiv & \frac{V_{if}}{\hbar} \square \quad \omega_0 \qquad \left| \Delta \right| \equiv \left| \omega - \omega_0 \right| \square \quad \omega_0 \\ \text{with } \omega_0 \equiv & \frac{E_f - E_i}{\hbar} \qquad \quad \Omega \text{ called Rabi frequency} \\ \Delta \text{ called detuning} \end{split}$$

Get Rabi oscillation:
$$P_2 = \frac{\Omega^2}{\Omega^2 + \Delta^2} \sin\left(\frac{\sqrt{\Omega^2 + \Delta^2}}{2}t\right)$$

Rabi Oscillation



If
$$\Delta = 0$$
, get perfect transfer at $t = \frac{\pi}{\Omega}$ called " π -pulse"

Make superposition
$$|\psi\rangle = \frac{1}{\sqrt{2}} (|1\rangle + |2\rangle)$$
 at $t = \frac{\pi}{2\Omega}$ called " $\pi/2$ -pulse"

Oscillation damps on time scale Γ^{-1} Approaches steady state value from saturation

Lasers

So what do we drive these transitions with?
Sometimes microwaves
More often lasers

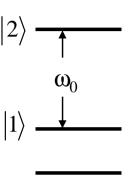
Lasers are basic tool for much modern physics Nearly all AMO physics

Involves three ingredients:

- Gain medium
- Pumping mechanism
- Mirrors

Gain medium = collection of atoms, molecules, etc with well a defined transition _____

Relatively narrow linewidth, long excited state life time



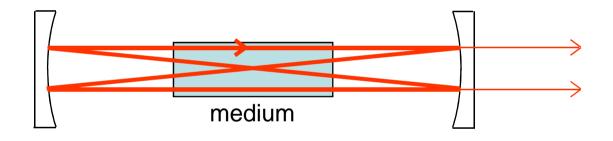
Pumping = energy source driving atoms to upper state Need inversion: $N_2 > N_1$

Easiest to achieve in "four level" system

$$|2\rangle \xrightarrow{\tau \downarrow \uparrow} |2\rangle \xrightarrow{\omega_0 R_{\text{pum}}} |1\rangle \xrightarrow{\tau'\downarrow} |2\rangle$$

As long as $N_2 > N_I$, light at ω drives $2 \to 1$ more than $1 \to 2$ Number of photons increases: gain!

Mirrors: light at ω passes through medium many times builds up to high intensity



Eventually saturates transition: gain \rightarrow 0 laser has steady state output

Types of lasers

Most important distinction: pulsed vs. continuous

Pulsed laser: output on for brief time

Various methods:

- Switch pumping on/off (flash pumped)
- Add "shutter" to cavity (Q-switched)
- Use nonlinear effects (mode-locked)

Can get short pulse durations:

- picoseconds (Q-switched)
- femtoseconds (mode-locked)

Typically use as strobe or marker to study fast events (chemical reactions, electronic motion, etc.)

Get high peak power

Continuous (CW) laser: on all the time

Typically frequency is very stable

Frequency determined by mirrors: need $n\lambda = L$ L = round trip path length

Put one mirror on piezoelectric translator "lock" to Fabry-Perot cavity and/or atomic transition

Readily get $\Delta f \sim$ few MHz As low as 1 Hz

Use for spectroscopy, laser cooling

Most types of laser can be pulsed or CW, depending on setup

Common lasers

Classify by gain medium:

- Solid state (Ti-Sapphire, Alexandrite) doped crystal
 Fiber laser doped optical fiber
- Gas laser (HeNe, Argon ion, CO₂) gas discharge
 Excimer laser reactive gases
- Dye laser liquid dye
- Diode laser semiconductor

Wavelengths: incomplete coverage from UV to IR

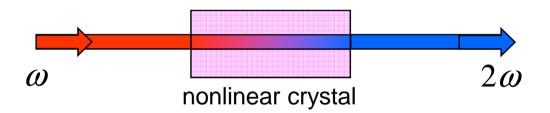
Power: typically 1-10 W average output lower for diode lasers, some gas lasers

Frequency Conversion

Incomplete wavelength coverage is annoying. Often no laser at the frequency you need

Solution: use nonlinear optics to change ω

Simplest example: second harmonic generation = frequency doubling



Can also combine, split, or mix frequencies:

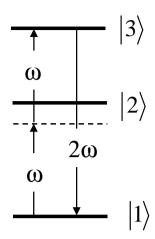
$$\omega_1 + \omega_2 \rightarrow \omega_3$$

$$\omega_3 \rightarrow \omega_2 + \omega_1$$

$$\omega_1 + \omega_2 \rightarrow \omega_3$$
 $\omega_3 \rightarrow \omega_2 + \omega_1$ $\omega_1 + \omega_2 \rightarrow \omega_3 + \omega_4$

Why does it work?

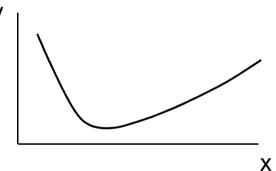
Quantum picture:
Off-resonant two-photon transition



Classical picture:

Electron in anharmonic potential

Response to drive at ω :



 $\begin{array}{c} \times \\ \\ \\ \end{array} \begin{array}{c} \\ \\ \end{array} \begin{array}{c}$

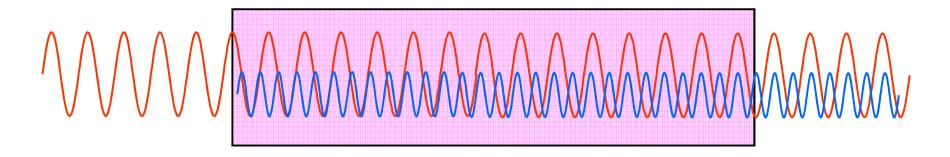
Component at 2ω

Why don't you see this all the time?

- 1. Need medium with large nonlinearity
- 2. Need high intensity at ω
- 3. Need phase matching

Phase matching:

Harmonic wave produced all along crystal



Waves travel at different speeds: $\lambda_1 \neq 2\lambda_2$ since $n = n(\omega)$

→ Waves get out of phase

Wave produced at one place cancels wave from another Net output small

A few ways to solve:

- Select and adjust crystal so $n(\omega) = n(2\omega)$ "phase matching"
- Engineer periodic structure in crystal to keep phase "periodically poled crystal"
- Use a thin crystal
 Need very high intensity → pulsed laser

Some common crystals: KDP, KTP, LBO, BBO, LiNbO₃

Can typically get ~10% conversion efficiency depends a lot on medium, phase matching, input power

AMO at UVA

Groups (in physics):

Gallagher: Rydberg atoms

- Microwave spectroscopy and control
- Multi-electron excitations
- Interactions between atoms

Bloomfield: Clusters

- Magnetism
- Structure and isomerization

Jones: Rydberg atoms

- Observation of electronic, nuclear motion
- Control of electronic, nuclear motion
- Simulation of collisions
- Ultrafast laser technology

Lehman: Spectroscopy (chemistry dept.)

- Molecular spectroscopy
- Molecular dynamics
- Molecules in He droplets

Pfister: Quantum optics

- Nonclassical light generation and characterization
- Quantum information
- Quantum measurements

Sackett: Bose-Einstein condensation

- Laser cooling
- Atom interferometry

Also

Arnold: Particle physics theory

- BEC phase transition

Cates: Nuclear and medical physics

- Optical pumping techniques

Kolomeisky: Condensed matter physics theory

- BEC theory
- Atomic properties

Pate: Chemistry

- Molecular dynanics
- Highly excited molecules

AMO Classes

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Phys 531 – Optics
Phys 532/822 – Photonics
Phys 842 – Atomic Physics
Phys 826 – Ultrafast Lasers
Phys 888 – Quantum Optics

Alternate years
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Also useful:

Phys 519 – Electronics

Phys 553 – Computational Physics

Other departments:

ECE 541 – Optics and Lasers

MAE 687 – Applied Engineering Optics

AMO Journals

General science:

Science, Nature

General physics:

Phys. Rev. Lett, Nature Physics

AMO specific:

Phys. Rev. A, J. Phys. B

AMO overlaps with physical chemistry:

Chem. Phys. Lett., J. Chem. Phys.

and optics:

Optics Lett., JOSA A, JOSA B, Applied Optics

General interest physics:

Physics Today – free with APS membership!

Good to look over periodically!

Conclusions

Tried to provided introduction/reminder about a few key ideas in AMO physics

- Atomic and molecular states
- Transitions
- Lasers and frequency conversion

and introduced opportunities for AMO here

As we have more seminars, keep these ideas in mind