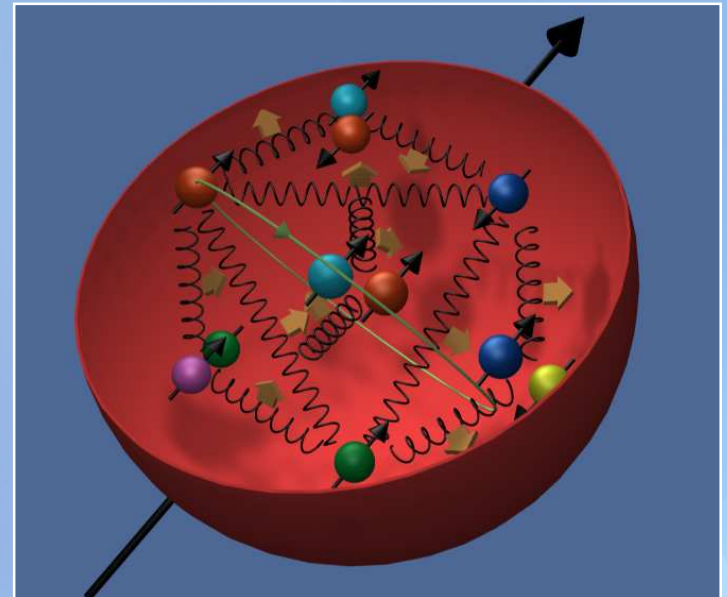
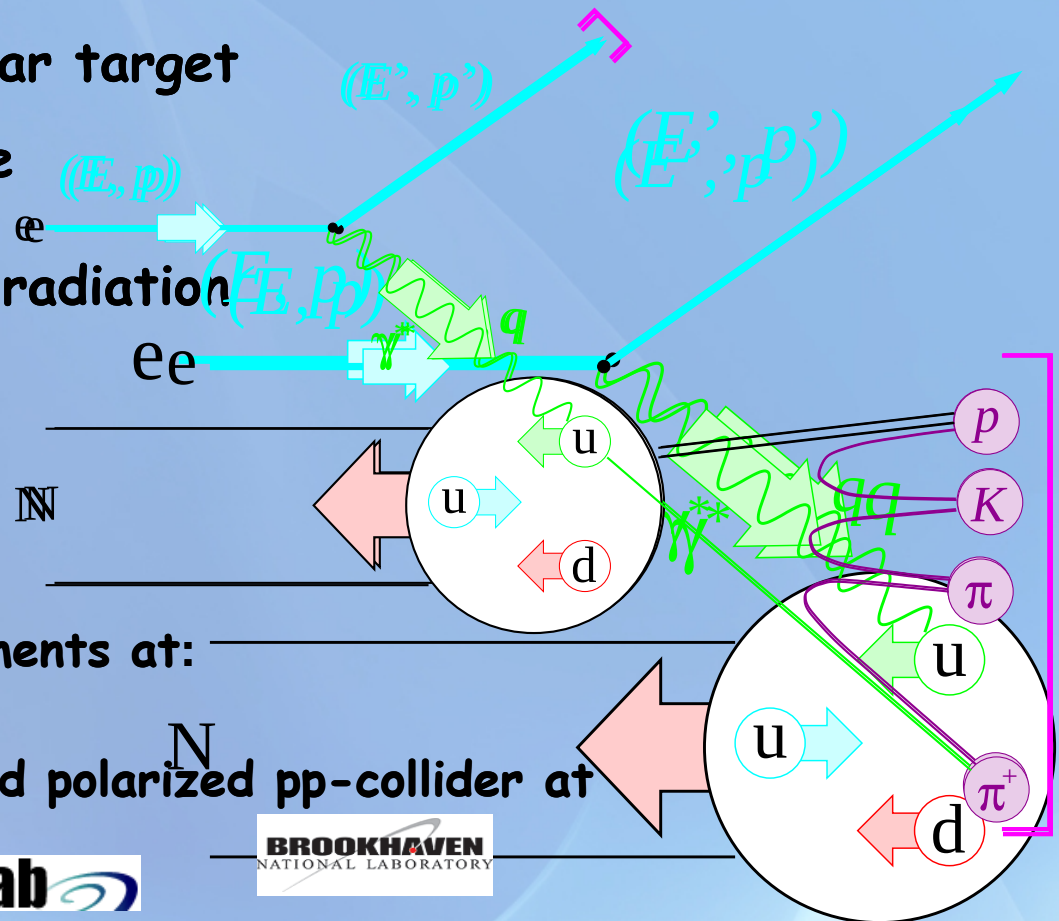


The Spin Structure of the Nucleon



Experimental Prerequisites

- longitudinally polarized lepton beam
- longitudinally polarized nuclear target
- large geometrical acceptance
 - small angles
limited by synchrotron radiation
 - big angles
limited by money
- good particle identification



So far only **"fixed target"** experiments at:



Jefferson Lab



and polarized N pp-collider at



E.C. Aschenauer

SFU, VanCouver, March 2007

tempora

Jefferson Lab
CLAS Detector

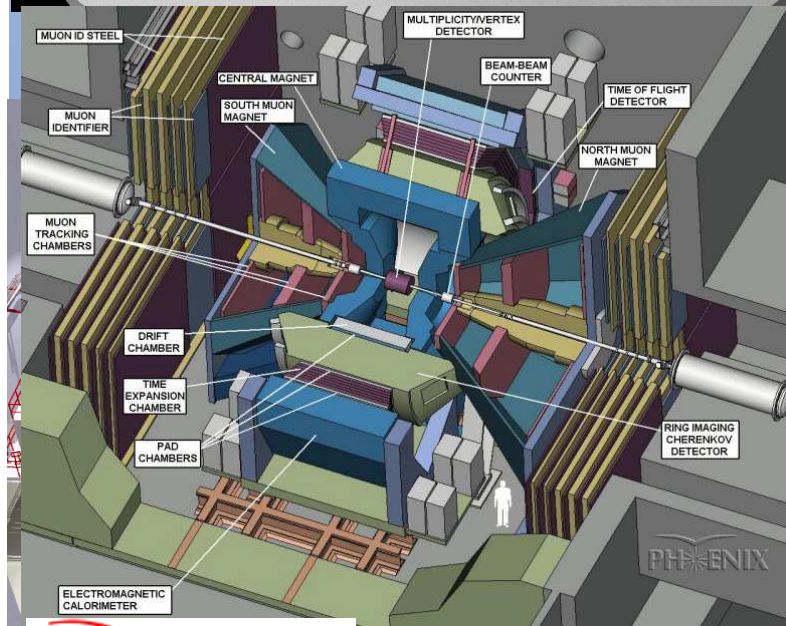
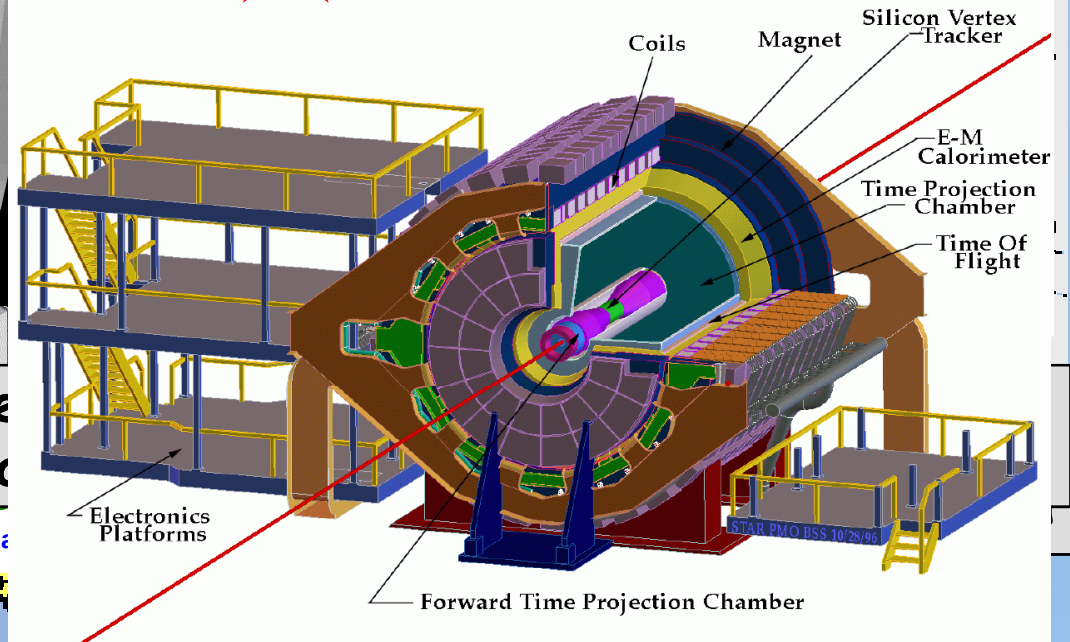
Hall A

Hall B



STAR Detector

Two high-resolution
4 GeV spectrometers



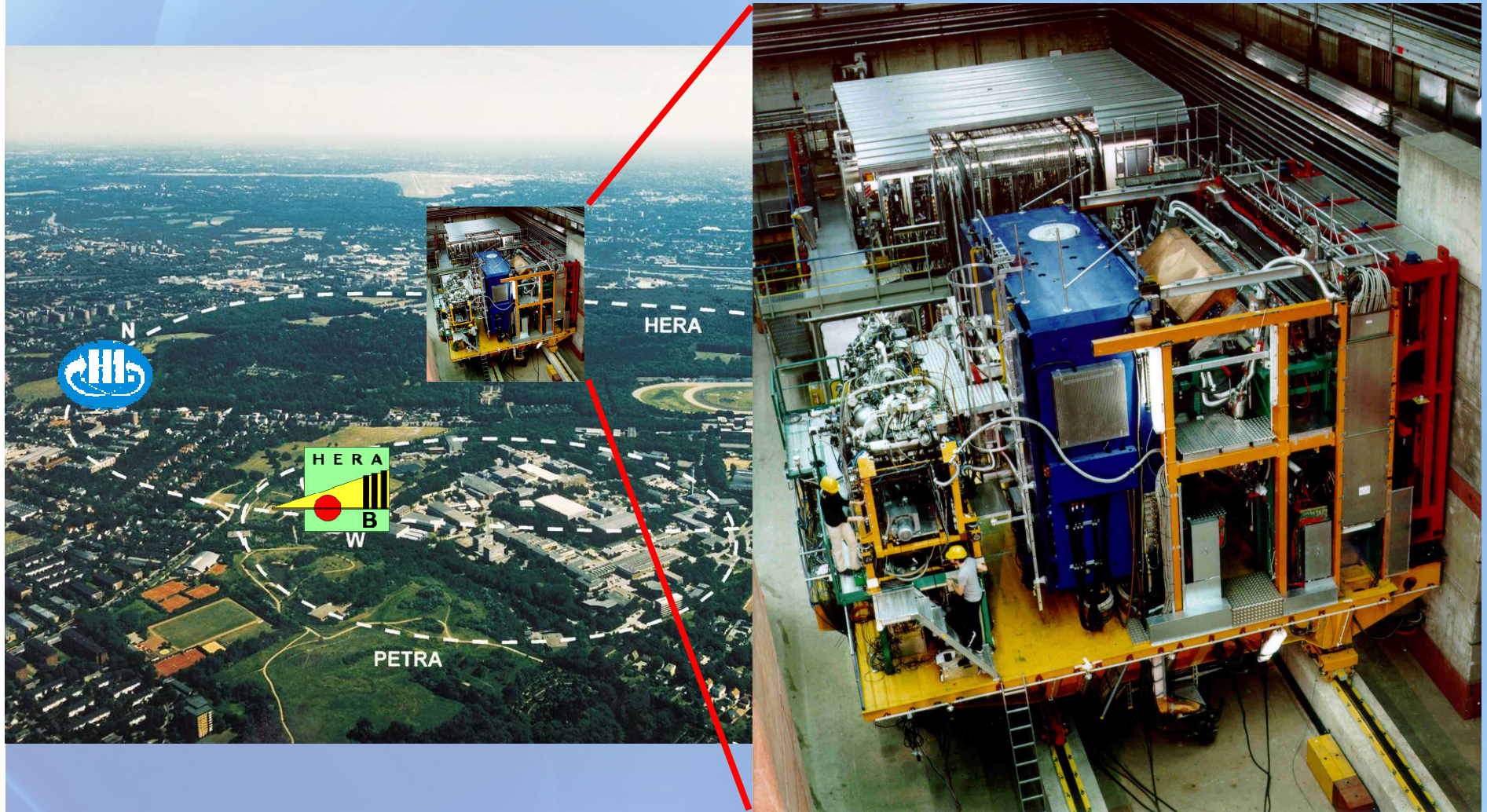
Beams: 250 GeV pp; <60>% polarization
Lumi: $1.2 \cdot 10^{31} \text{cm}^{-2}\text{s}^{-1}$

PHENIX ion experiments



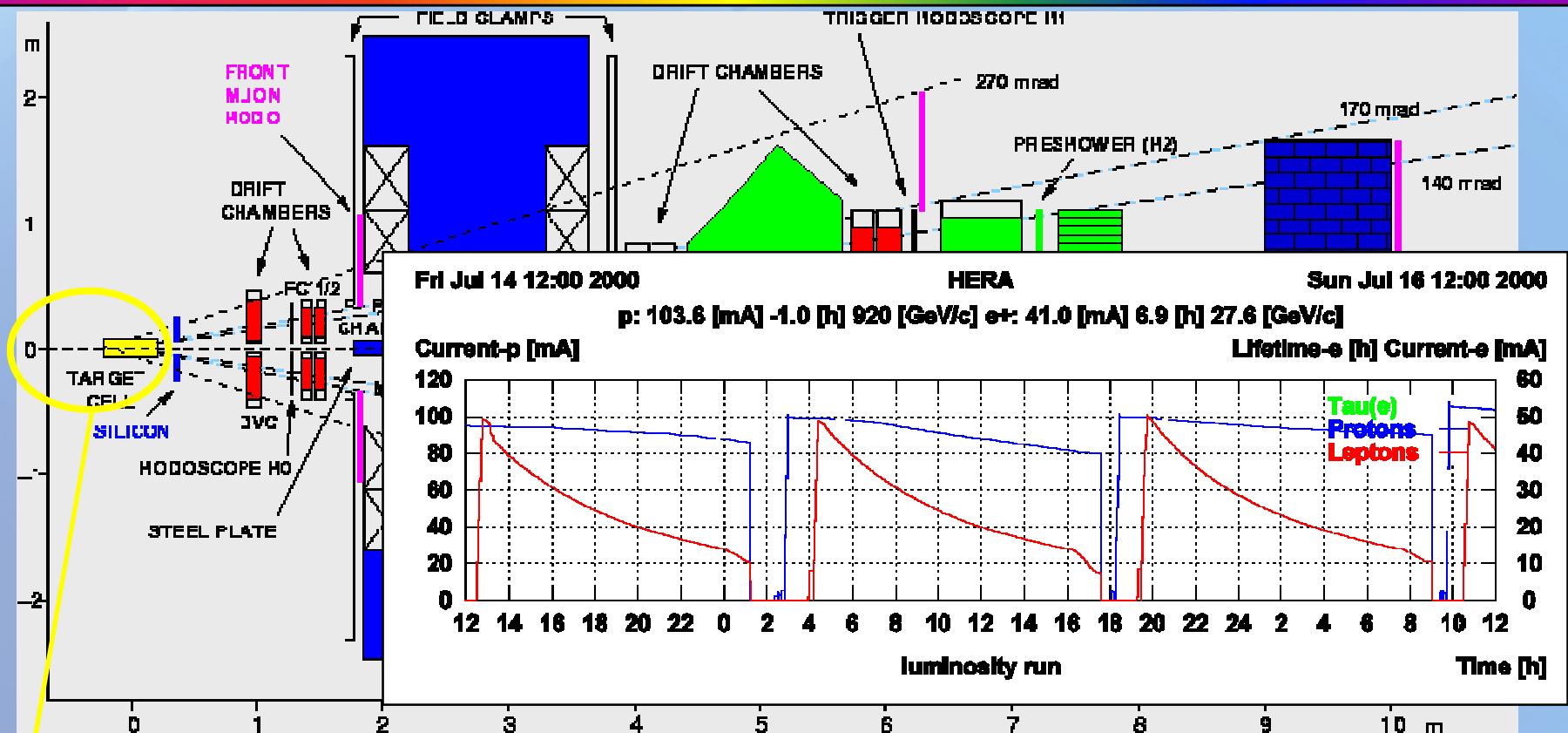
E.C. Aschenauer

SFU, VanCouver, March 2007



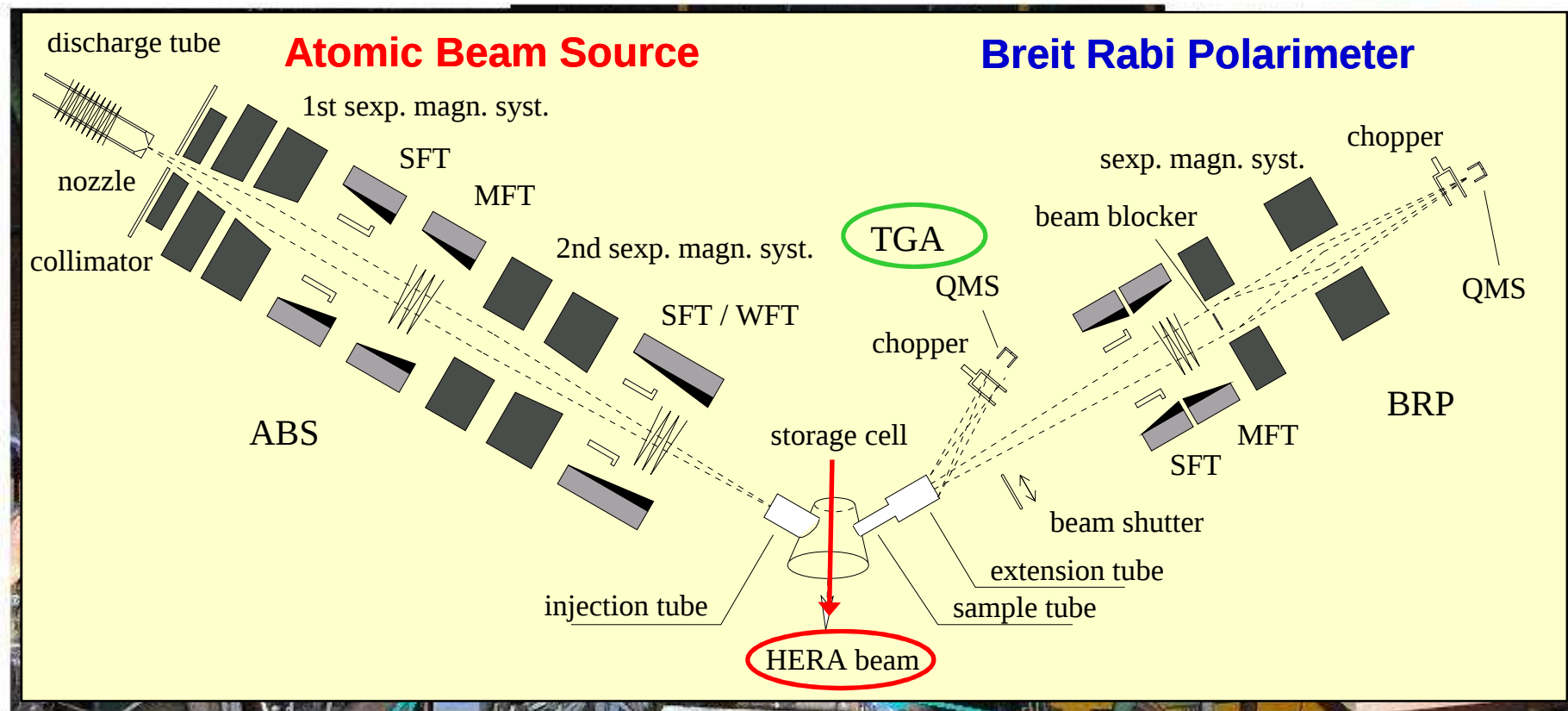
HERA: e^+/e^- (27GeV) - proton (920 GeV) collider

The HERMES Spectrometer



Internal Gas Target: $\vec{He}$, $\vec{H}$, $\vec{D}$, $H\uparrow$ unpol: $H_2, D_2, He, N_2, Ne, Kr, Xe$

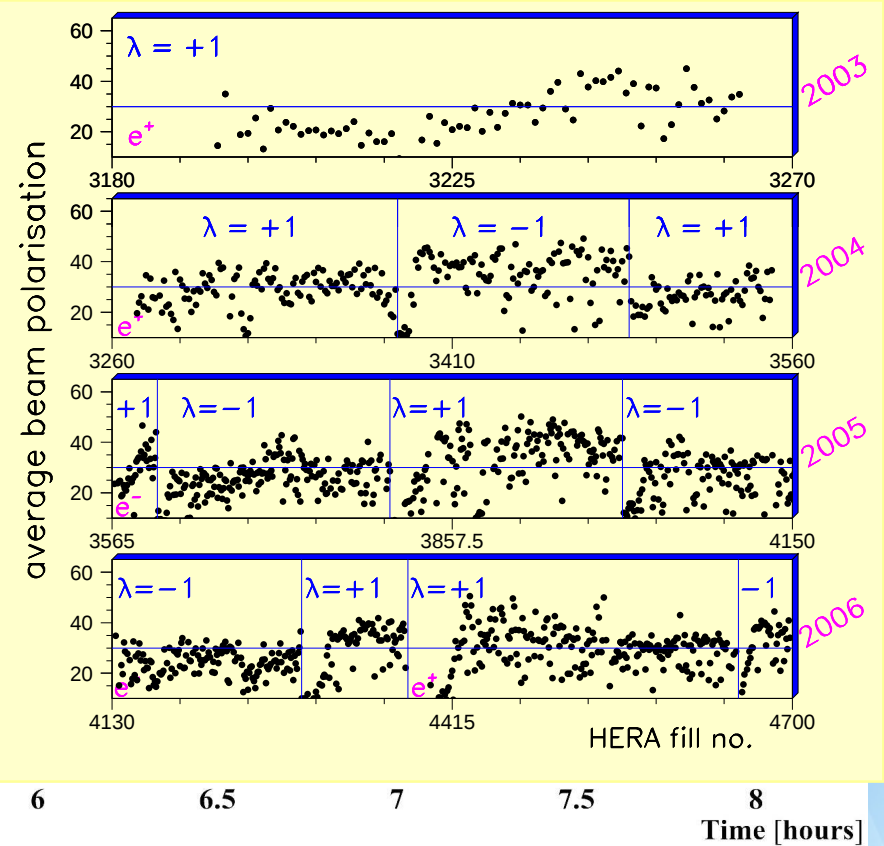
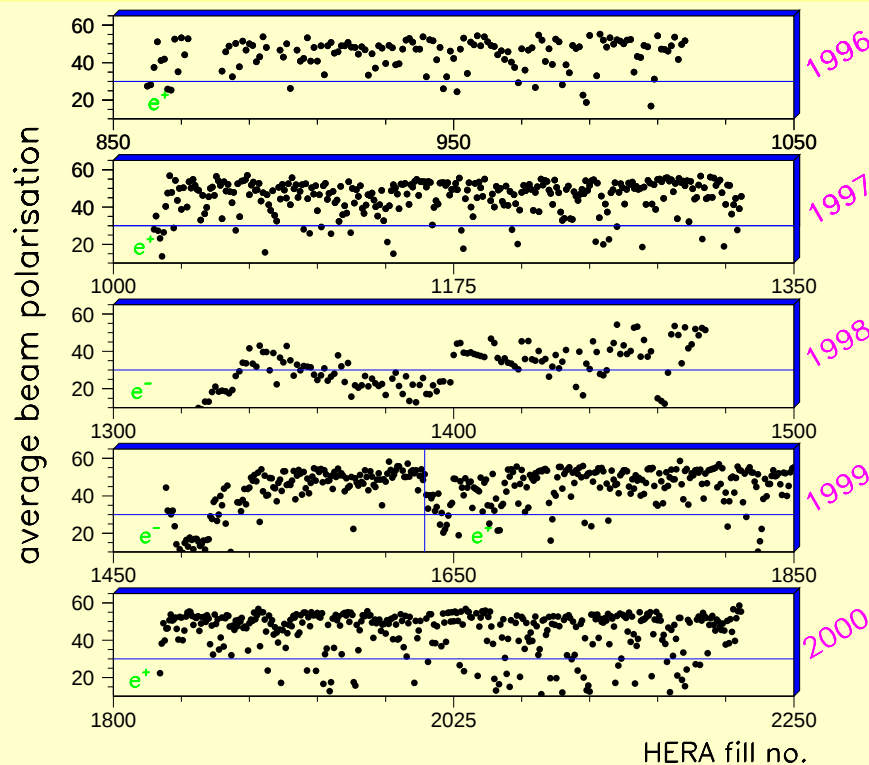
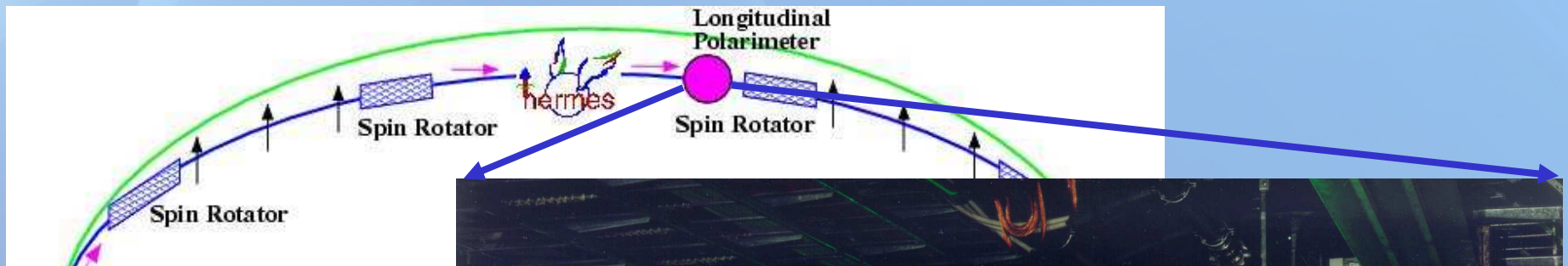
The HERMES polarized gas target



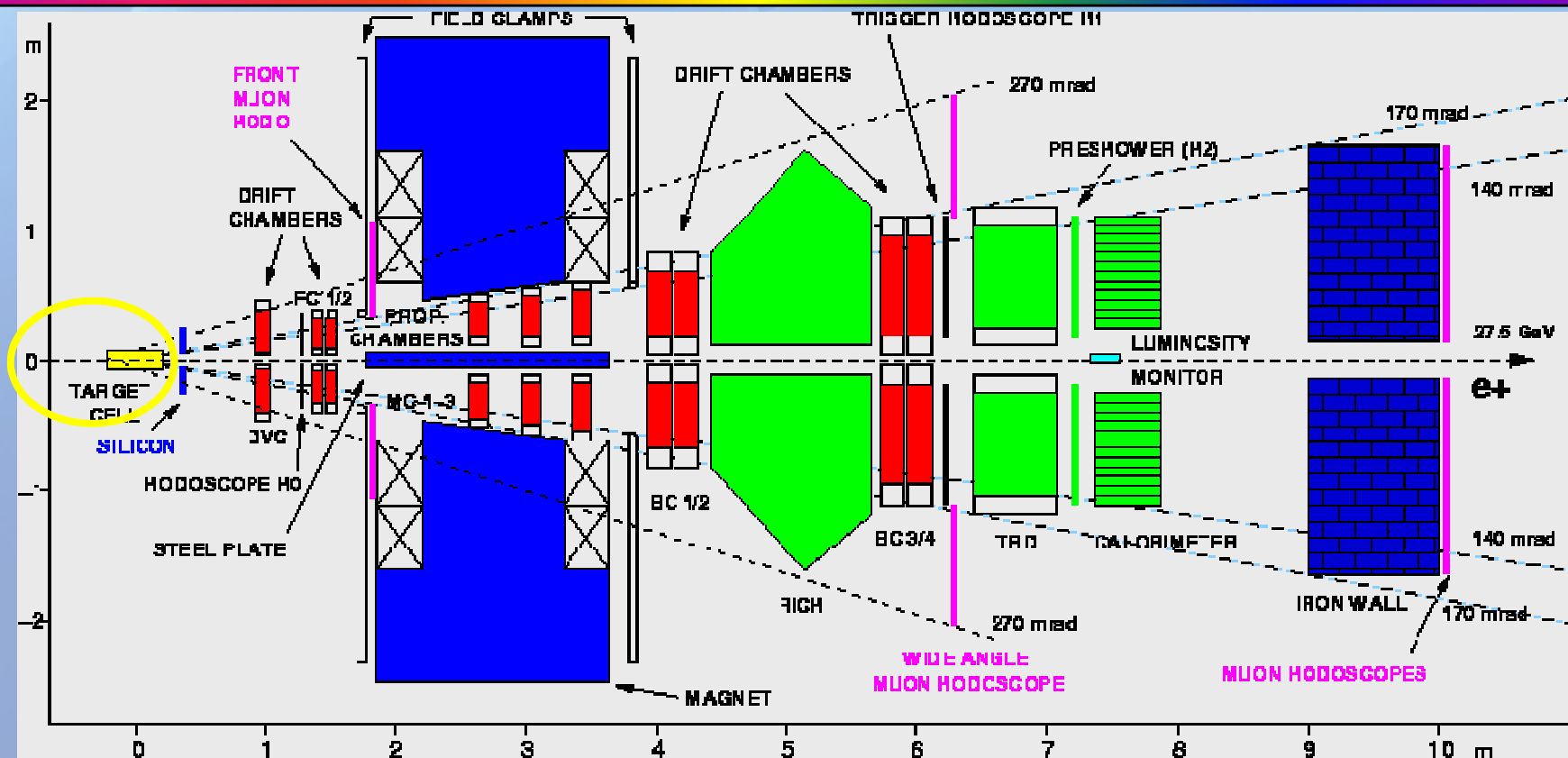
ADVANTAGES:

- no dilution; all the material is polarized
- no radiation damage
- rapid inversion of polarization direction every 90s in less than 0.5s

The HERA Polarimeters



The HERMES Spectrometer



Kinematic Range: $0.02 < x < 0.8$ at $Q^2 > 1 \text{ GeV}^2$ and $W > 2 \text{ GeV}$

Reconstruction: $\Delta p/p < 2\%$, $\Delta\theta < 1 \text{ mrad}$

Internal Gas Target: unpol: $\text{H}_2, \text{D}_2, \text{He}, \text{N}, \text{Ne}, \text{Kr}, \text{Xe}$ $\vec{\text{He}}, \vec{\text{H}}, \vec{\text{D}}, \text{H}^\uparrow$

Particle ID: TRD, Preshower, Calorimeter

→ 1997: Cherenkov, 1998 →: **RICH** + **Muon-ID**

Which Hadron (π, K, p) is Which

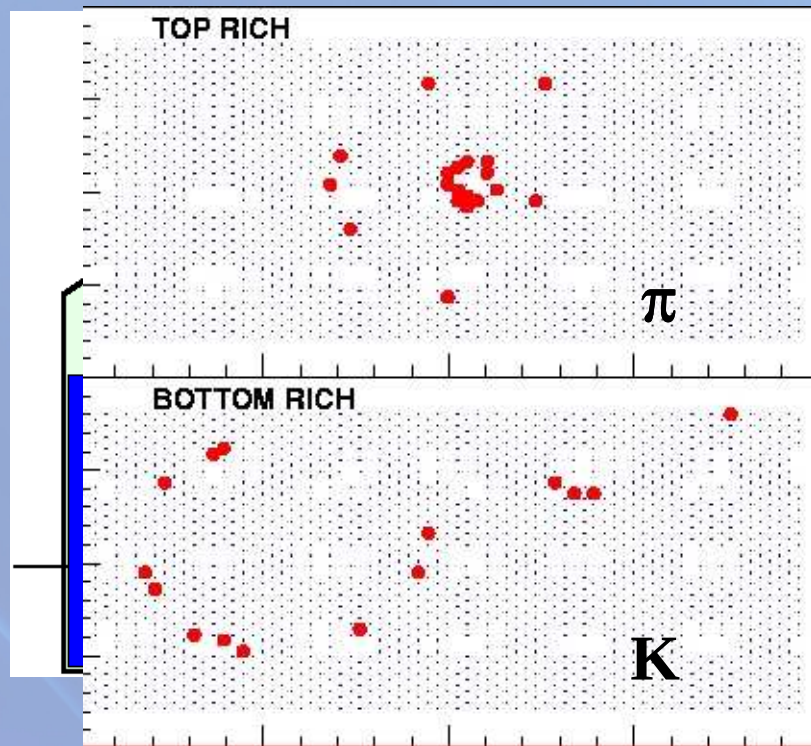
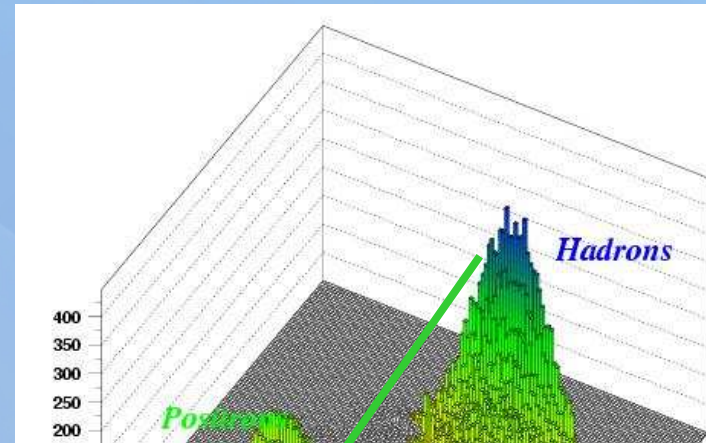
hadron/positron separation

combining signals from:

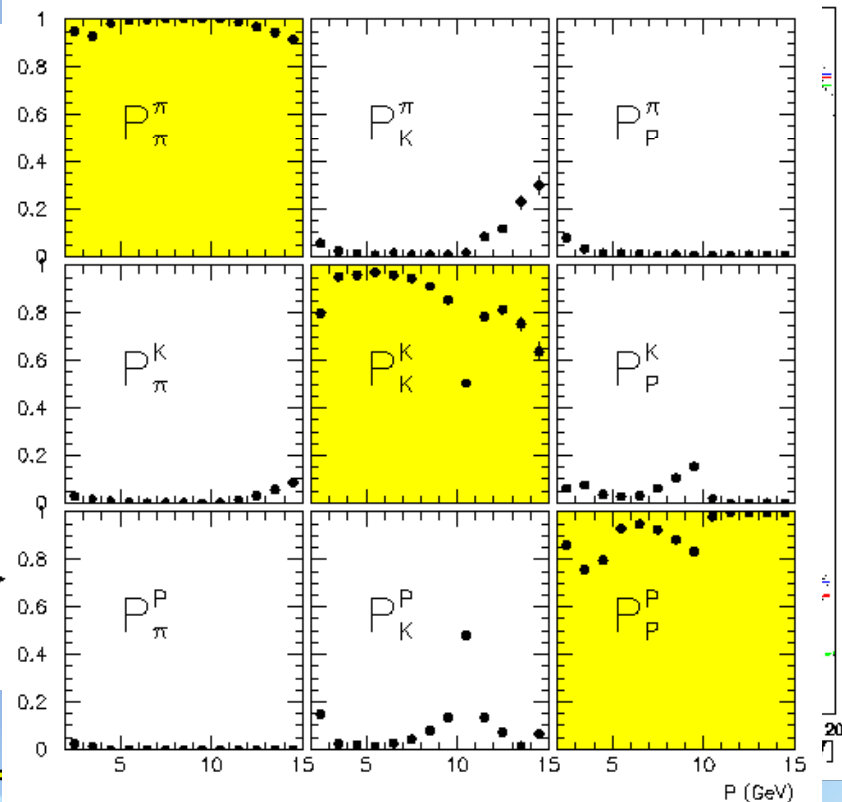
TRD, calorimeter, preshower

hadron separation

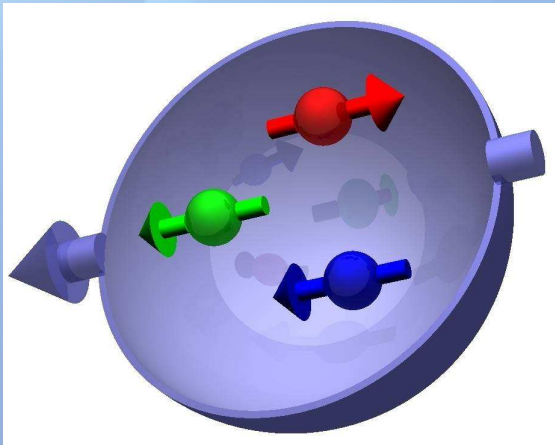
Dual radiator RICH for π, K, p



title
ck



News on the spin structure of the nucleon



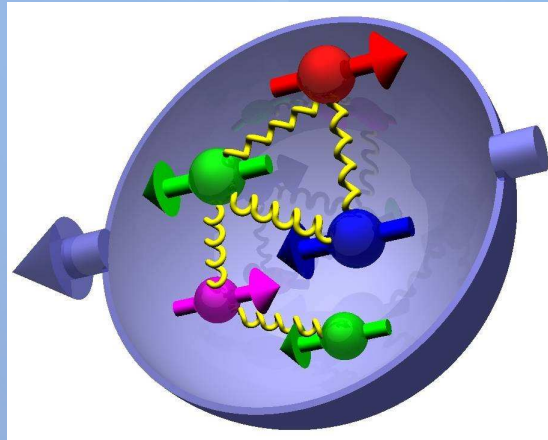
Naïve parton model

$$\Delta u_v = \frac{4}{3} \quad \Delta d_v = -\frac{1}{3}$$

BUT

1989 EMC measured
 $\Sigma = 0.120 \pm 0.094 \pm 0.138$

⇒ Spin Puzzle

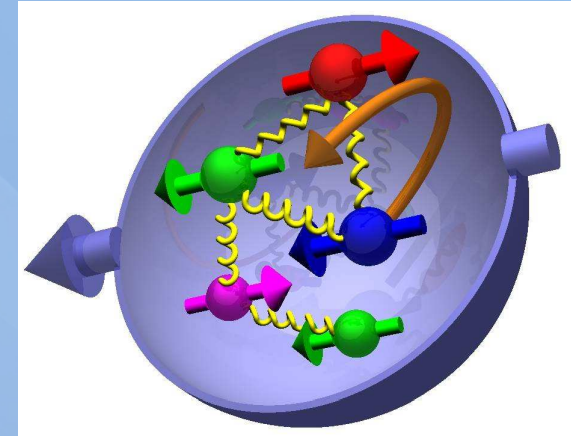


Unpolarised structure fct.

Gluons are important !

⇒ ΔG

⇒ Sea quarks Δq_s

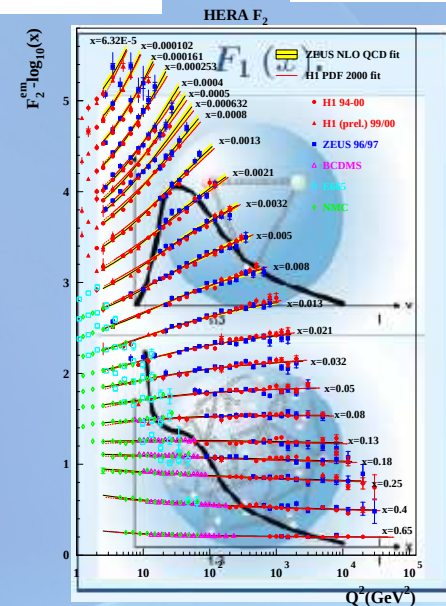
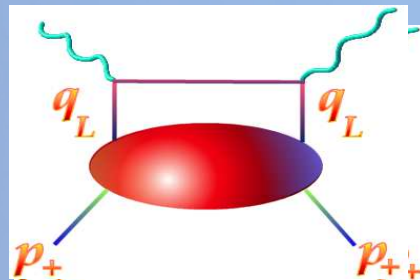
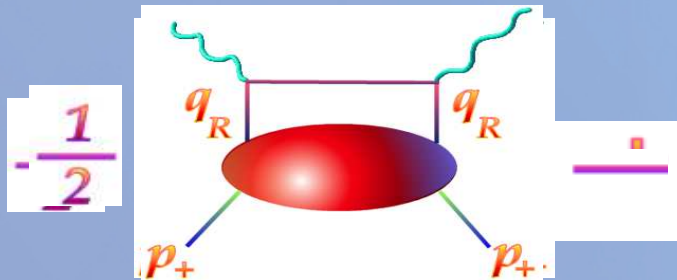
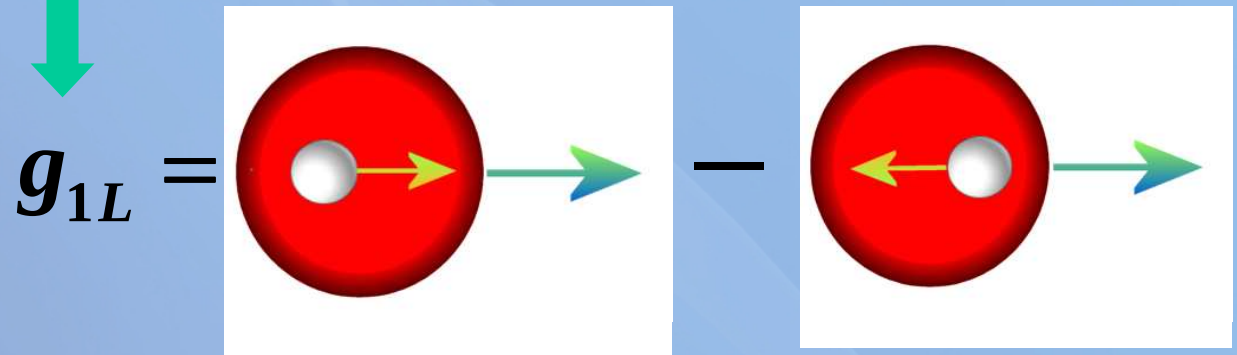
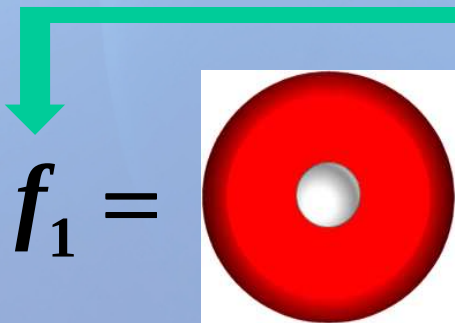


Full description of J_q and J_g
 needs
 orbital angular momentum

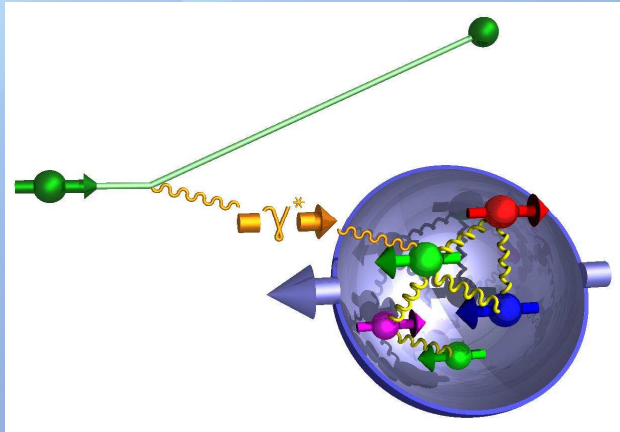
$$\frac{11}{22} = \frac{11}{22} \left(\underbrace{\left(\frac{1}{2} (\Delta u_v + \Delta d_v) + \Delta \Sigma \right)}_{(\Delta u_s + \Delta d_s + \Delta \bar{u} + \Delta \bar{d} + \Delta s + \Delta \bar{s})} + \Delta L_q + \Delta G + L_g \right)$$

The quark content of the nucleon

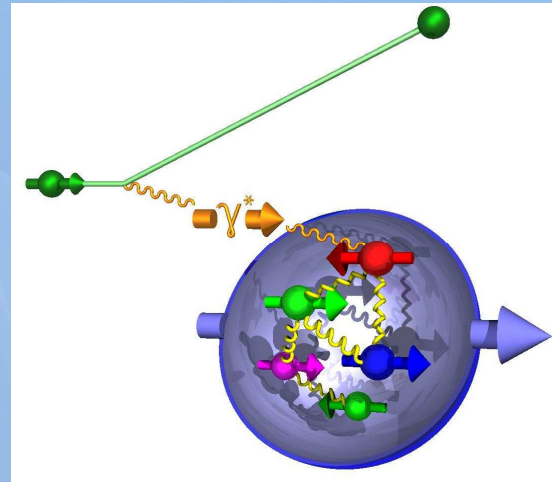
$$\Phi_{Corr}^{Tw2}(x) = \frac{1}{2} \left\{ f_1(x) + S_L g_1(x) \gamma_5 + h_1(x) \gamma_5 \gamma^1 S_T \right\} n^+$$



How to measure Quark Polarizations



$$\begin{aligned}\vec{s}_{\gamma^*} + \vec{s}_N &= 1/2 \\ \vec{s}_N &= \vec{s}_q \\ \sigma_{1/2} &\sim q^+(x)\end{aligned}$$



$$\begin{aligned}\vec{s}_{\gamma^*} + \vec{s}_N &= 3/2 \\ \vec{s}_N &= -\vec{s}_q \\ \sigma_{3/2} &\sim q^-(x)\end{aligned}$$

- Virtual photon γ^* can only couple to quarks of opposite helicity
- Select $q^+(x)$ or $q^-(x)$ by changing the orientation of target nucleon spin or helicity of incident lepton beam

$$\Delta q = q^+ - q^-$$

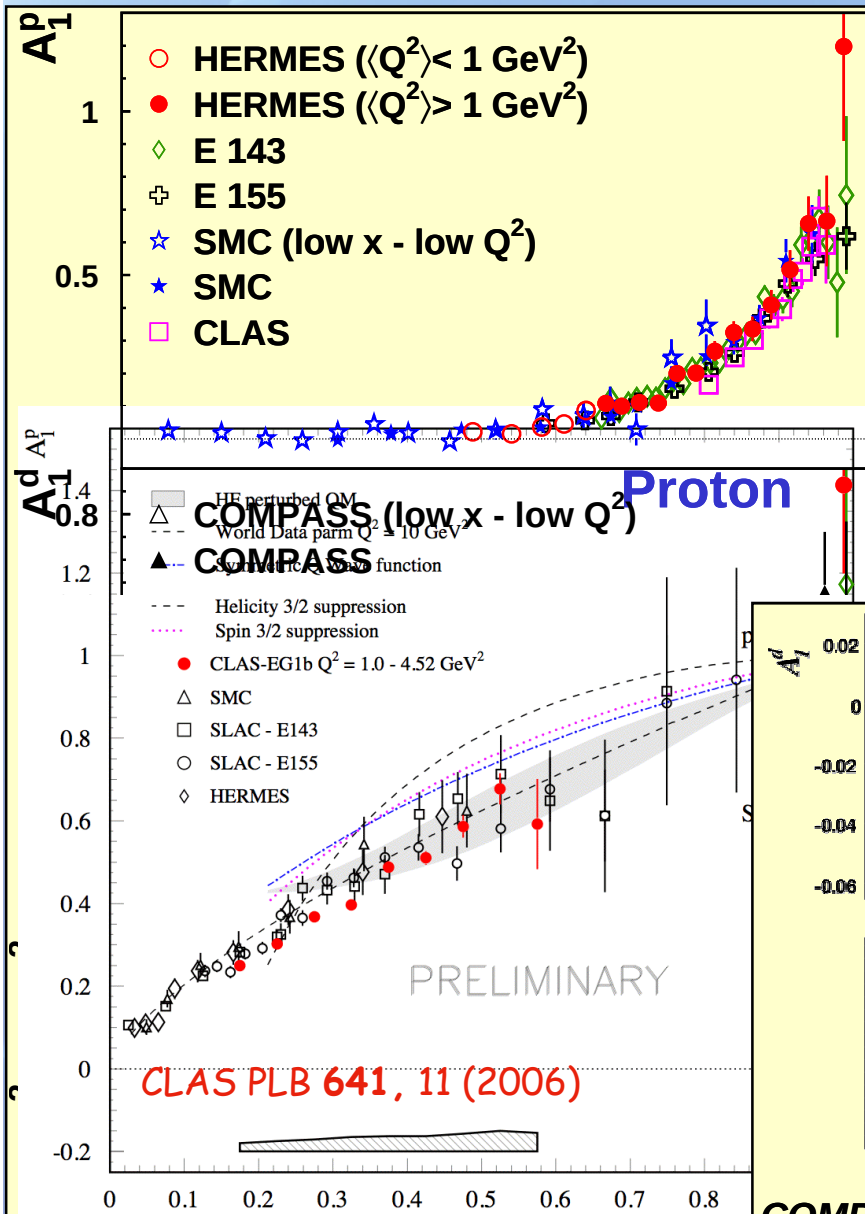
Asymmetry definition:

$$A_{\parallel}^{e',h} \sim \frac{\sigma_{1/2}^{e',h} - \sigma_{3/2}^{e',h}}{\sigma_{1/2}^{e',h} + \sigma_{3/2}^{e',h}} \quad A_{\parallel}^{e',h} = \frac{1}{\langle P_B P_T \rangle} \frac{N_{e',h}^{\Rightarrow} L^{\Rightarrow} - N_{e',h}^{\Rightarrow} L^{\Leftarrow}}{N_{e',h}^{\Rightarrow} L^{\Rightarrow} + N_{e',h}^{\Rightarrow} L^{\Leftarrow}}$$

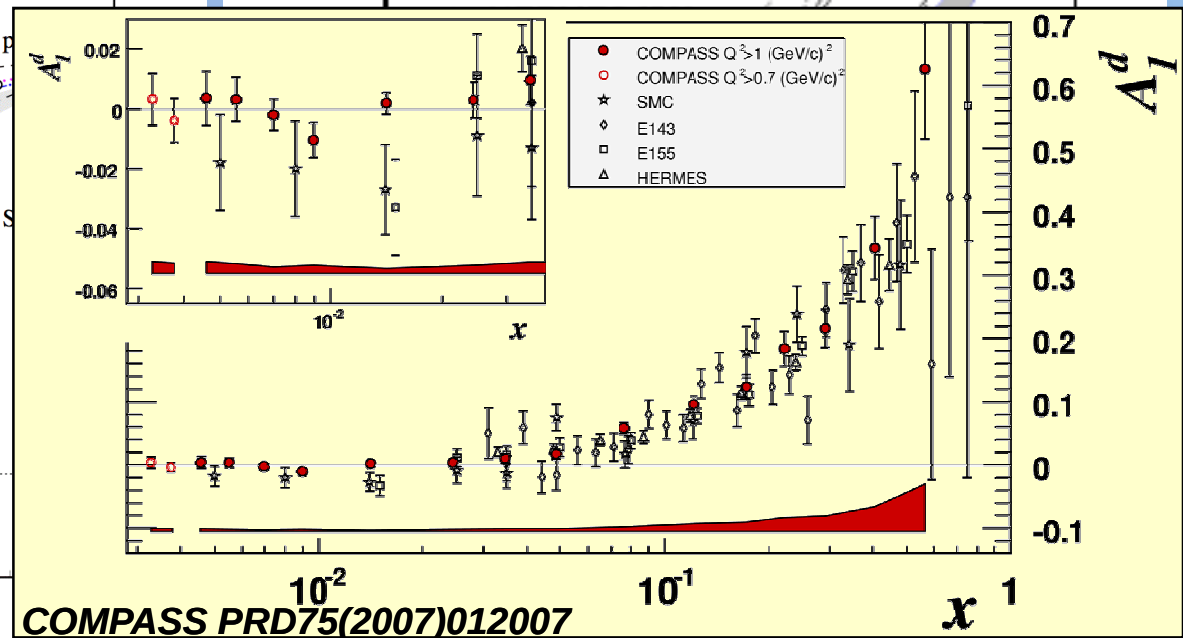
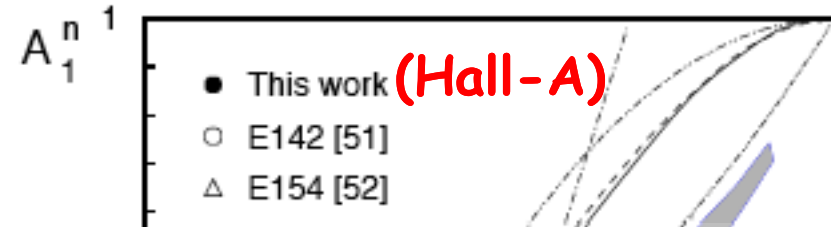
inclusive DIS: only e' info used

semi-inclusive DIS: $e'+h$ info used

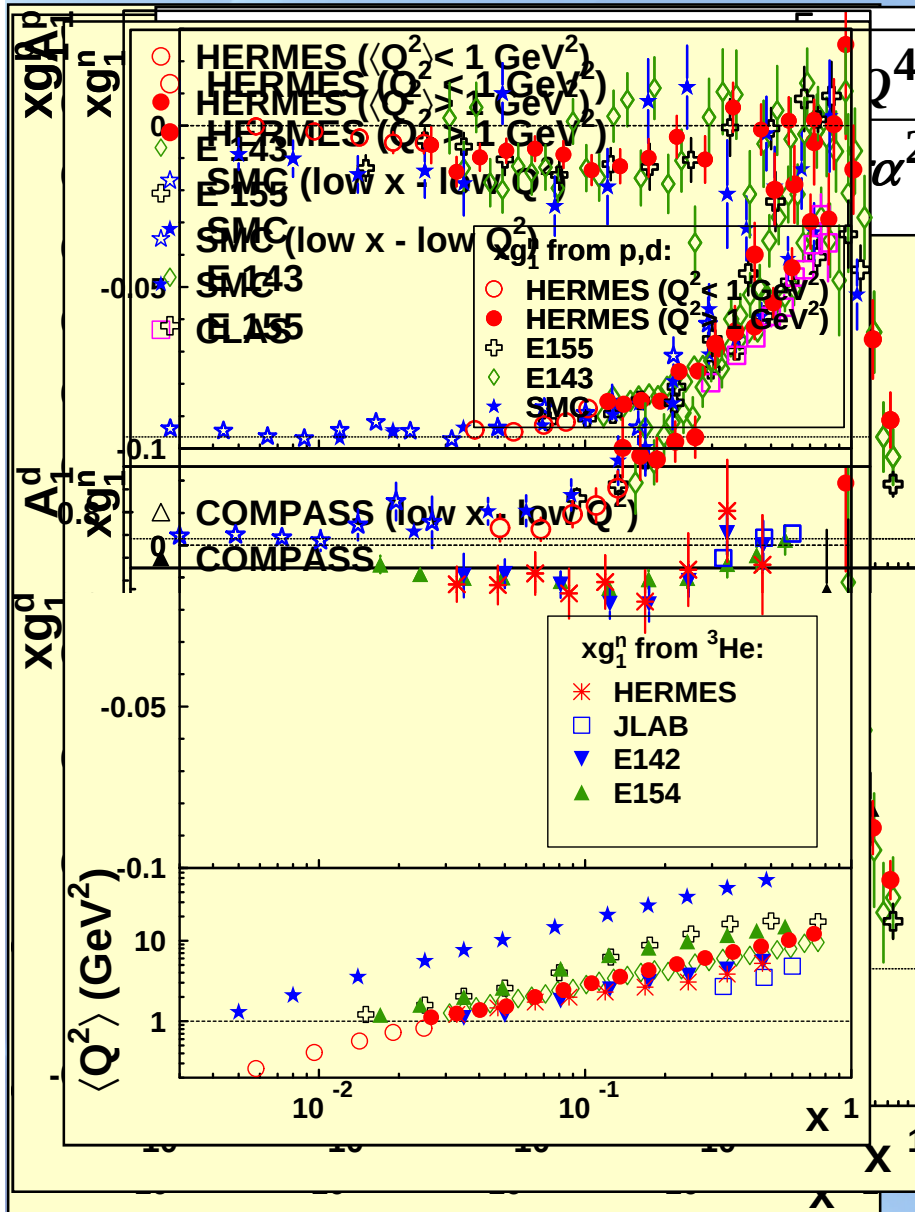
World data on inclusive DIS



- New data from COMPASS, HERMES & JLab very precise
- high x -behaviour consistent with $A \rightarrow 1$ with $x \rightarrow 1$



World data on inclusive DIS



$$\frac{d^2\sigma_{unpol}}{dx dQ^2} A_1(x, Q^2) + \frac{y}{2} \gamma^2 g_2(x, Q^2)$$

Combine p and d to get n:

$$g_1^d = \frac{1}{2} (g_1^p + g_1^n) (1 - \frac{3}{2} w_D)$$

or ^3He

$$g_1^{He^3} \approx P_n g_1^n + 2P_p g_1^p$$

► What can we learn on the PDFs

$$g_1 \sim \langle e^2 \rangle [\Delta C_\Sigma \otimes \Delta \Sigma + \Delta C_G \otimes \Delta G + \Delta C_{NS} \otimes \Delta q_{NS}^{p,n}]$$

$$\stackrel{\text{LO-QCD}}{=} \frac{1}{2} \langle e^2 \rangle [\Delta \Sigma + \Delta q_{NS}^{p,n}]$$

Compass: hep-ex/0609038

Hermes: hep-ex/0609039



HERMES: Integrals

$$\int_{0.021}^{0.9} dx g_1^d = 0.0436 \pm 0.0012(\text{stat}) \pm 0.0018(\text{syst}) \pm 0.0008(\text{par}) \pm 0.0026(\text{evol})$$

Saturation in deuteron integral is assumed

→ use only deuterium

$$a_0 = \frac{1}{\Delta C_s} \left[\frac{9\Gamma_1^d}{(1 - \frac{3}{2}w_D)} - \frac{1}{4} a_8 \Delta C_{NS} \right]$$

From hyperon beta decay

$$a_8 = 0.586 \pm 0.031$$

From neutron beta decay
 $a_3 = 1.269 \pm 0.003$

$$\Delta u + \Delta \bar{u} = \frac{1}{6} [2a_0 + a_8 + 3a_3]$$

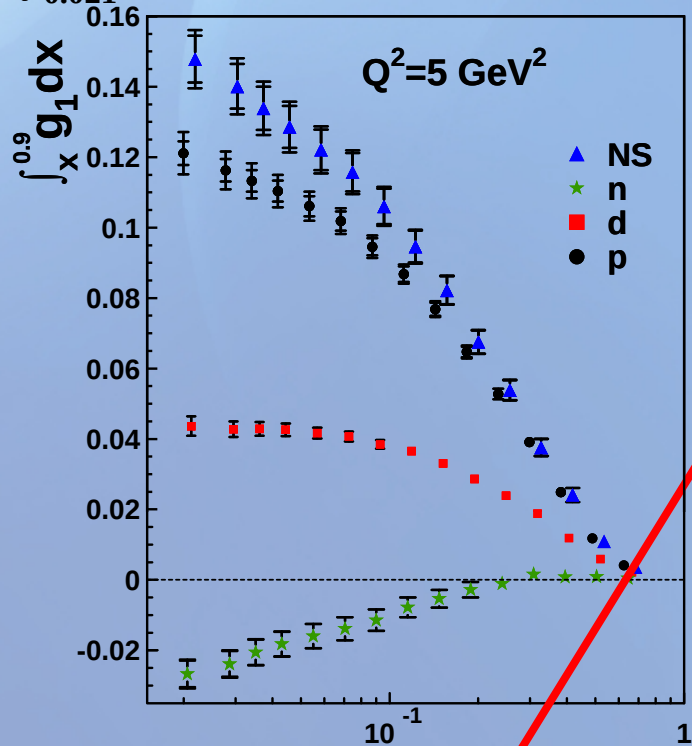
$$\Delta d + \Delta \bar{d} = \frac{1}{6} [2a_0 + a_8 - 3a_3]$$

$$\Delta s + \Delta \bar{s} = \frac{1}{3} [a_0 - a_8]$$

→ COMPASS:

$$a_0 = 0.33 \pm 0.03 \pm 0.05$$

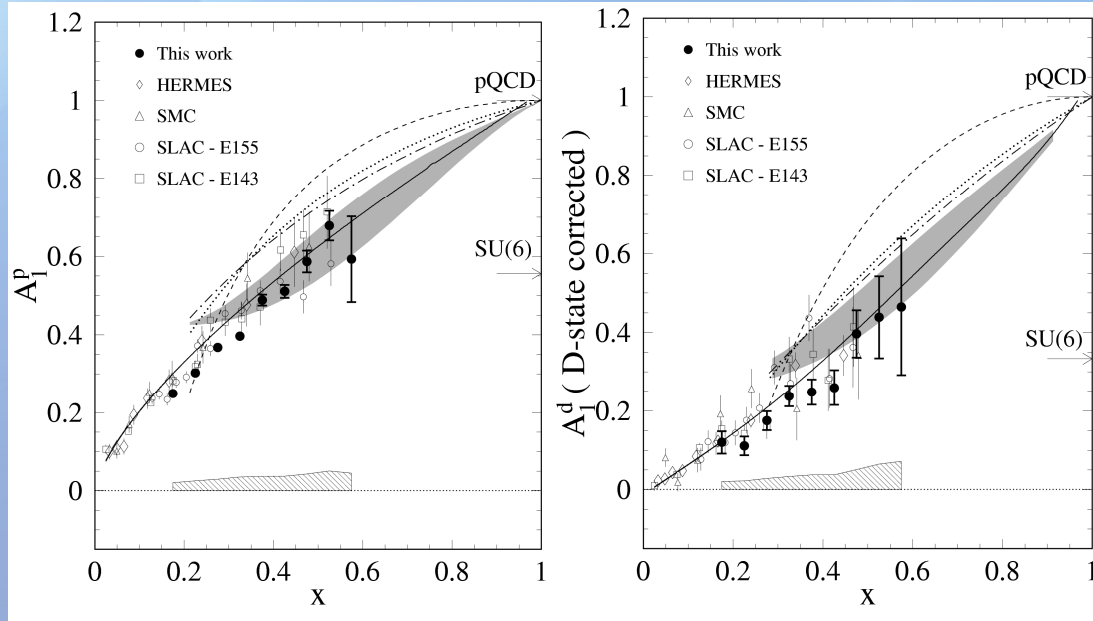
$$\Delta s + \Delta \bar{s} = -0.08 \pm 0.01 \pm 0.02$$



	central value	uncertainties		
		theor.	exp.	evol.
a_0	0.330	0.011	0.025	0.028
$\Delta u + \Delta \bar{u}$	0.842	0.004	0.008	0.009
$\Delta d + \Delta \bar{d}$	-0.427	0.004	0.008	0.009
$\Delta s + \Delta \bar{s}$	-0.085	0.013	0.008	0.009

$Q^2 = 5 \text{ GeV}^2$, NNLO in $\overline{\text{MS}}$ scheme

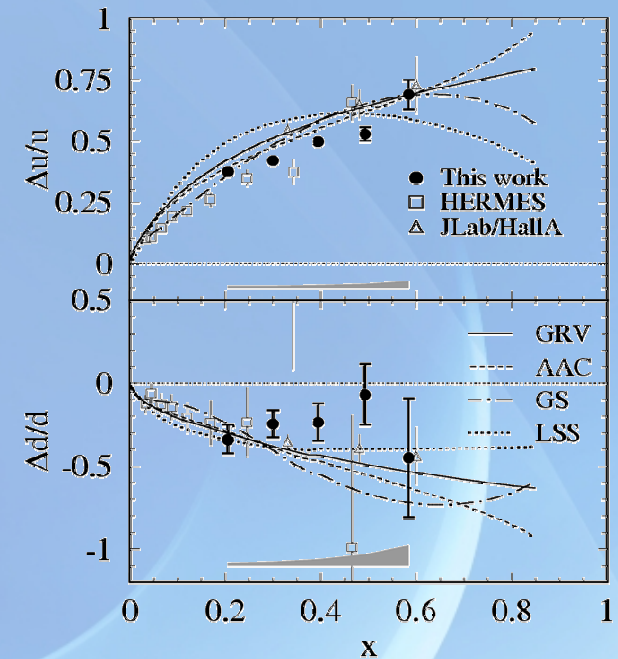




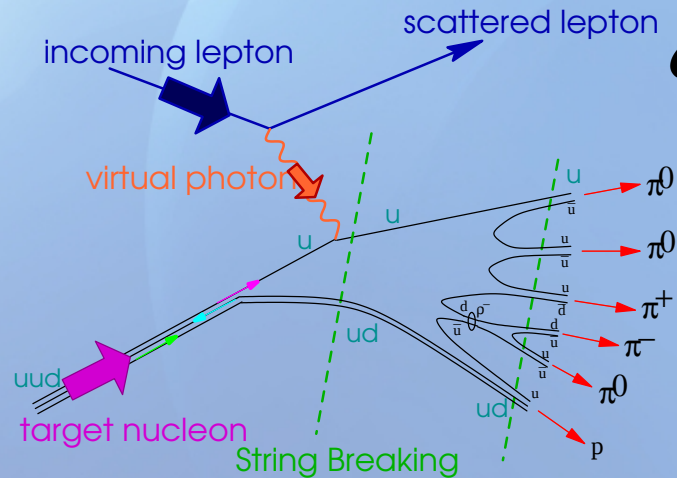
● Using:

$$\frac{\Delta u}{u} \approx \frac{5g_1^p - 2g_1^d / (1 - 1.5w_D)}{5F_1^p - 2F_1^d}$$

$$\frac{\Delta d}{d} \approx \frac{8g_1^d / (1 - 1.5w_D) - 5g_1^p}{8F_1^d - 5F_1^p}$$



Polarised quark distributions



Correlation between detected hadron and struck q_f
 $\Rightarrow$ "Flavor - Separation"

Inclusive DIS:

$$\Delta\Sigma = (\Delta u + \Delta d + \Delta\bar{u} + \Delta\bar{d} + \Delta s + \Delta\bar{s})$$

Semi-inclusive DIS:

$$(\Delta u, \Delta d, \Delta\bar{u}, \Delta\bar{d}, \Delta s, \Delta\bar{s})$$

In LO-QCD:

$$A_1^h(x, Q^2, z) = \frac{\sigma_{1/2}^h - \sigma_{3/2}^h}{\sigma_{1/2}^h + \sigma_{3/2}^h} \sim C \sum_q \frac{e_q^2 q(x, Q^2) D_q^h(z, Q^2)}{\sum_{q'} e_{q'}^2 q'(x, Q^2) D_{q'}^h(z, Q^2)} \frac{\Delta q(x, Q^2)}{q(x, Q^2)}$$

COMPASS: Valence PDFs

$$A^{+/-} = \frac{(\sigma_{h^+}^{\Rightarrow} - \sigma_{h^-}^{\Rightarrow}) - (\sigma_{h^+}^{\Leftarrow} - \sigma_{h^-}^{\Leftarrow})}{(\sigma_{h^+}^{\Rightarrow} - \sigma_{h^-}^{\Rightarrow}) + (\sigma_{h^+}^{\Leftarrow} - \sigma_{h^-}^{\Leftarrow})}$$

● **For LO:**

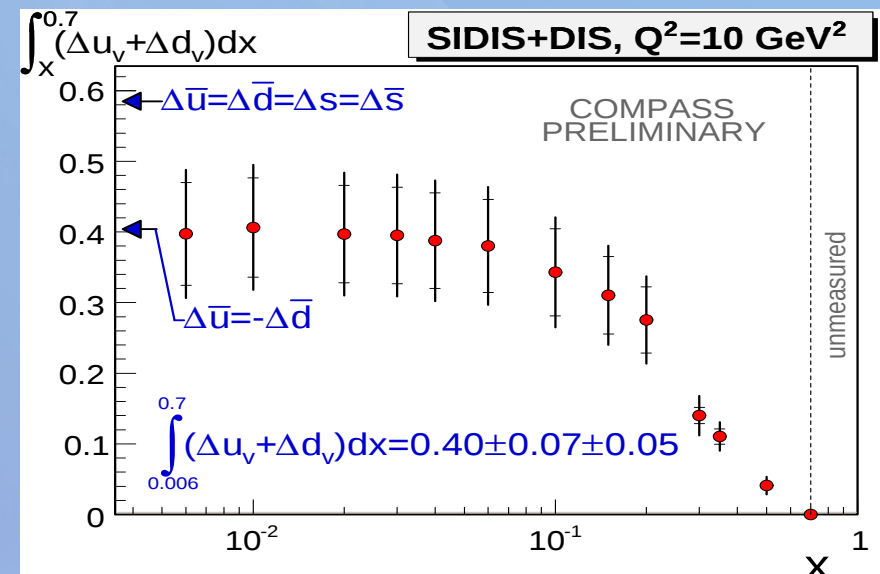
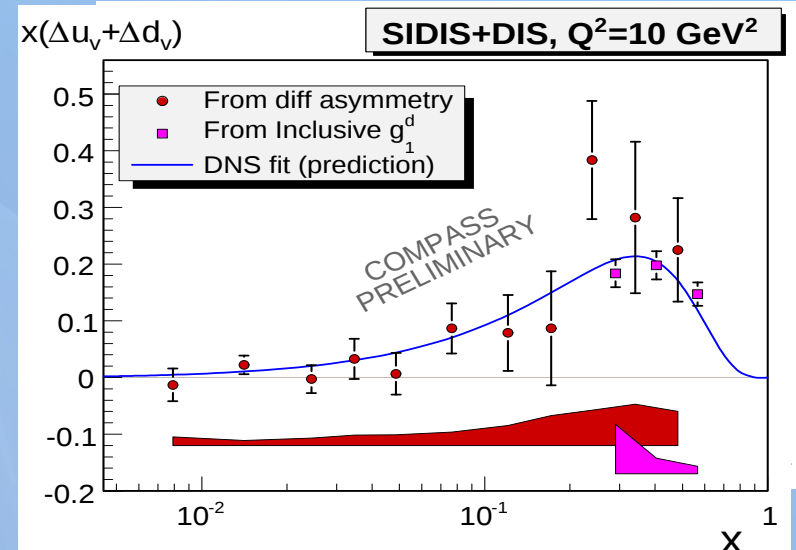
$$A_d^{\pi^+-\pi^-}(x) = A_d^{K^+-K^-} = \frac{\Delta u_v(x) + \Delta d_v(x)}{u_v(x) + d_v(x)}$$

● **Assuming:**

$$\Gamma_1^N = \frac{1}{9}(a_0 + \frac{1}{4}a_8)$$

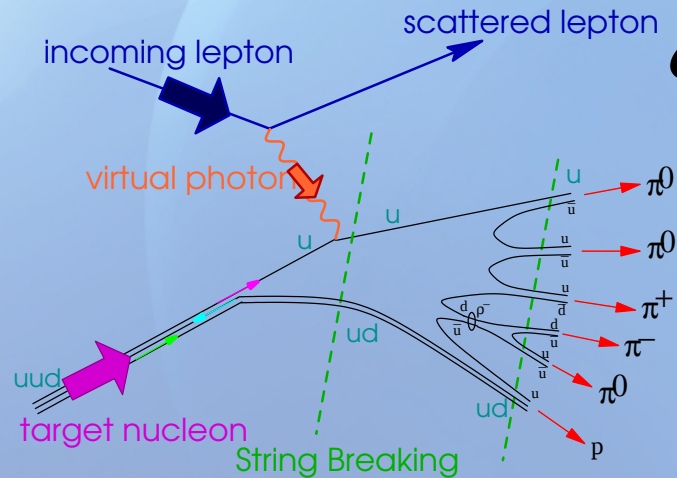
$$\Gamma_v = \int_0^1 (\Delta u_v(x) + \Delta d_v(x)) dx$$

$$\begin{aligned} \Delta \bar{u} + \Delta \bar{d} &= 3\Gamma_1^N - \frac{1}{2}\Gamma_v + \frac{1}{12}a_8 \\ &= (\Delta s + \Delta \bar{s}) + \frac{1}{2}(a_8 - \Gamma_v) \end{aligned}$$



➡ Γ_v is $2.5\sigma_{\text{stat}}$ away from flavor symmetric sea scenario

Polarised quark distributions



Correlation between detected hadron and struck q_f
 ⇒ "Flavor - Separation"

Inclusive DIS:

$$\Delta\Sigma = (\Delta u + \Delta d + \Delta\bar{u} + \Delta\bar{d} + \Delta s + \Delta\bar{s})$$

Semi-inclusive DIS:

$$(\Delta u, \Delta d, \Delta\bar{u}, \Delta\bar{d}, \Delta s, \Delta\bar{s})$$

In LO-QCD:

$$A_1^h(x, Q^2, z) = \frac{\sigma_{1/2}^h - \sigma_{3/2}^h}{\sigma_{1/2}^h + \sigma_{3/2}^h} \sim C \sum_q \frac{e_q^2 q(x, Q^2) D_q^h(z, Q^2)}{\underbrace{\sum_{q'} e_{q'}^2 q'(x, Q^2) D_{q'}^h(z, Q^2)}_{P_q^h(x, Q^2, z)}} \frac{\Delta q(x, Q^2)}{q(x, Q^2)}$$

$P_q^h(x, Q^2, z) \leftarrow MC$

Extract Δq by solving:

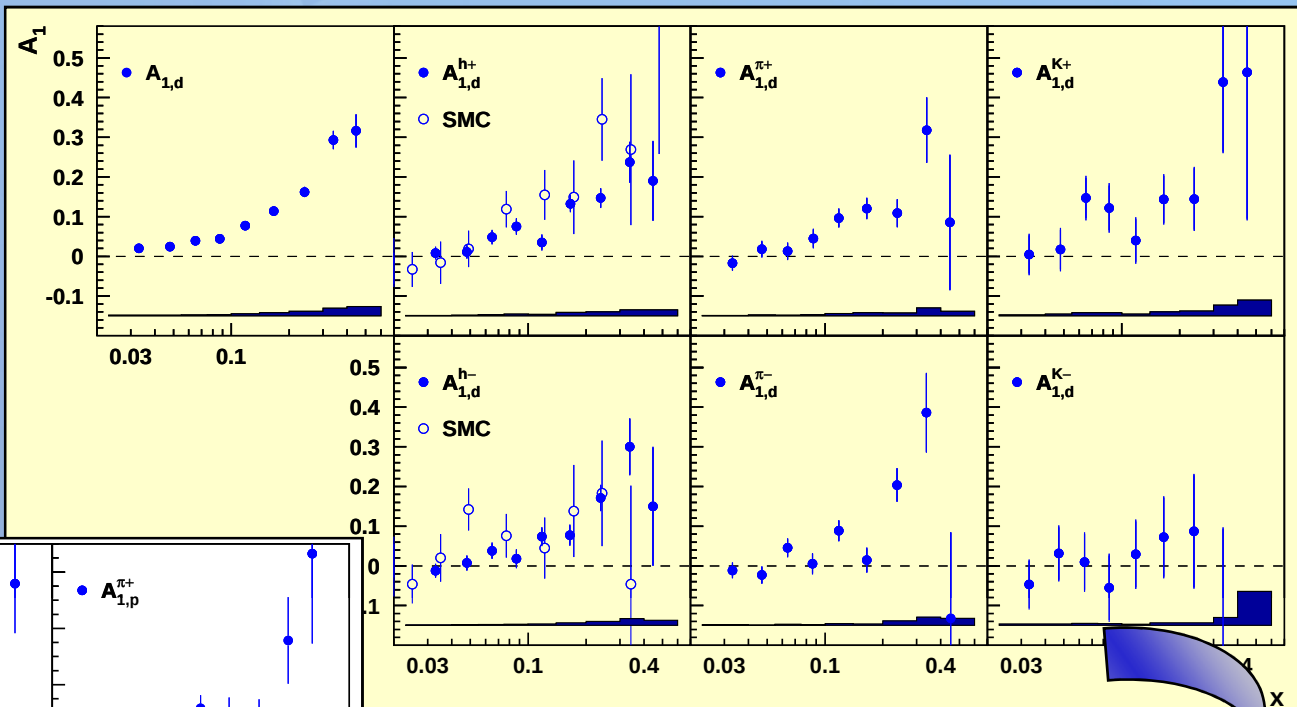
$$\vec{A} = P \vec{Q}$$

$$\vec{A} = (A_{1,p}(x), A_{1,p}^{\pi^\pm}(x), A_{1,d}(x), A_{1,d}^{\pi^\pm}(x), A_{1,d}^{K^\pm}(x))$$

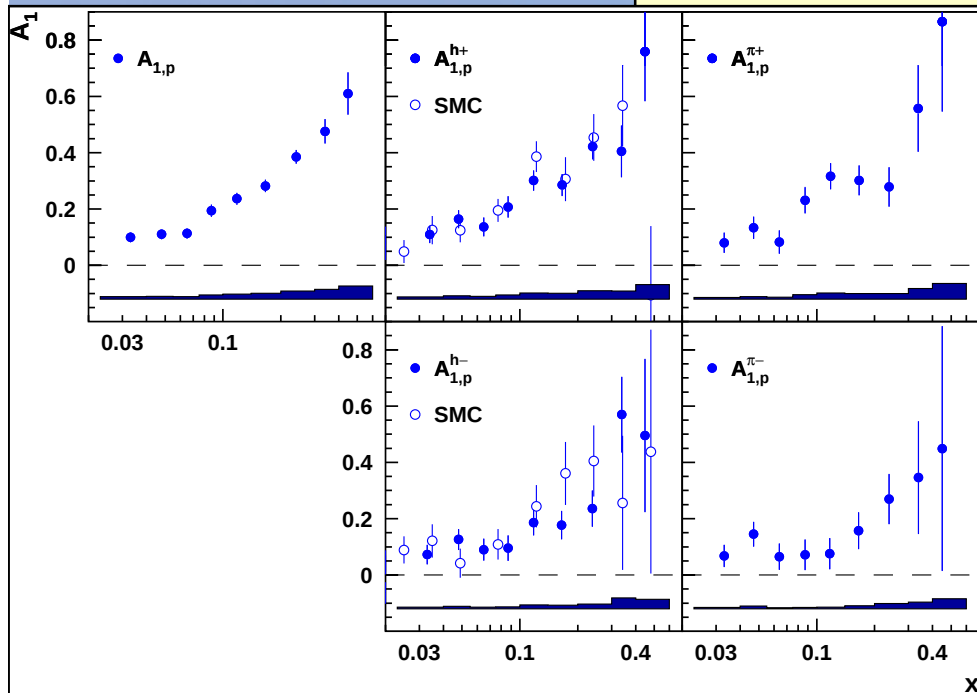
$$\vec{Q} = \left(\frac{\Delta u}{u}, \frac{\Delta d}{d}, \frac{\Delta\bar{u}}{u}, \frac{\Delta\bar{d}}{d}, \frac{\Delta s}{s}, \frac{\Delta\bar{s}}{s} \right)$$

Measured Asymmetries

Deuterium



Proton

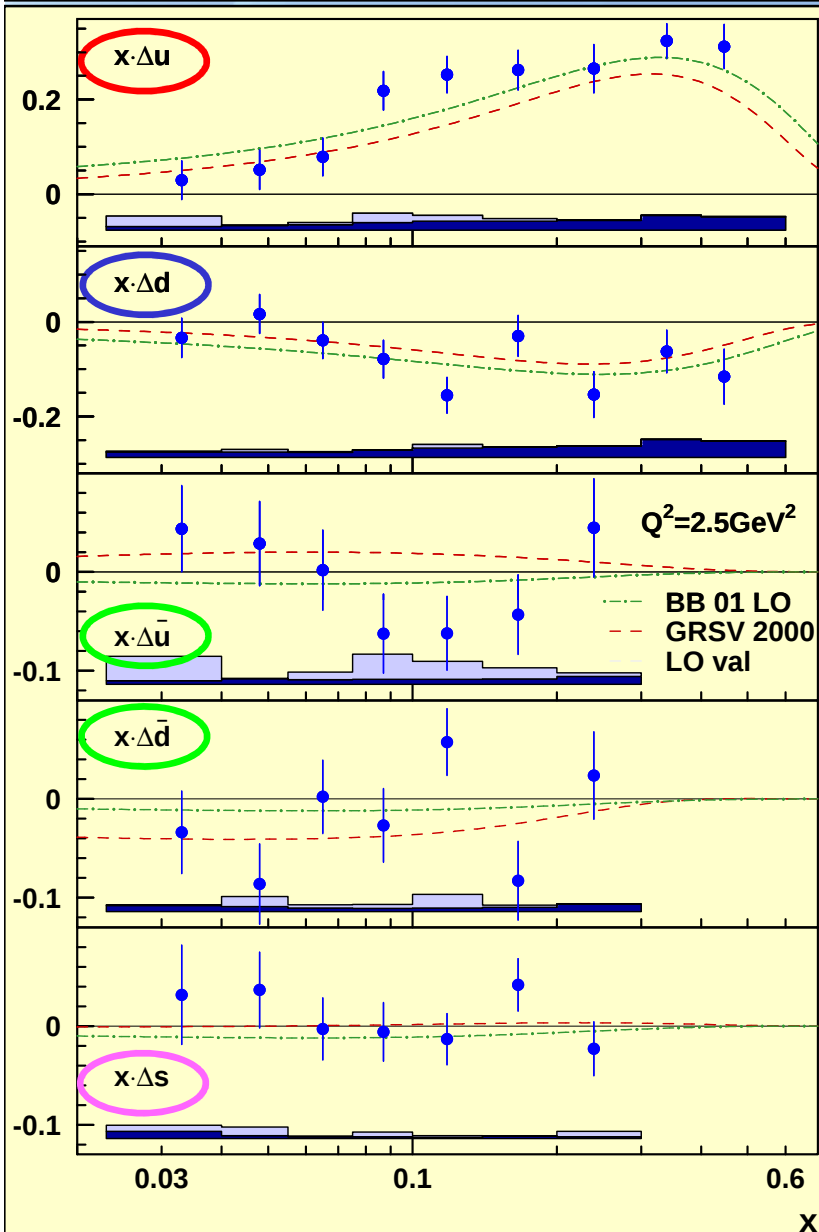


$$A_1^{K^-} \sim 0 \Rightarrow K^- = (\bar{u}s)$$

● statistics sufficient for
5 parameter fit

$$\vec{Q} = \left(\frac{\Delta u}{u}, \frac{\Delta d}{d}, \frac{\Delta \bar{u}}{\bar{u}}, \frac{\Delta \bar{d}}{\bar{d}}, \frac{\Delta s}{s}, \frac{\Delta \bar{s}}{\bar{s}} = 0 \right)$$

Polarized Quark Densities



- First complete separation of pol. PDFs without assumption on sea polarization
- $\Delta u(x) > 0$ $\Delta d(x) < 0$
- Polarised opposite to proton spin
- $\Delta \bar{u}(x), \Delta \bar{d}(x) \sim 0$
- No indication for $\Delta s(x) < 0$
- In measured range (0.023 – 0.6)

$$\int \Delta \bar{u} = -0.002 \pm 0.043$$

$$\int \Delta \bar{d} = -0.054 \pm 0.035$$

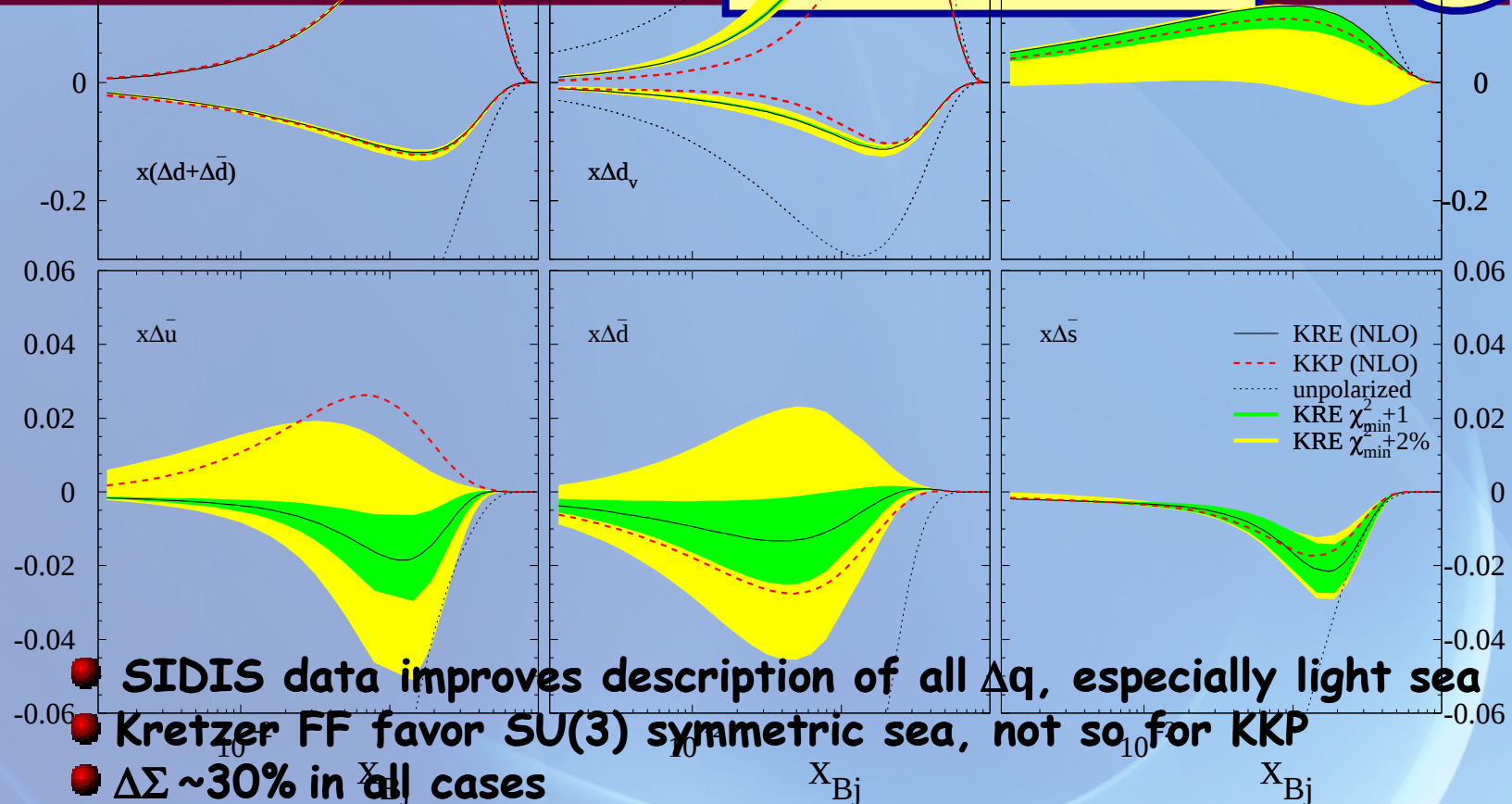
$$\int \Delta s = +0.028 \pm 0.034$$

NLO FIT to DIS & SIDIS Data

D. De Florian et al. hep-ph/0504155

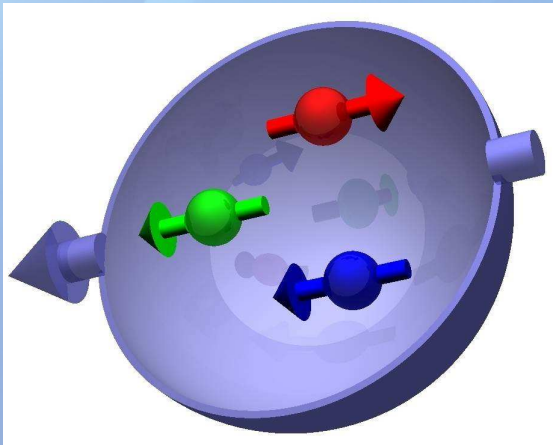
NLO @ $Q^2=10 \text{ GeV}^2$

	χ^2_{DIS}	χ^2_{SIDIS}	Δu_v	Δd_v	$\Delta \bar{u}$	$\Delta \bar{d}$	Δs	Δg
Kretzer	206	225	0.94	0.34				0.68
KKP	206	231	0.70	-0.26				0.57



- SIDIS data improves description of all Δq , especially light sea
- Kretzer FF favor SU(3) symmetric sea, not so for KKP
- $\Delta \Sigma \sim 30\%$ in all cases

The Gluon Polarization



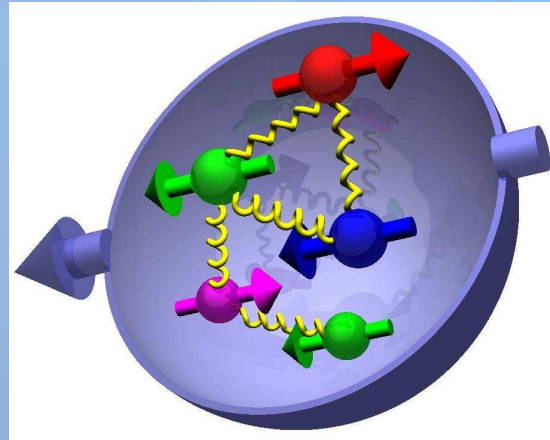
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BUT

1989 EMC measured
 $\Sigma = 0.120 \pm 0.094 \pm 0.138$

⇒ Spin Puzzle



Unpolarised structure fct.

Gluons are important !

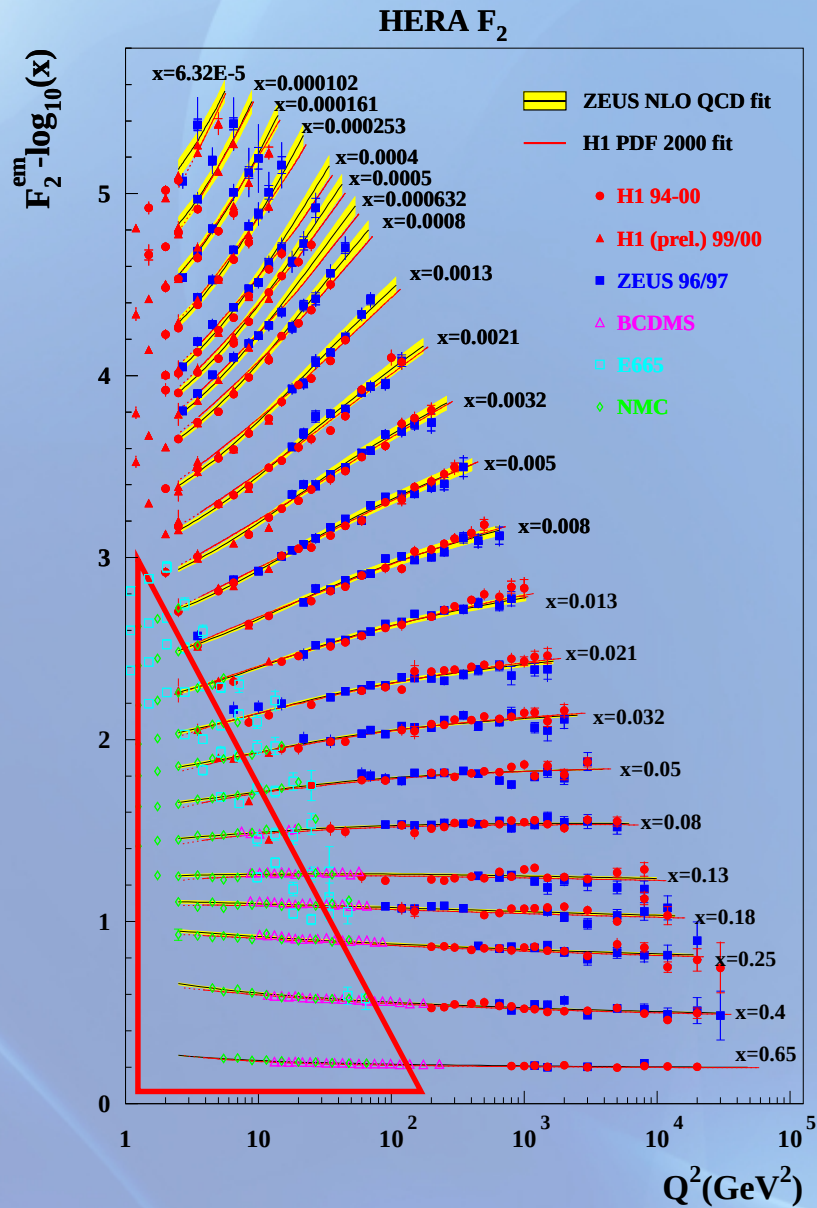
⇒ ΔG

⇒ Sea quarks Δq_s

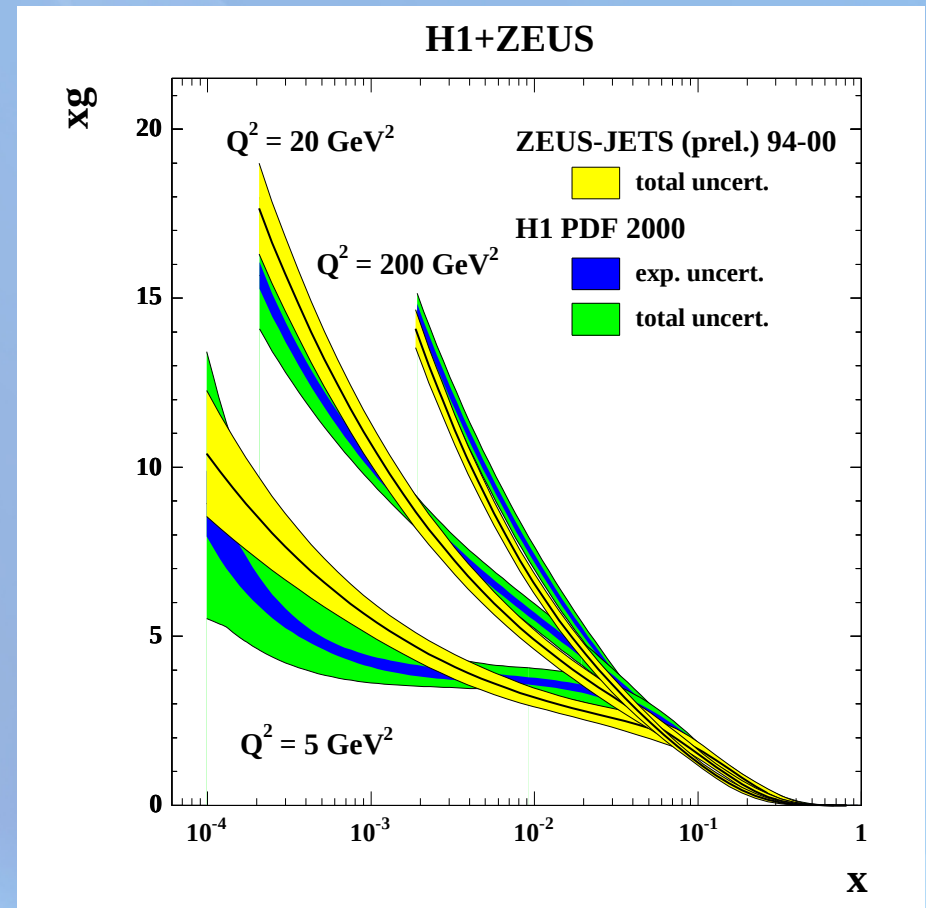
$$\frac{1}{2} = \frac{1}{2} (\Delta u_v + \Delta d_v + \Delta q_s) + \Delta G$$

$$(\Delta u_s + \Delta d_s + \Delta \bar{u} + \Delta \bar{d} + \Delta s + \Delta \bar{s})$$

Unpolarized Gluon Distribution

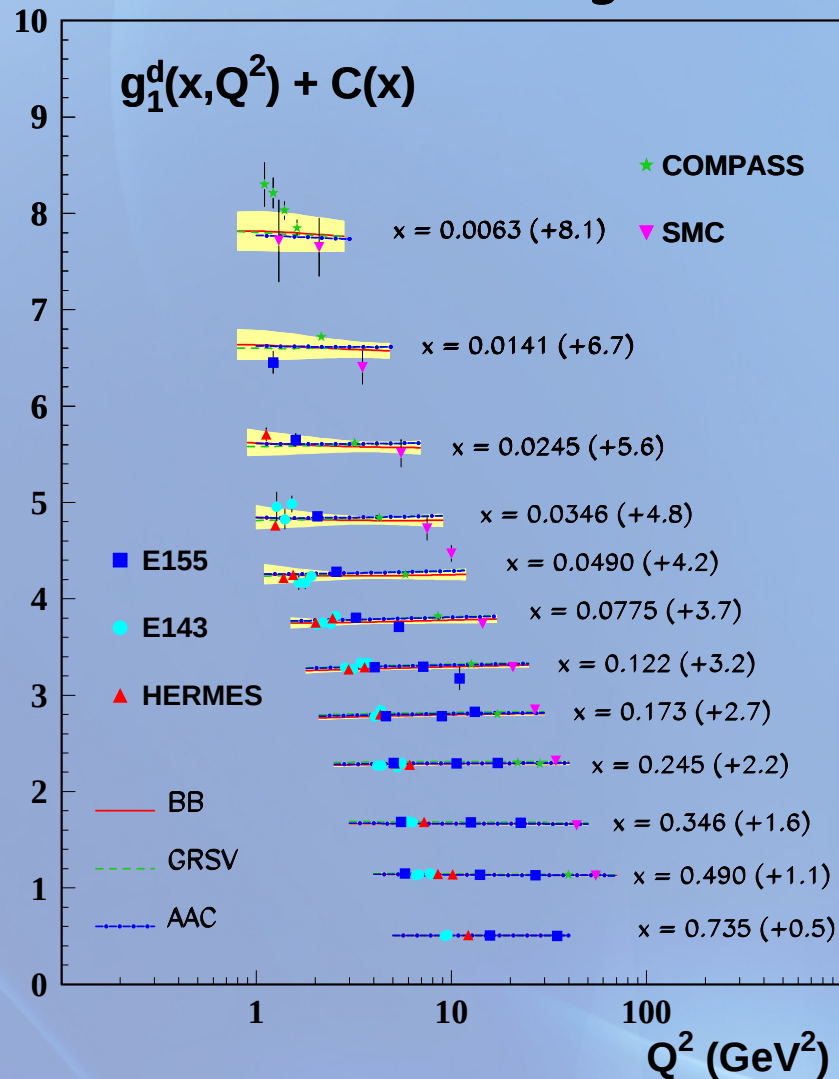


● big Q^2 - x lever arm
 ➡ very accurate $G(x)$

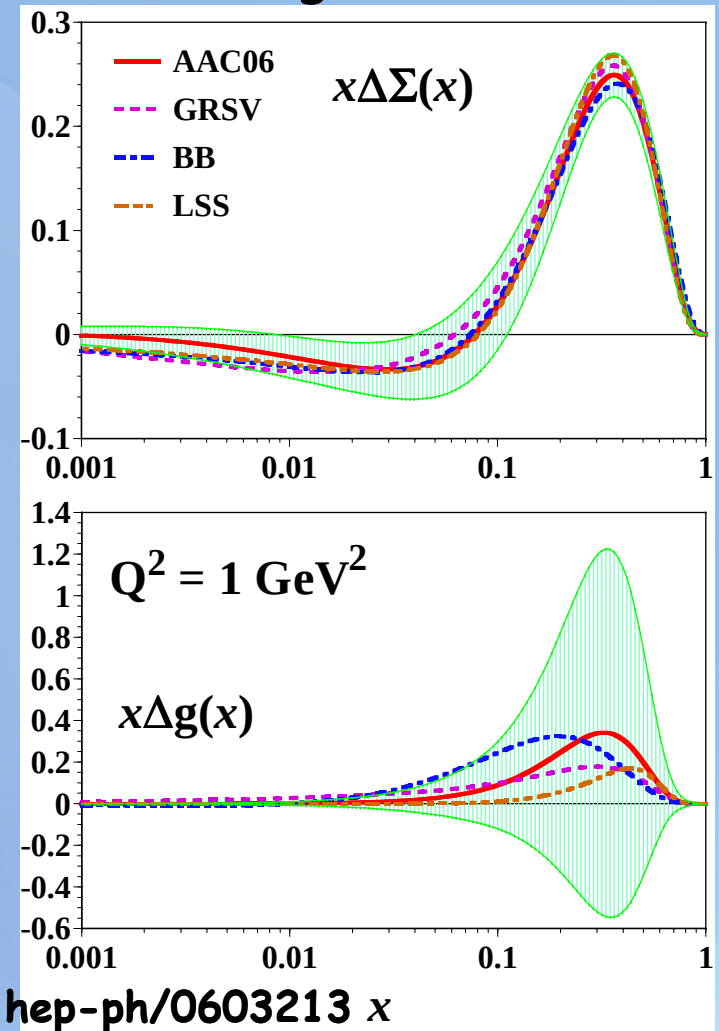


How to measure ΔG

● Indirect from scaling violation

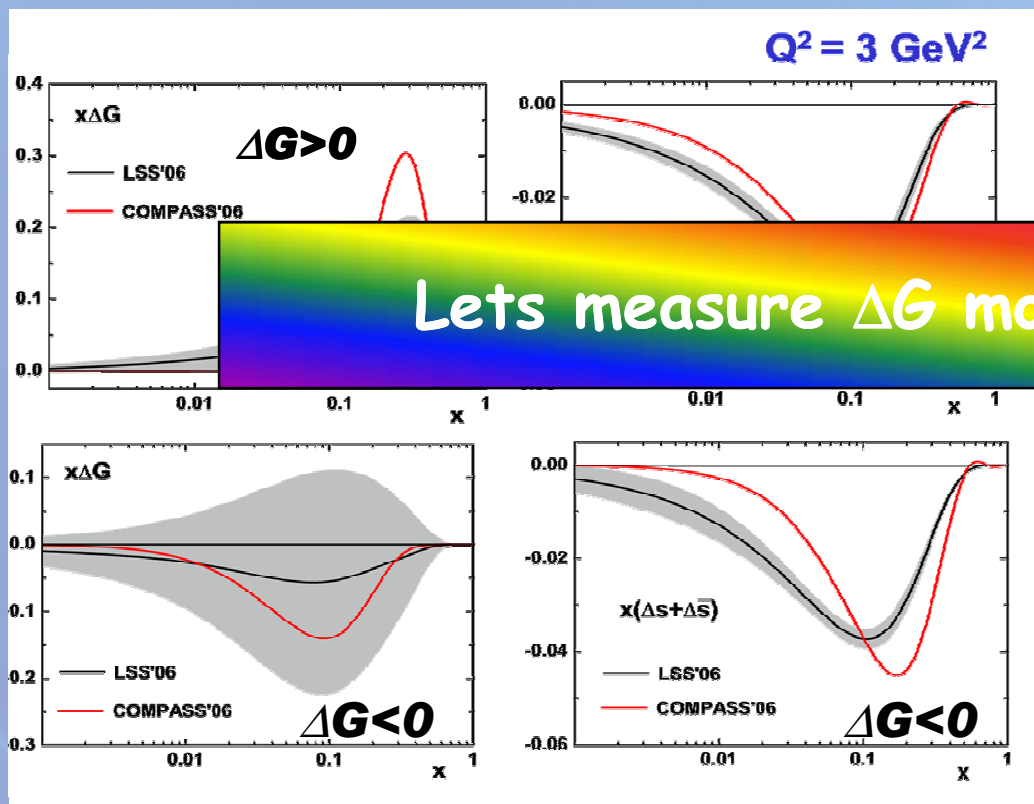


- fixed target experiments
- ➡ small Q^2 - x lever arm
- ➡ even sign of ΔG unknown

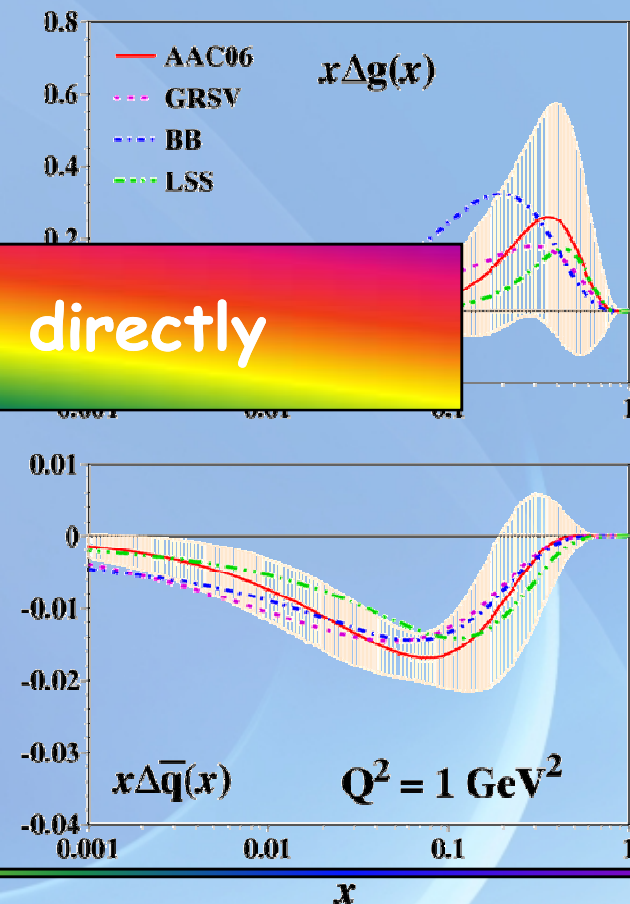


Several more fits

- using mainly only inclusive data or a combination of inclusive and some semi-inclusive data
- Results for ΔG still completely all over the place
- Need a consistent approach for fit and uncertainty determination with all world data taken into account



Lets measure ΔG more directly



The golden channels

Idea: Direct measurement of ΔG

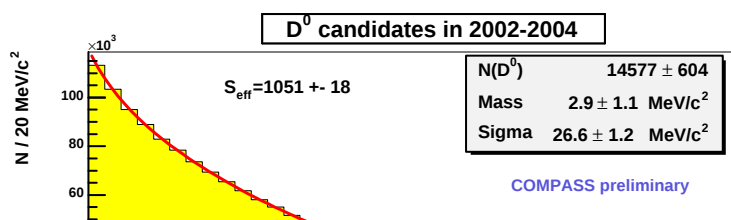
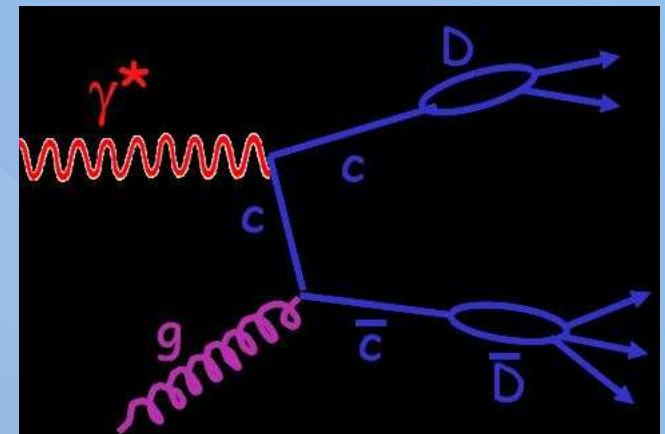
➡ Isolate the photon gluon fusion process (PGF)

● Open Charm production

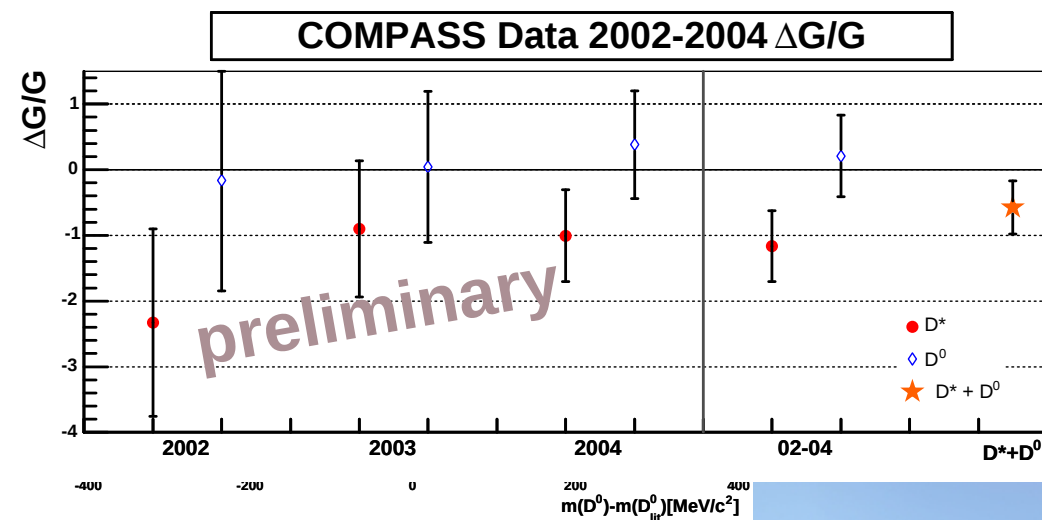
➡ Reaction: $\gamma^* p \rightarrow \bar{D}^0 + D^0 + X$

➡ $c\bar{c} \Rightarrow$ reconstruct D^*, D^0

$$A_{\gamma N}^{PGF} = \frac{\int d\hat{s} \Delta\sigma^{PGF} \Delta G(x_g, \hat{s})}{\int d\hat{s} \Delta\sigma^{PGF}} \approx \langle a_{LL}^{PGF} \rangle \frac{\Delta G}{G}$$



LO-MC: Aroma



$$\left\langle \frac{\Delta G}{G} \right\rangle = -0.57 \pm 0.41(\text{stat.}) \pm 0.17$$

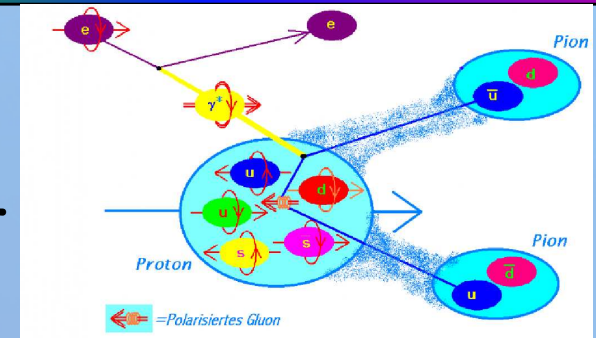
$$x_G \square 0.15$$

$$\mu^2 = 13 \text{ GeV}^2$$

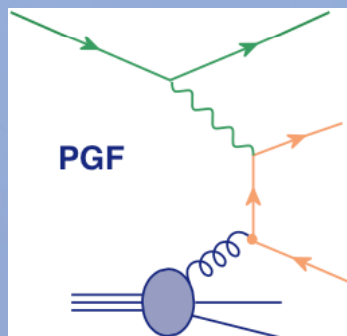
The golden channels

Idea: Direct measurement of ΔG

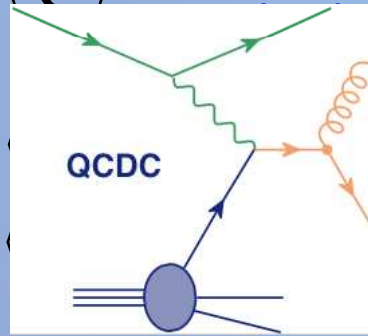
- ➡ Isolate the photon gluon fusion process
- detection of hadronic final states with high p_T
 - ➡ high p_T pairs of hadrons
 - ➡ single high p_T hadrons



- Several possible contributions to the measured asymmetry
 - MC needed to determine R and a_{LL}
- $\langle Q^2 \rangle > 1$ less sub-processes contributing ☺
 $\langle Q^2 \rangle < 0.1$ more sub-processes contributing ☹

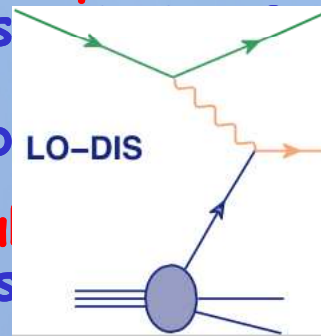


+



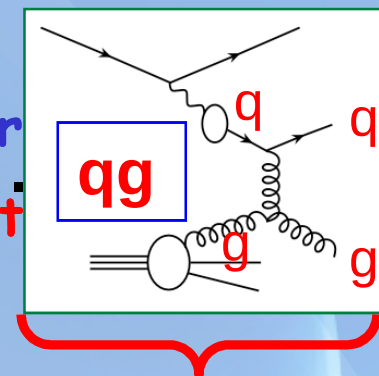
sub
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er s

+



contr
cont

+



Important
at $Q^2 < 0.1$

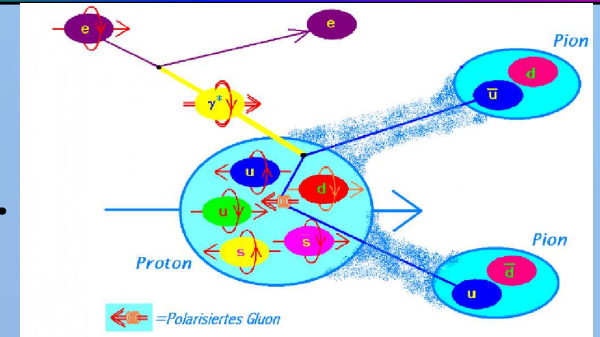
$$A_{||}^{meas}(p_t) = \sum_i f_i A_{||}^i = f_{Bg} A_{||}^{Bg} + f_{Sig} A_{||}^{Sig}; \quad f_i = \frac{\sigma_i}{\sigma_{tot}}$$

$$\left\langle \frac{\Delta G}{G} \right\rangle = \frac{1}{f_{Sig} \langle \hat{a} \rangle} [A_{||}^{meas} - f_{Bg} A_{||}^{Bg}]$$

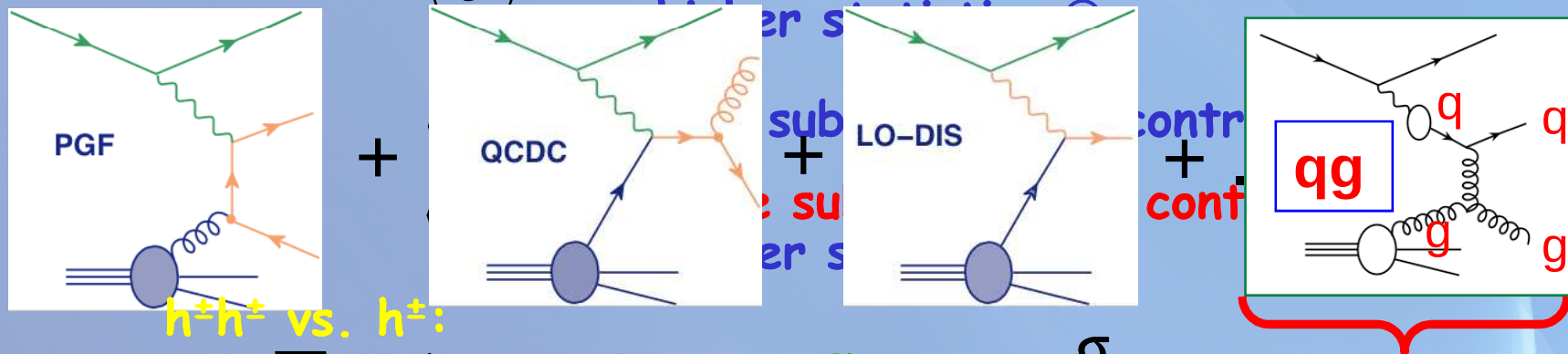
The golden channels

Idea: Direct measurement of ΔG

- Isolate the photon gluon fusion process
- detection of hadronic final states with high p_T
 - high p_T pairs of hadrons
 - single high p_T hadrons



- Several possible contributions to the measured asymmetry
- MC needed to determine R and a_1



$A_{||}^{meas}(h^\pm) = \sum_i f_i A_{||}^{i,meas} \rightarrow pQCD \text{ NLO calculations (easier) possible at } Q^2 < 0.1$

$$\left\langle \frac{\Delta G}{G} \right\rangle = \frac{1}{f_{Sig} \langle \hat{a} \rangle} [A_{||}^{meas} - f_{Bg} A_{||}^{Bg}]$$

COMPASS Results

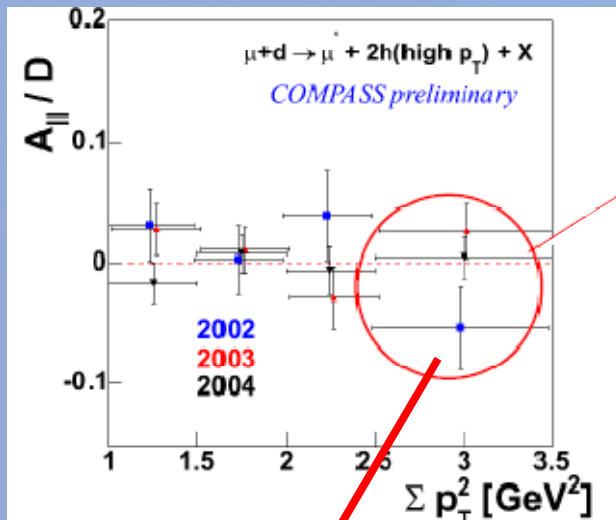
● Channel: $h^\pm h^\pm$

● Cuts: $Q^2 > 1 \text{ GeV}^2$

$$\frac{A_{\parallel}}{D} \Big|_{p_T^1, p_T^2 > 0.7 \text{ GeV}, p_{\parallel}^2 + p_{\perp}^2 \geq 2.5 \text{ GeV}^2, D x_F > 0.1, z > 0.1} = 0.015 \pm 0.080(\text{stat}) \pm 0.013(\text{syst})$$

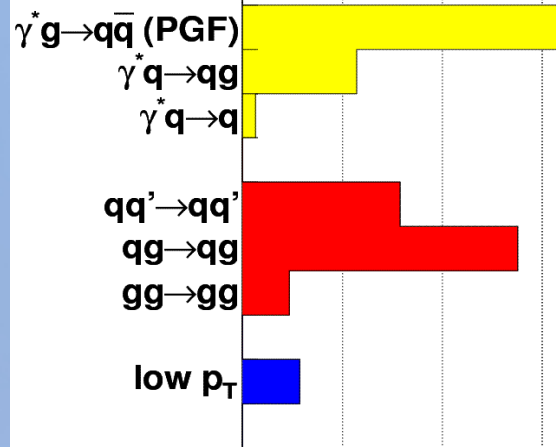
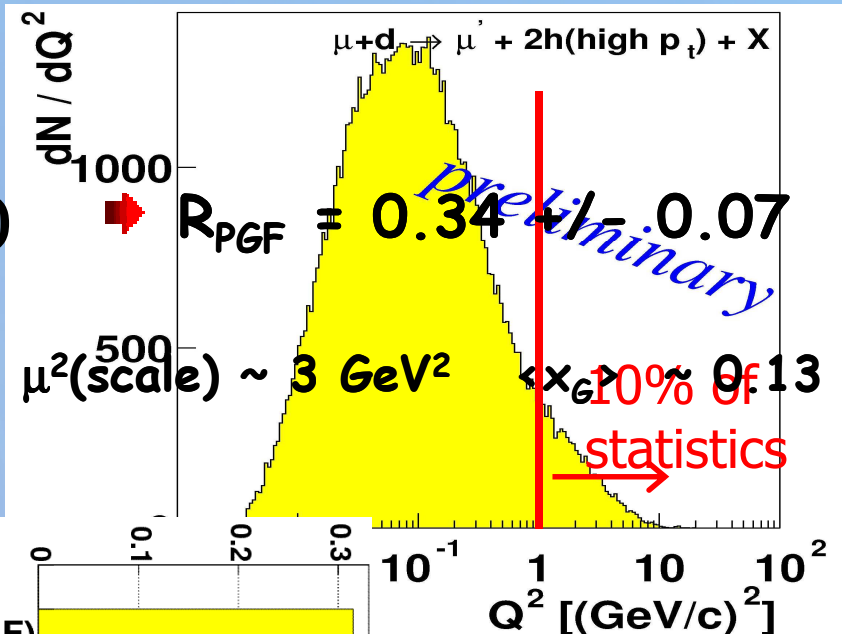
$$\frac{\Delta g}{g} \Big|_{m(h_1, h_2) > 1.5 \text{ GeV}} = 0.06 \pm 0.31(\text{stat}) \pm 0.06(\text{syst})$$

$Q^2 < 1 \text{ GeV}^2$



used for $\Delta g/g$ extraction

&



$$\Delta g/g = 0.016 \pm 0.058(\text{stat.}) \pm 0.055(\text{syst.})$$

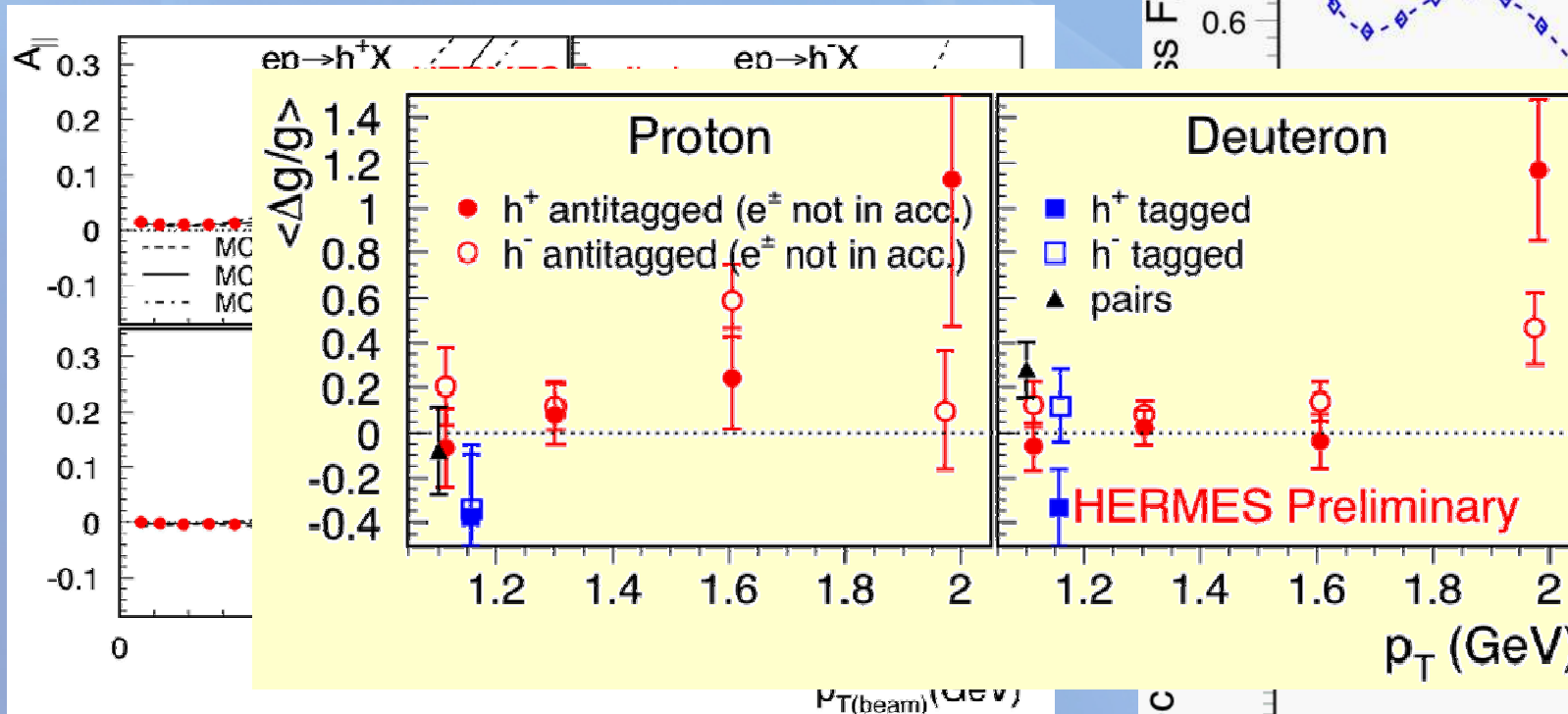
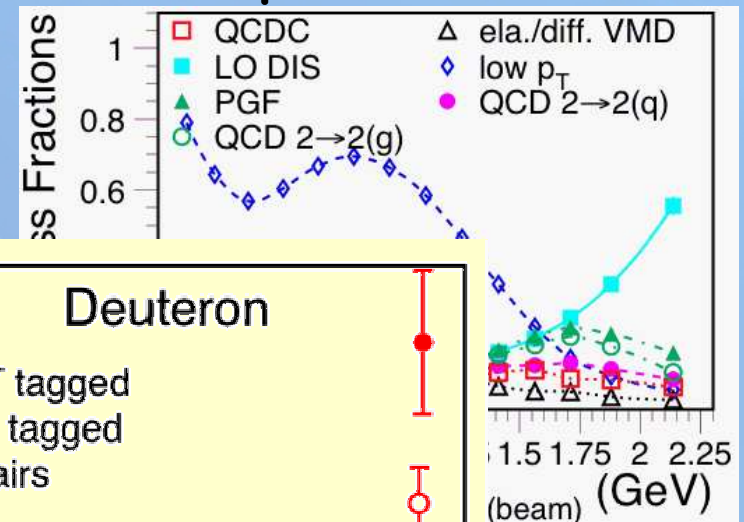
$$\langle x_g \rangle = 0.085^{+0.07}_{-0.035} \quad \mu^2 = 3.0 \text{ GeV}^2$$



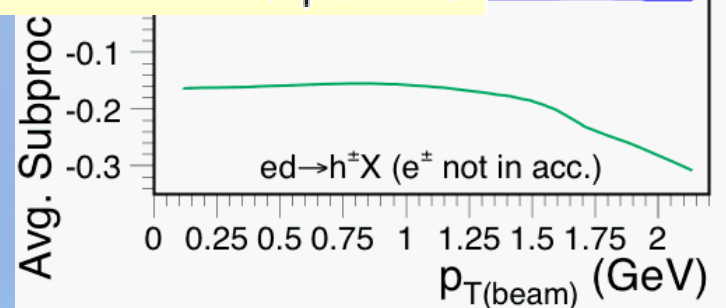
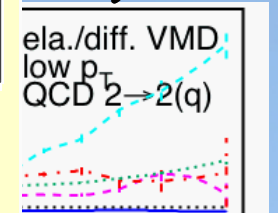
HERMES Results

● Channels: $h^\pm h^\pm$, $h^\pm$

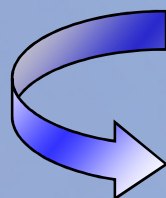
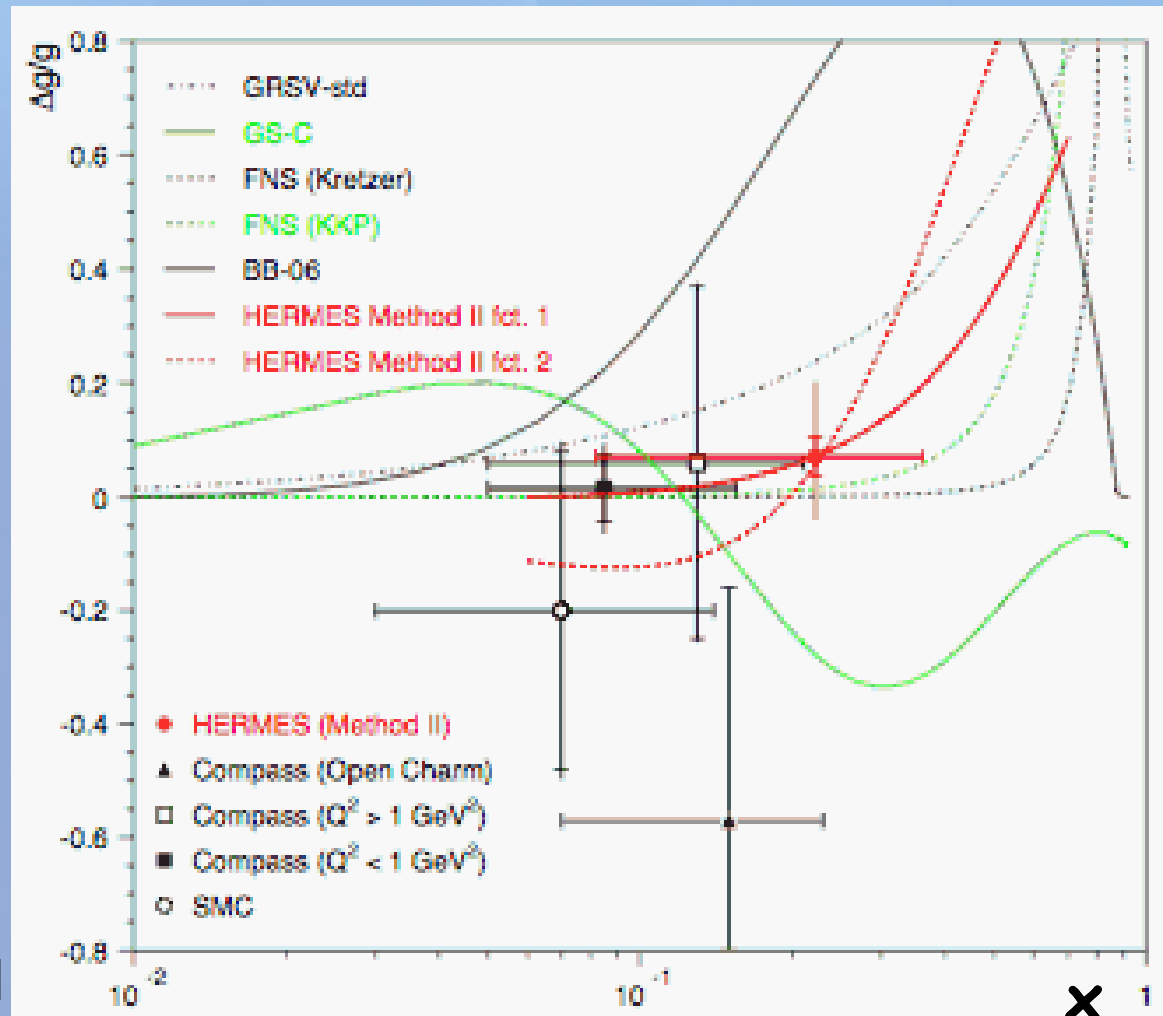
Subprocess Fractions



metries
std.)

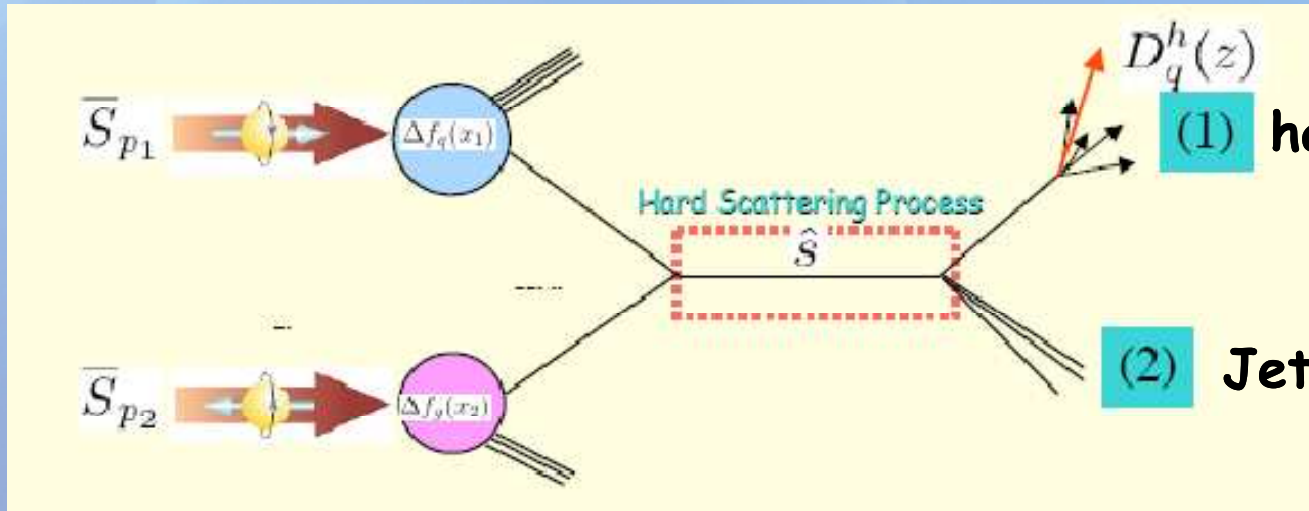


World Data on $\Delta G/G$

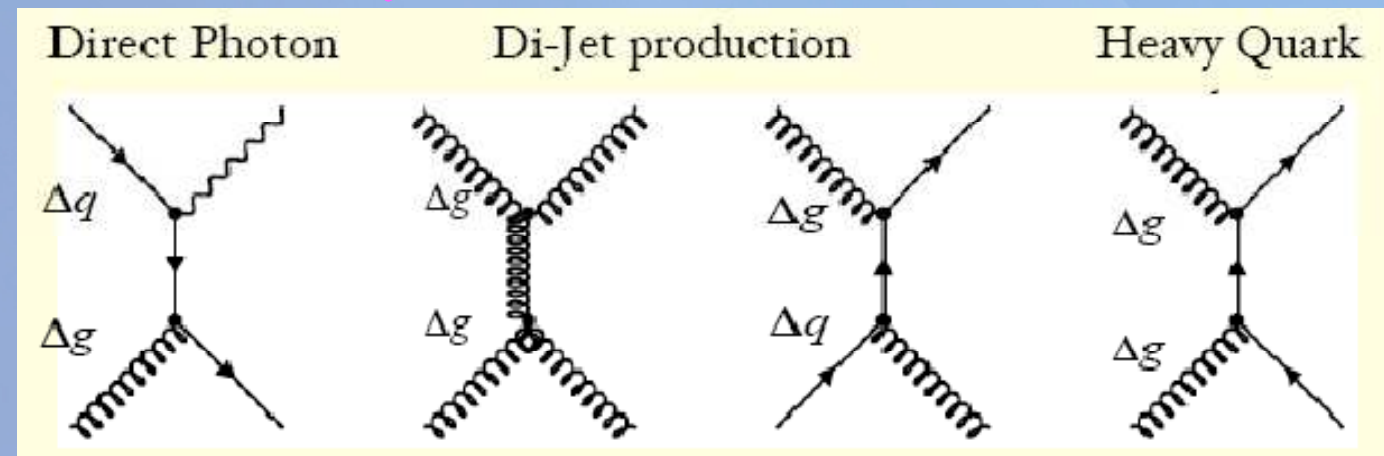


Long way to go till $\Delta g(x)$

Reaction mechanism:



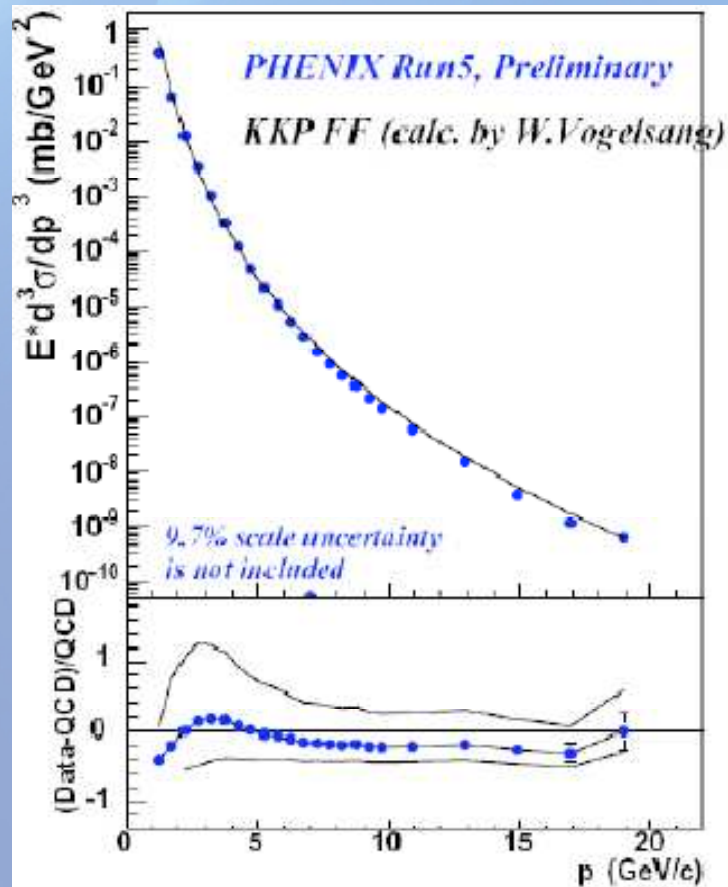
Hard sub processes:



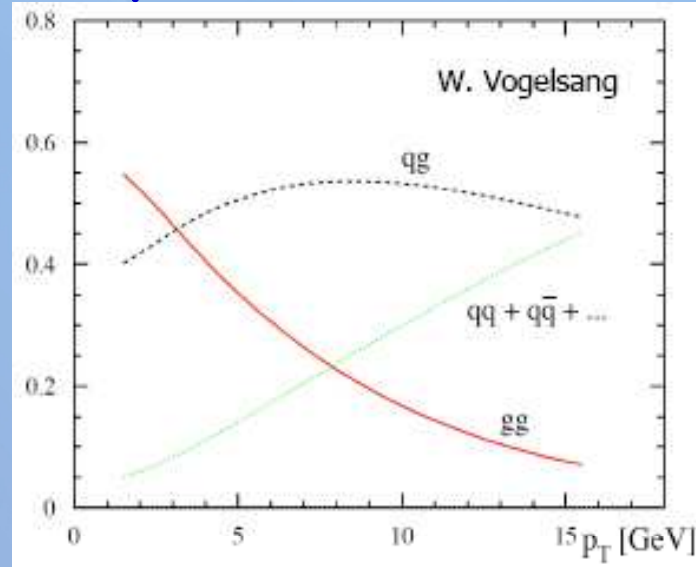
π^0 Production



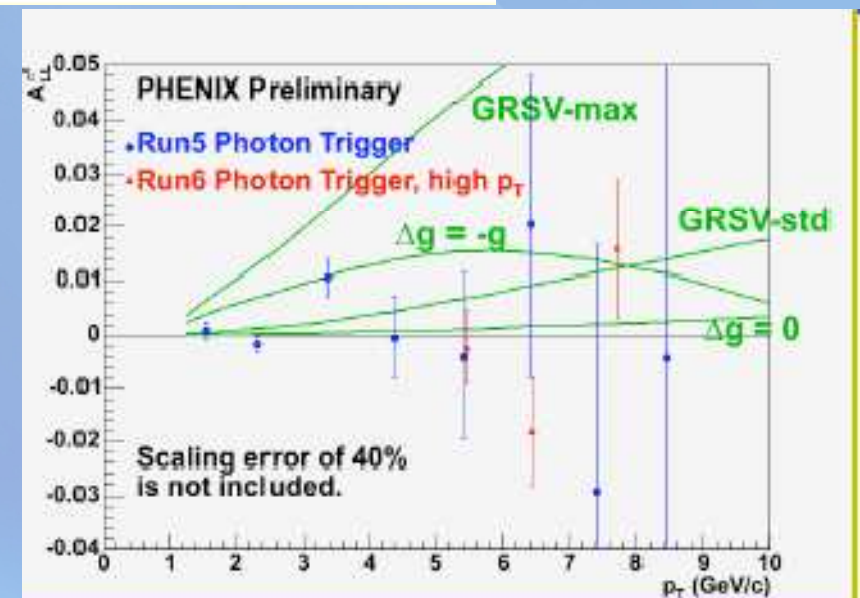
unpolarized cross section:



sub processes:



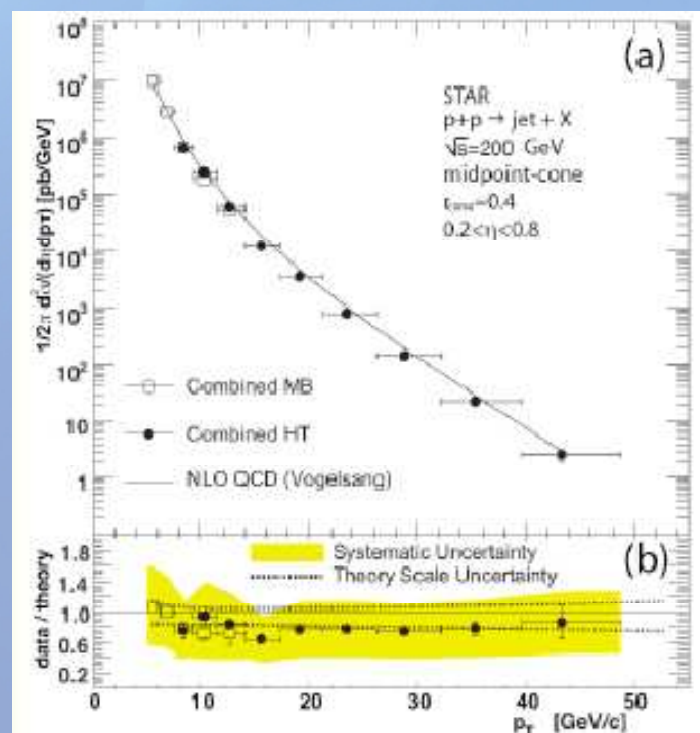
➡ good agreement with NLO pQCD



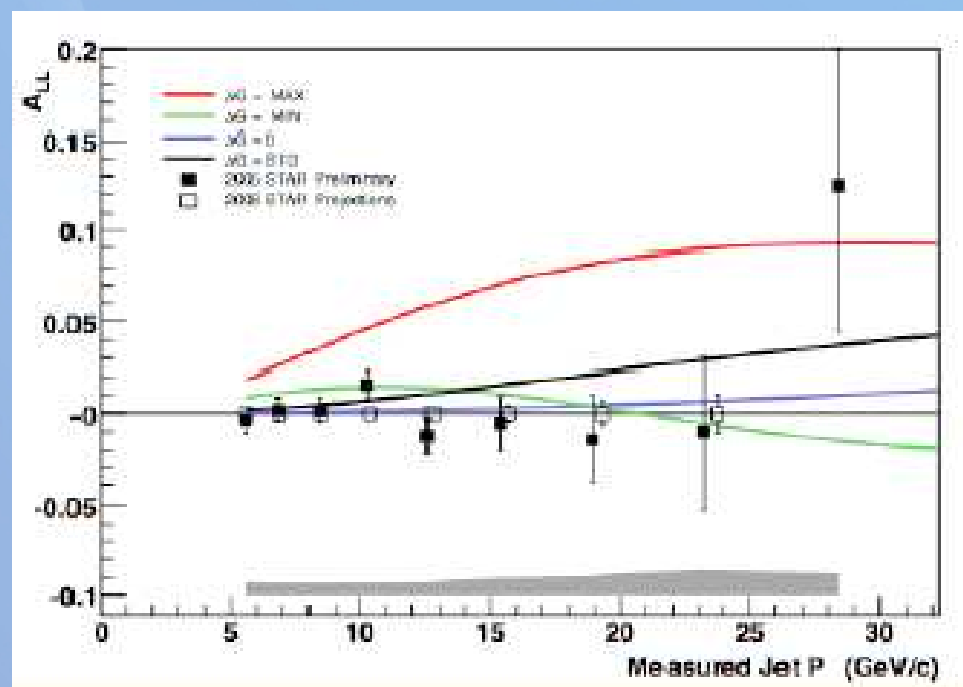
Jet Production



unpolarized cross section:

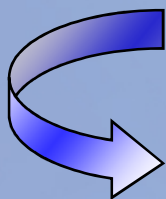
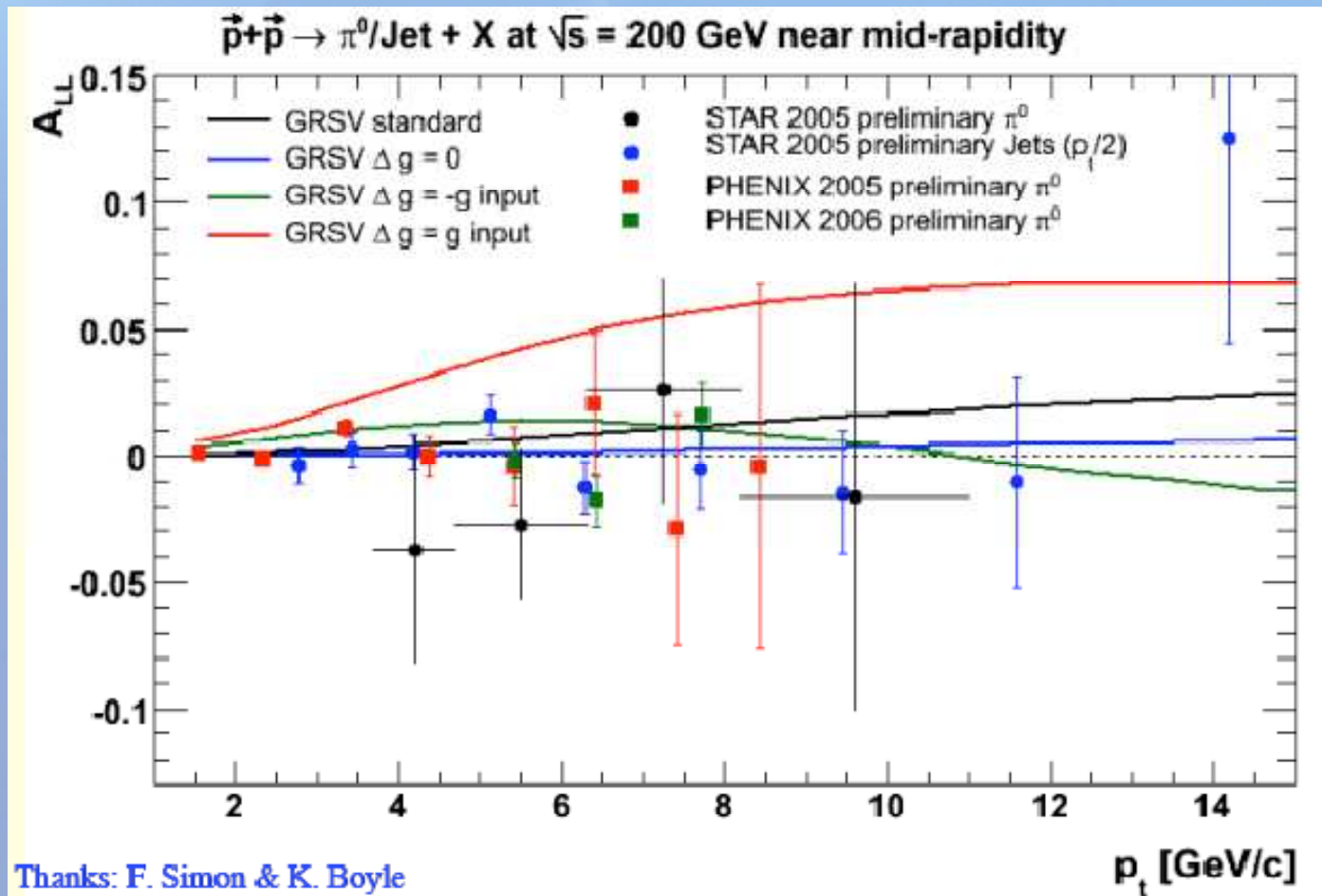


Asymmetry:



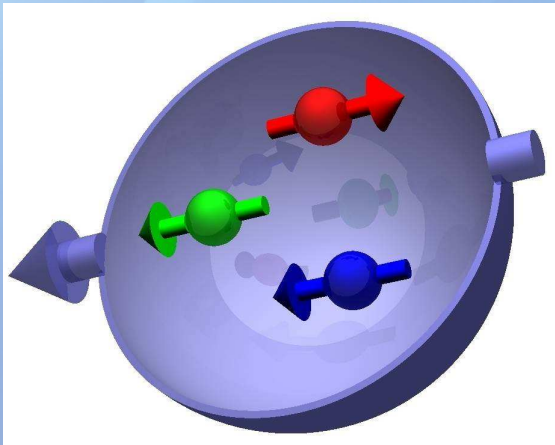
➡ good agreement with NLO pQCD

➡ ΔG seems small



maybe the way is a bit shorter till $\Delta g(x)$

Quark Orbital Angular Momentum



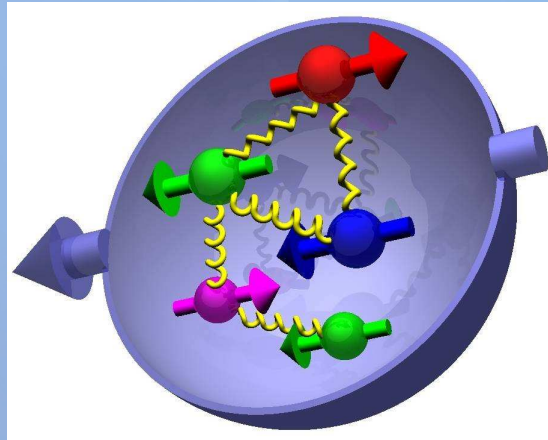
Naïve parton model

$$\Delta u_v = \frac{4}{3} \quad \Delta d_v = -\frac{1}{3}$$

BUT

1989 EMC measured
 $\Sigma = 0.120 \pm 0.094 \pm 0.138$

⇒ Spin Puzzle

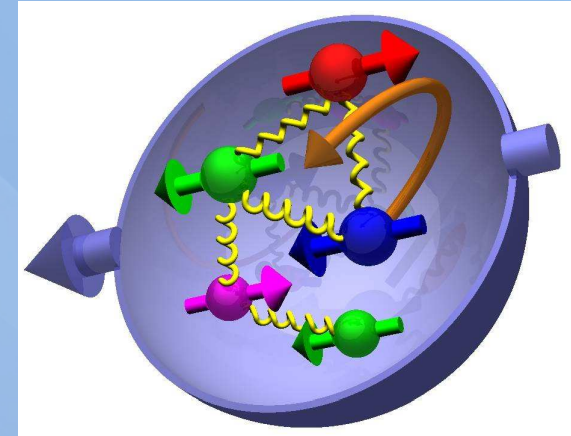


Unpolarised structure fct.

Gluons are important !

⇒ ΔG

⇒ Sea quarks Δq_s



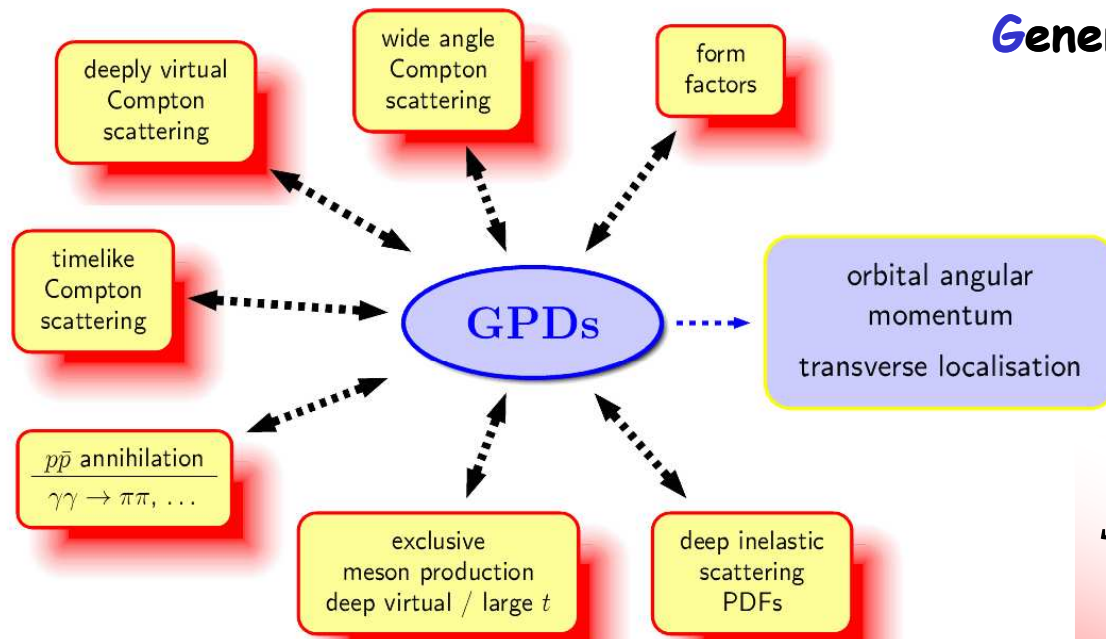
Full description of J_q and J_g
 needs
 orbital angular momentum

$$\frac{1}{2} = \frac{1}{2} \left(\underbrace{\Delta u_v + \Delta d_v + \Delta q_s}_{(\Delta u_s + \Delta d_s + \Delta \bar{u} + \Delta \bar{d} + \Delta s + \Delta \bar{s})} \right) + \Delta G_q + \Delta G_g + L_g$$

The Hunt for L_q

Study of hard **exclusive processes** leads to a new class of PDFs

Generalized Parton Distributions
 $H^q, E^q, \tilde{H}^q, \tilde{E}^q$



possible access to orbital angular momentum

$$J_q = \frac{1}{2} \left(\int_{-1}^1 x dx (H^q + E^q) \right)_{t \rightarrow 0}$$

$$J_q = \frac{1}{2} \Delta\Sigma + L_q$$

exclusive: all products of the reaction are detected
 $\Rightarrow$ missing energy (ΔE) and missing Mass (M_x) = 0

from DIS:
 HERMES ~0.3

GPDs Introduction

What does GPDs characterize?

unpolarized

polarized

$$H^q(x, \xi, t)$$

$$\tilde{H}^q(x, \xi, t)$$

conserve nucleon helicity

$$H^q(x, 0, 0) = q, \tilde{H}^q(x, 0, 0) = \Delta q$$

$$E^q(x, \xi, t)$$

$$\tilde{E}^q(x, \xi, t)$$

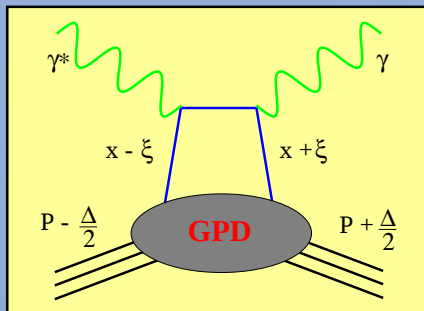
flip nucleon helicity

not accessible in DIS

quantum numbers of final state

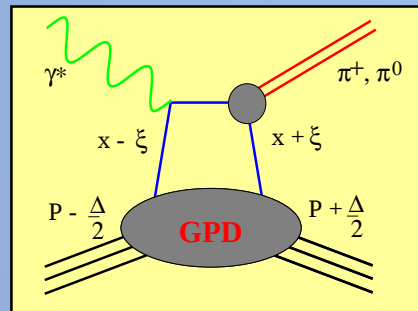


select different GPD



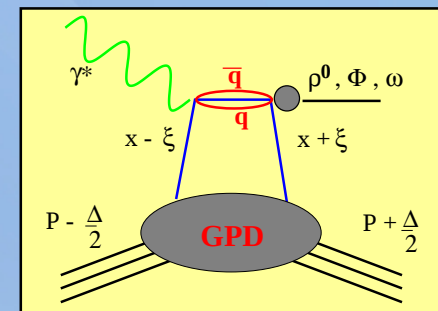
DVCS

$$\underbrace{H^q}_{A_C, A_{LU}}, \underbrace{E^q}_{A_{UT}}, \underbrace{\tilde{H}^q, \tilde{E}^q}_{A_{UL}}$$



pseudo-scalar mesons

$$\underbrace{\tilde{H}^q, \tilde{E}^q}_{A_{UT}, \sigma_{\pi^+}}$$

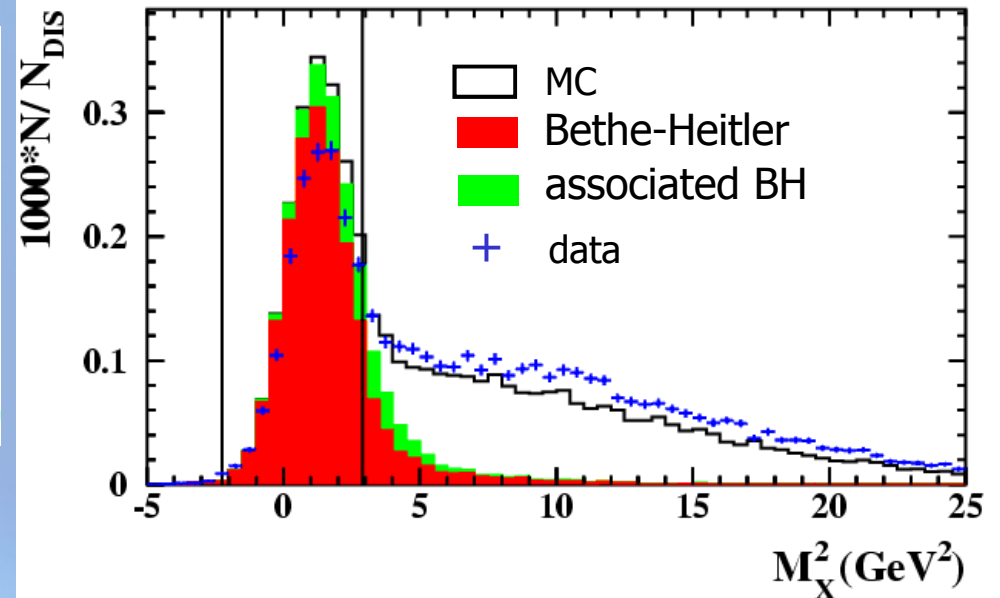
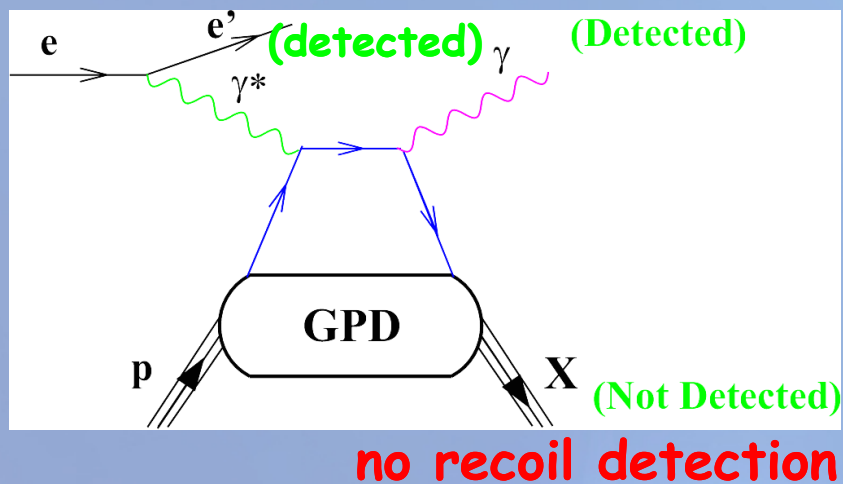
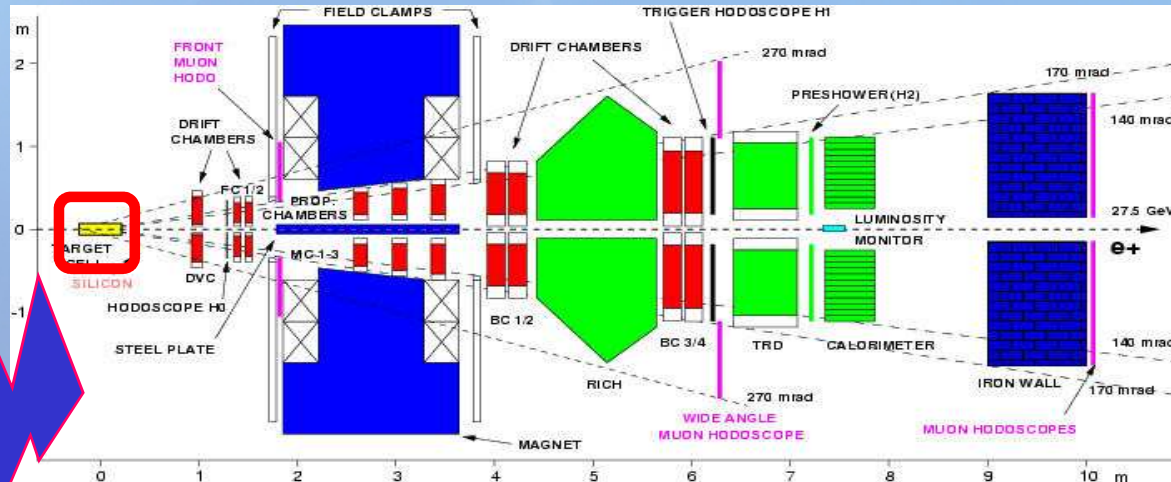


vector mesons

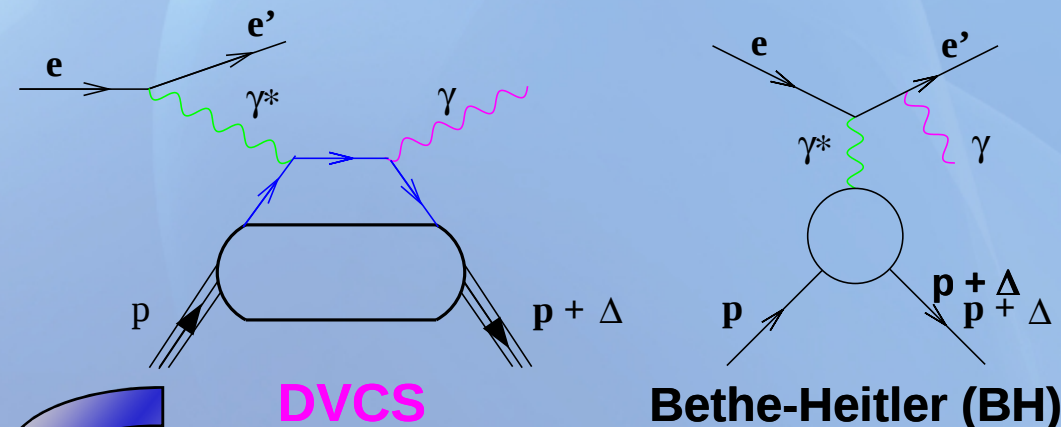
$$\underbrace{H^q, E^q}_{A_{UT}, \sigma_{\rho, \Phi, \omega}}$$

Exclusivity @ HERMES

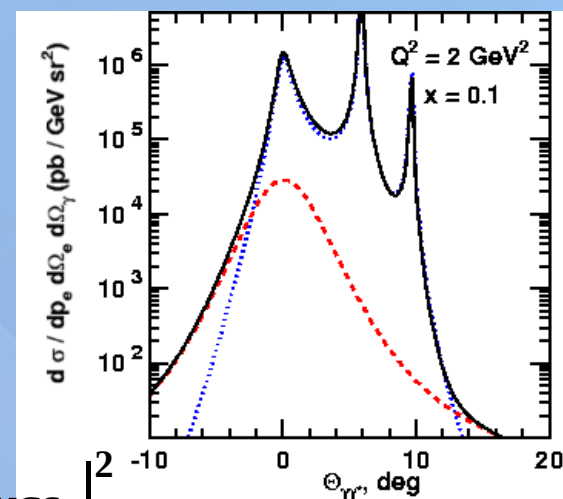
e^+/e^-
 $\rightarrow$
 27.5 GeV
 $P_b = 55\%$



two experimentally undistinguishable processes:



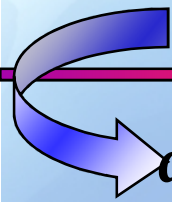
HERMES / JLAB
kinematics: BH $\gg$ DVCS



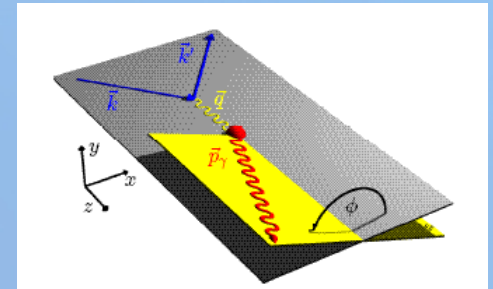
$$d\sigma \sim \left(\tau_{BH}^* \tau_{DVCS} + \tau_{DVCS}^* \tau_{BH} \right) + |\tau_{BH}|^2 + |\tau_{DVCS}|^2$$

isolate BH-DVCS interference term $\Rightarrow$ non-zero azimuthal asymmetries

DVCS ASYMMETRIES



$$d\sigma \sim \left(\tau_{BH}^* \tau_{DVCS} + \tau_{DVCS}^* \tau_{BH} \right) + |\tau_{BH}|^2 + |\tau_{DVCS}|^2$$



→ different charges: $e^+ e^-$ (only @HERA!):

$$\Delta\sigma_c \sim \cos\phi \cdot \text{Re}\{ \mathbf{H} + \xi \tilde{\mathbf{H}} + \dots \}$$

$$\Rightarrow \mathbf{H}$$

→ polarization observables:

$$\Delta\sigma_{LU} \sim \sin\phi \cdot \text{Im}\{ \mathbf{H} + \xi \tilde{\mathbf{H}} + k \mathbf{E} \}$$

$$\Rightarrow \sigma_{UT}$$

$$\Delta\sigma_{UL} \sim \sin\phi \cdot \text{Im}\{ \tilde{\mathbf{H}} + \xi \mathbf{H} + \dots \}$$

$$\begin{array}{cc} \swarrow & \searrow \\ \text{beam} & \tilde{\text{target}} \\ \Rightarrow & \mathbf{H} \end{array}$$

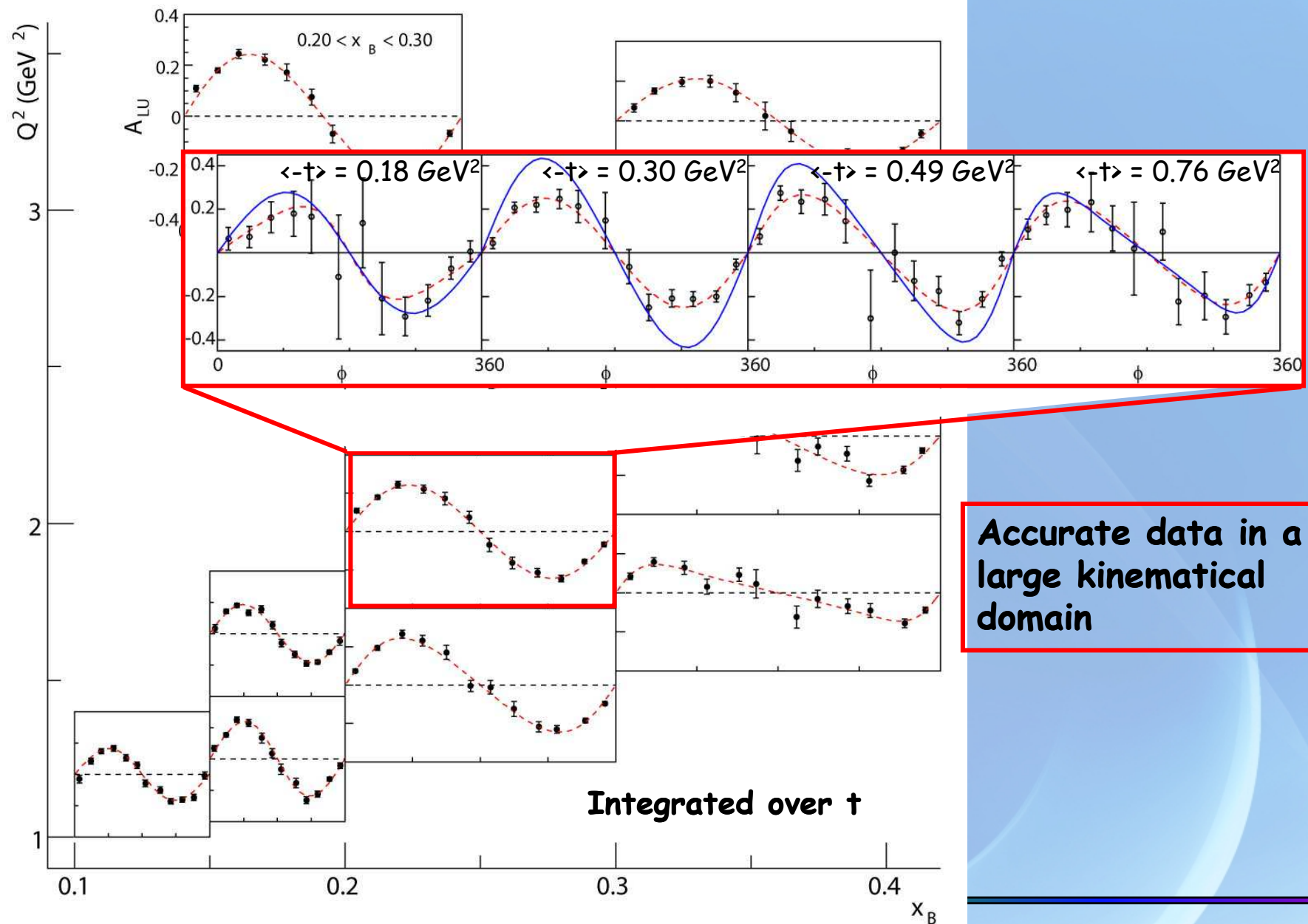
$$\Delta\sigma_{UT} \sim \sin\phi \cdot \text{Im}\{ k(\mathbf{H} - \mathbf{E}) + \dots \}$$

$$\Rightarrow \mathbf{H}, \mathbf{E}$$

$$\xi = x_B/(2-x_B), k = t/4M^2$$

kinematically suppressed

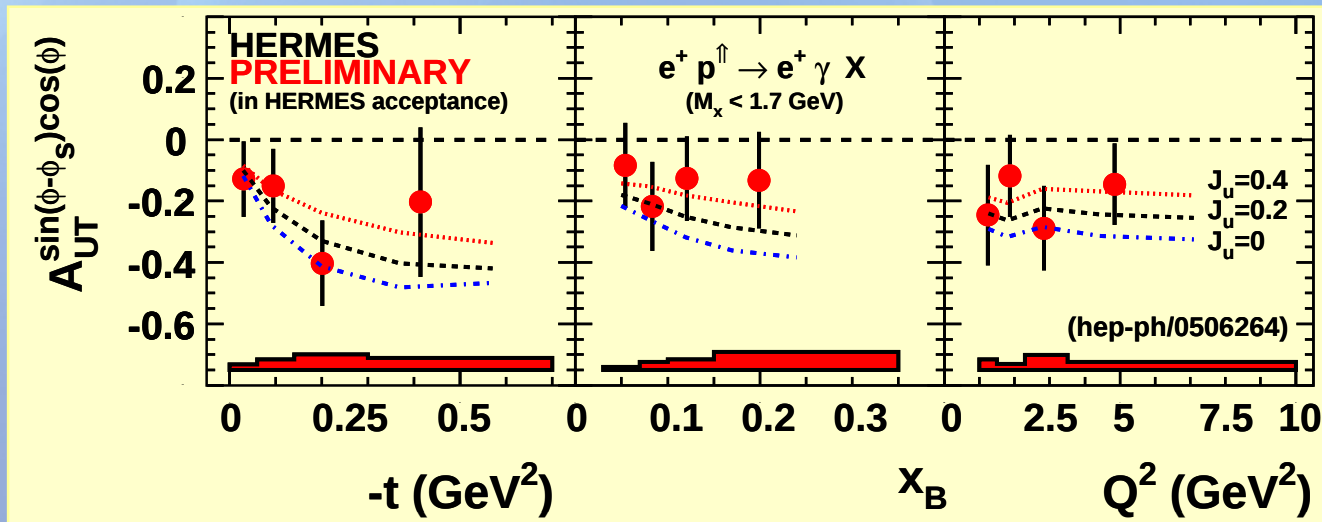
CLAS: DVCS - BSA



Accurate data in a large kinematical domain

A way to E and $J_u - J_d$

Hermes DVCS-TTSA:



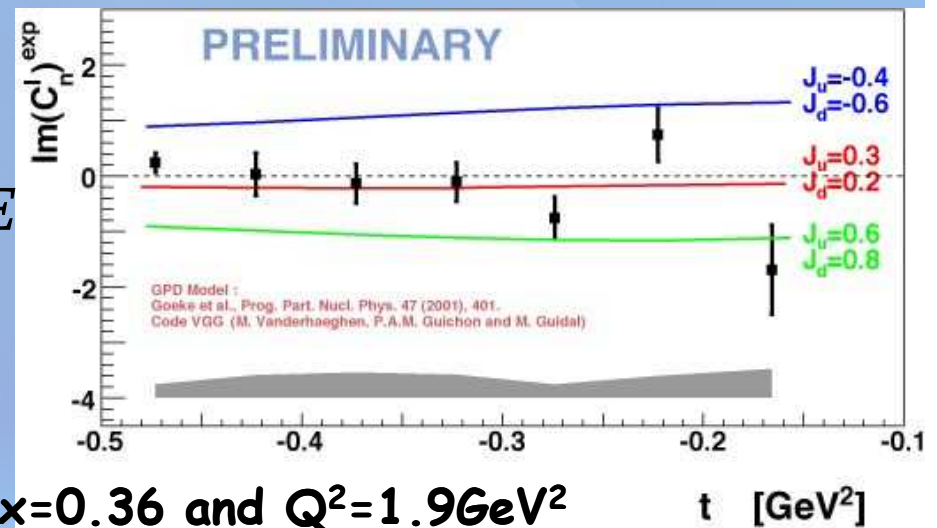
$$A_{UT}^{\sin(\phi - \phi_s) \cos(\phi)} \sim \text{Im}(F_2 H - F_1 E)$$

Hall A nDVCS-BSA:

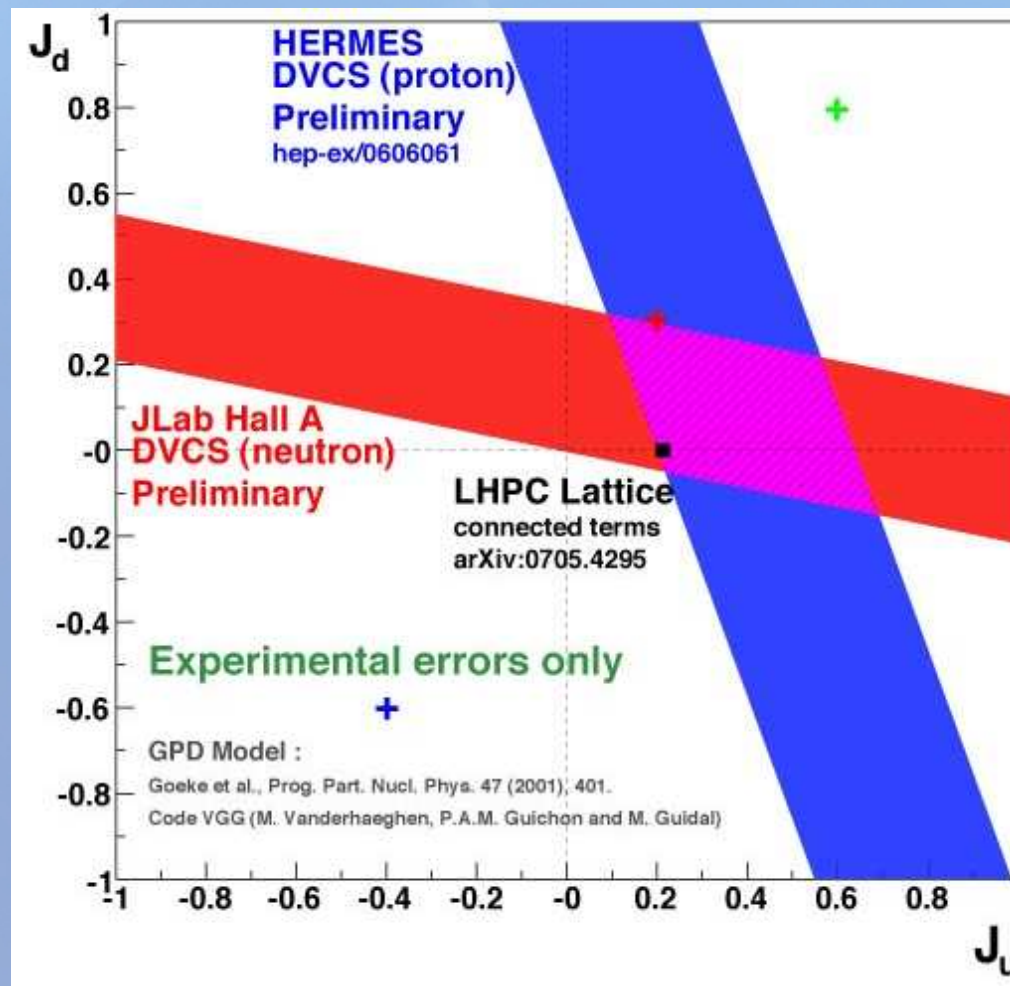
- Neutron obtained combining deuterium and proton

$$A_{LU} \sim F_1 H + x(F_1 + F_2) \tilde{H} - \frac{t}{4M^2} F_2 E$$

- F_1 small u & d cancel in $\tilde{H}$

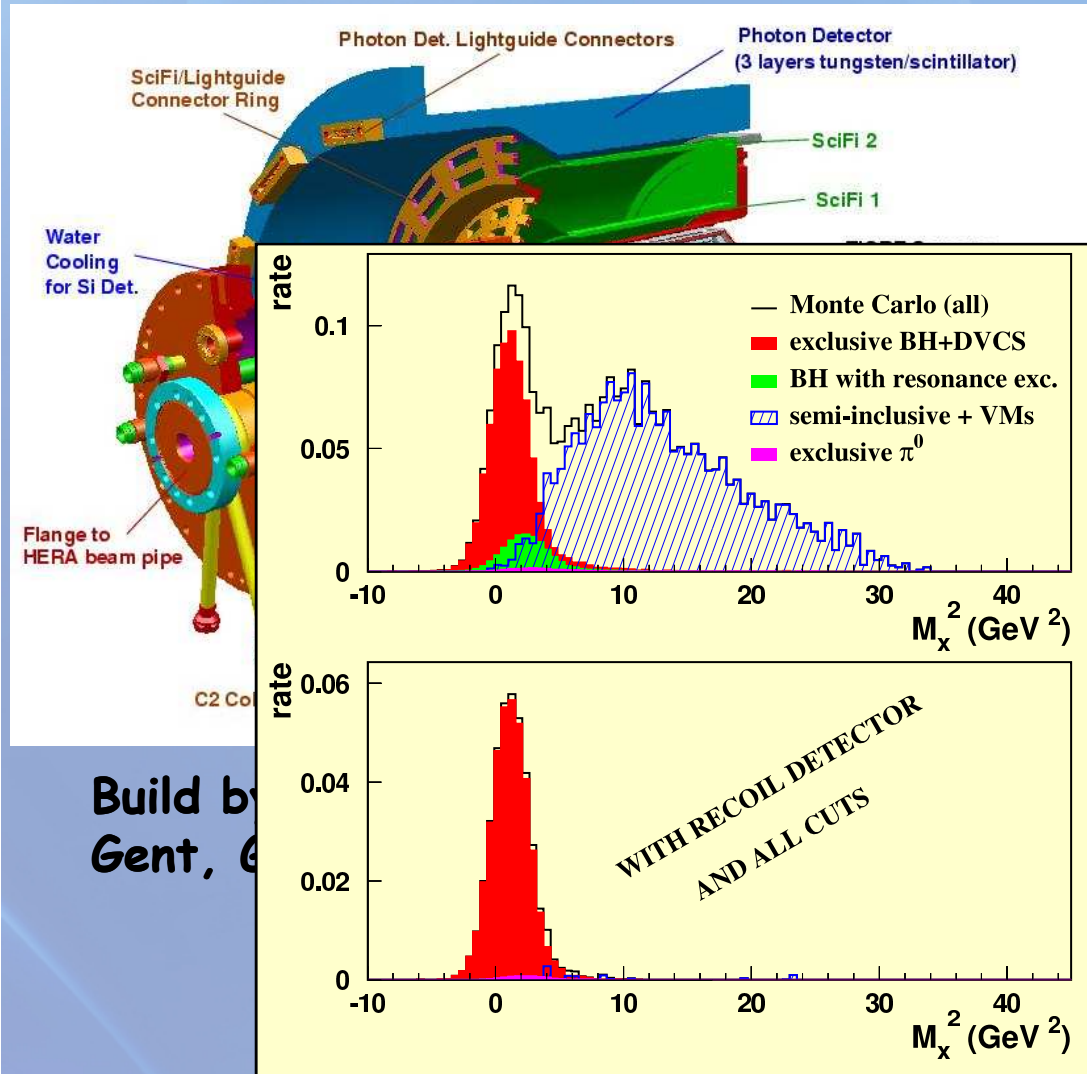


Can we constrain ($J_u - J_d$)



- first model dependent extraction of $J_u - J_d$ possible
VGG-Code: GPD-model: LO/Regge/D-term=0

The HERMES Recoil Detector

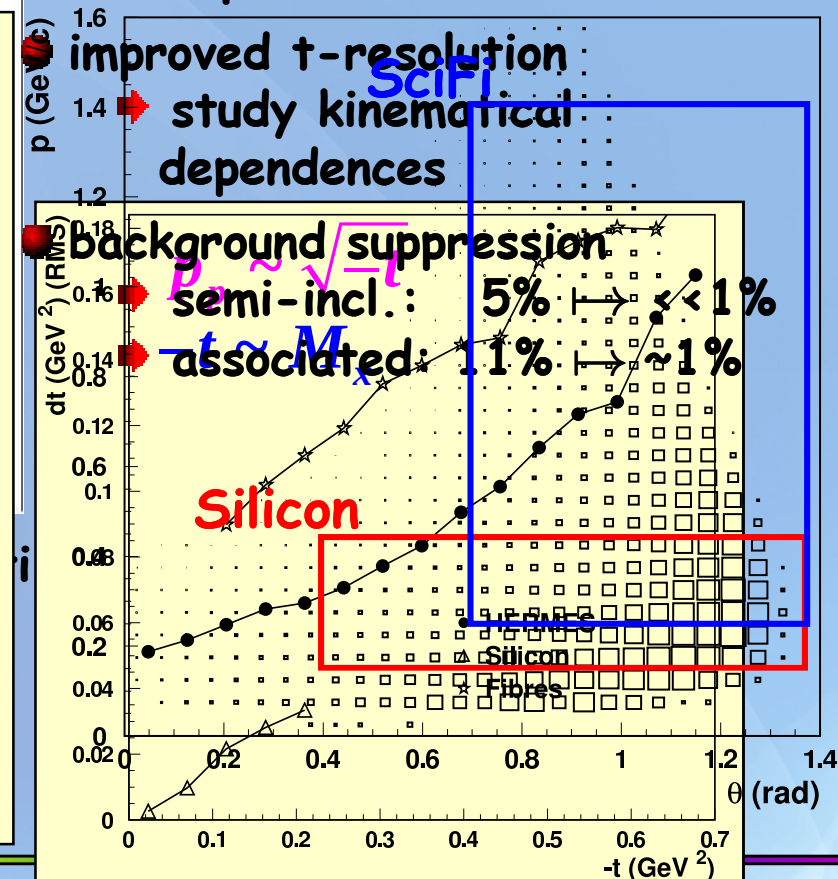


● detection of the recoiling proton

➡ p ; 135 - 1200 MeV/c

➡ 76% ϕ acceptance

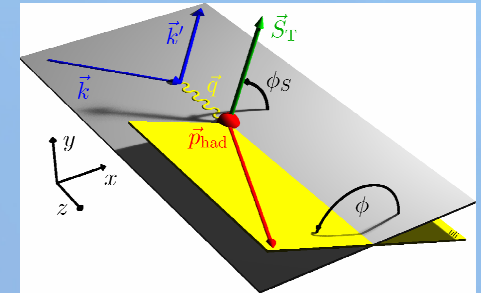
➡ π/p PID via dE/dx



Access to L_q in semi-inclusive scattering

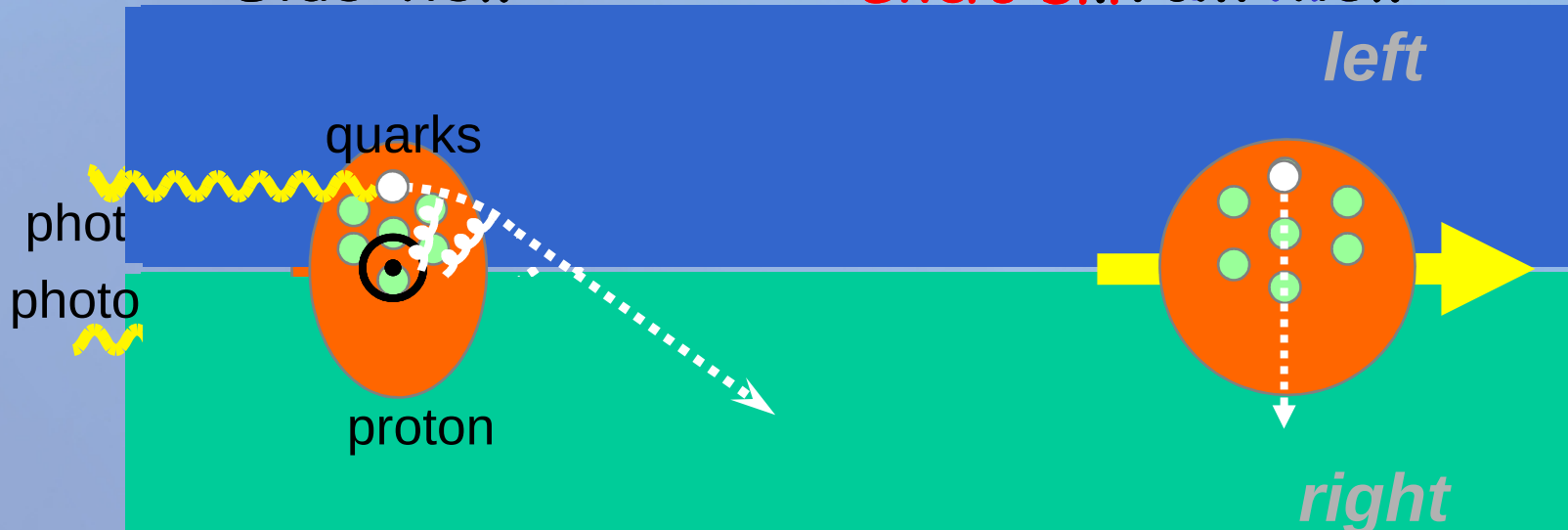
- New structure function accessible with SSA

$$F_{UU,T}^{\sin(\phi_h - \phi_s)}(x, z, P_{h\perp}^2) = C \left[-\frac{\hat{h} \cdot p_T}{M} f_{1T}^\perp D_1 \right]$$



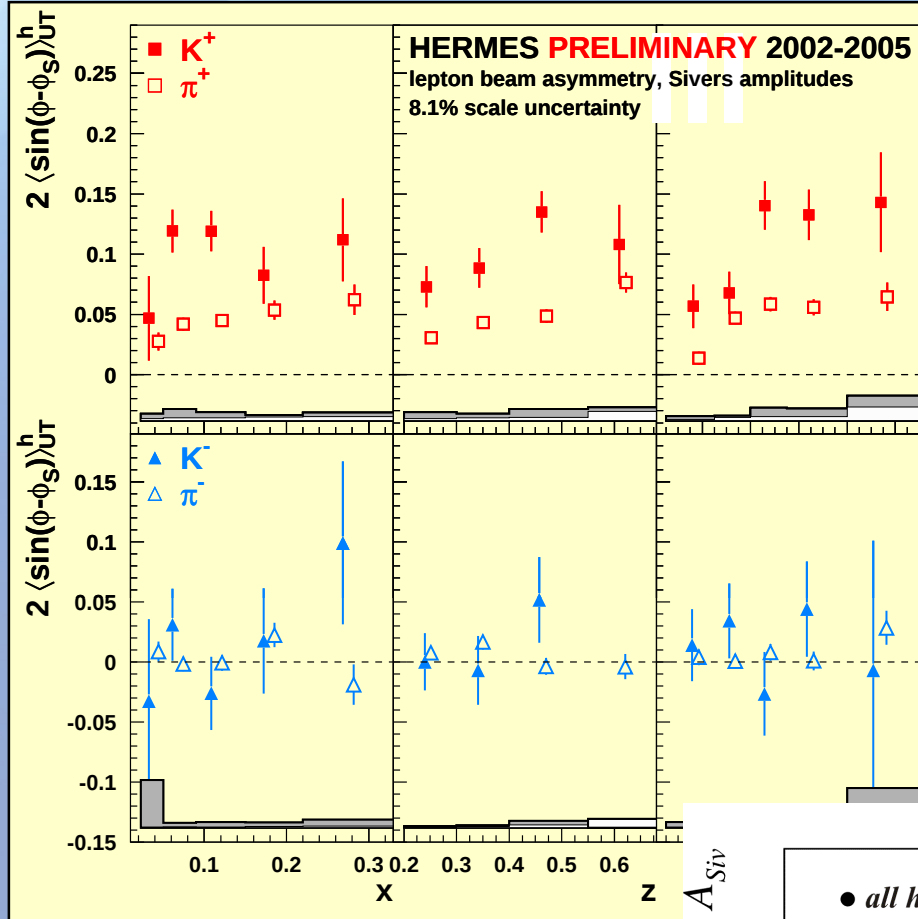
Side view

Sivers DFF Front view



- Not predicted by spin and twist-2 distribution of quarks in transverse space (orbital angular momentum is required)
- A distortion in the distribution of quarks in transverse space can give rise to a nonzero Sivers function

HERMES & COMPASS Measurements

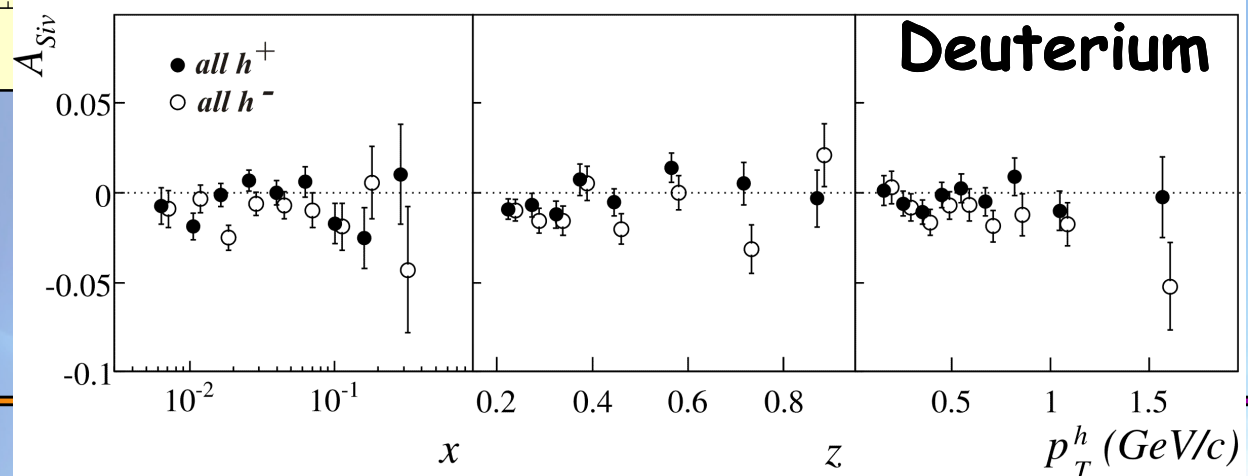


$$A_{Sivers} \propto f_{1T}^\perp(x) D_1(z)$$

- Proton:
- Sivers moment:
 - $\pi^+ > 0$ $\pi^- \sim 0$
 - $K^+ > 0$ $K^- \sim 0$
 - $K^+ > \pi^+$
- ➡ sea quarks important
- Deuterium ~ 0
 - ➡ u and d quark cancel

Proton

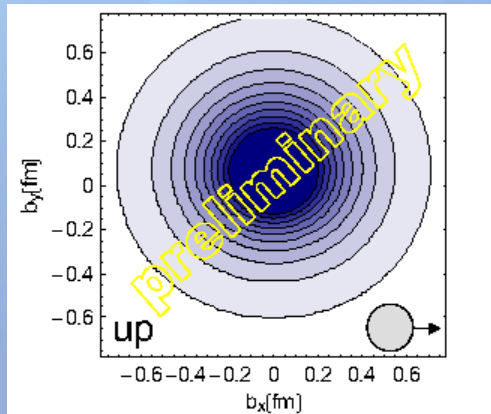
COMPASS 2002-2004



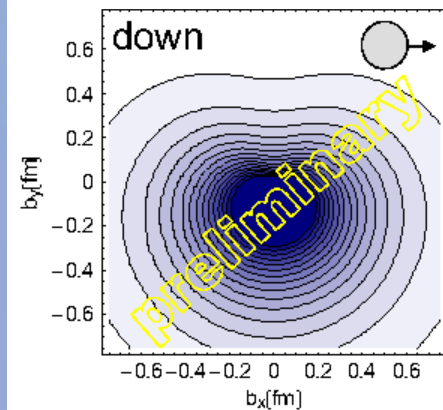
E.C. Aschenauer

Results from theory

Lattice QCD: Sivers



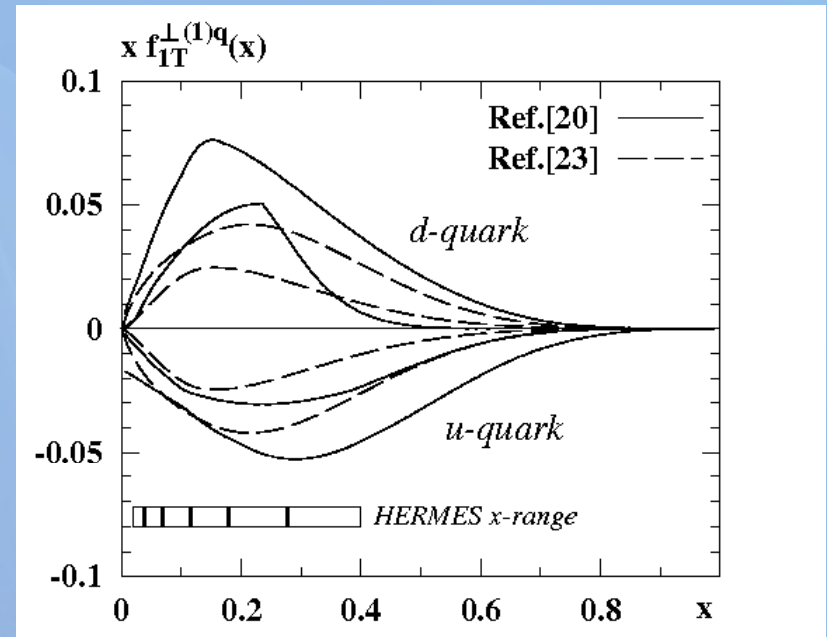
negative for up quarks



positive for down quarks

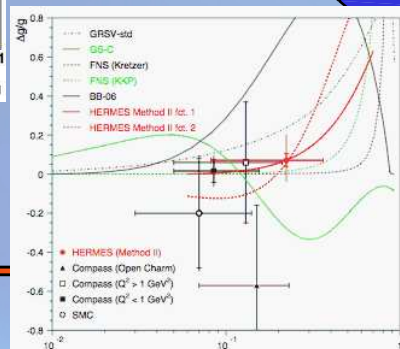
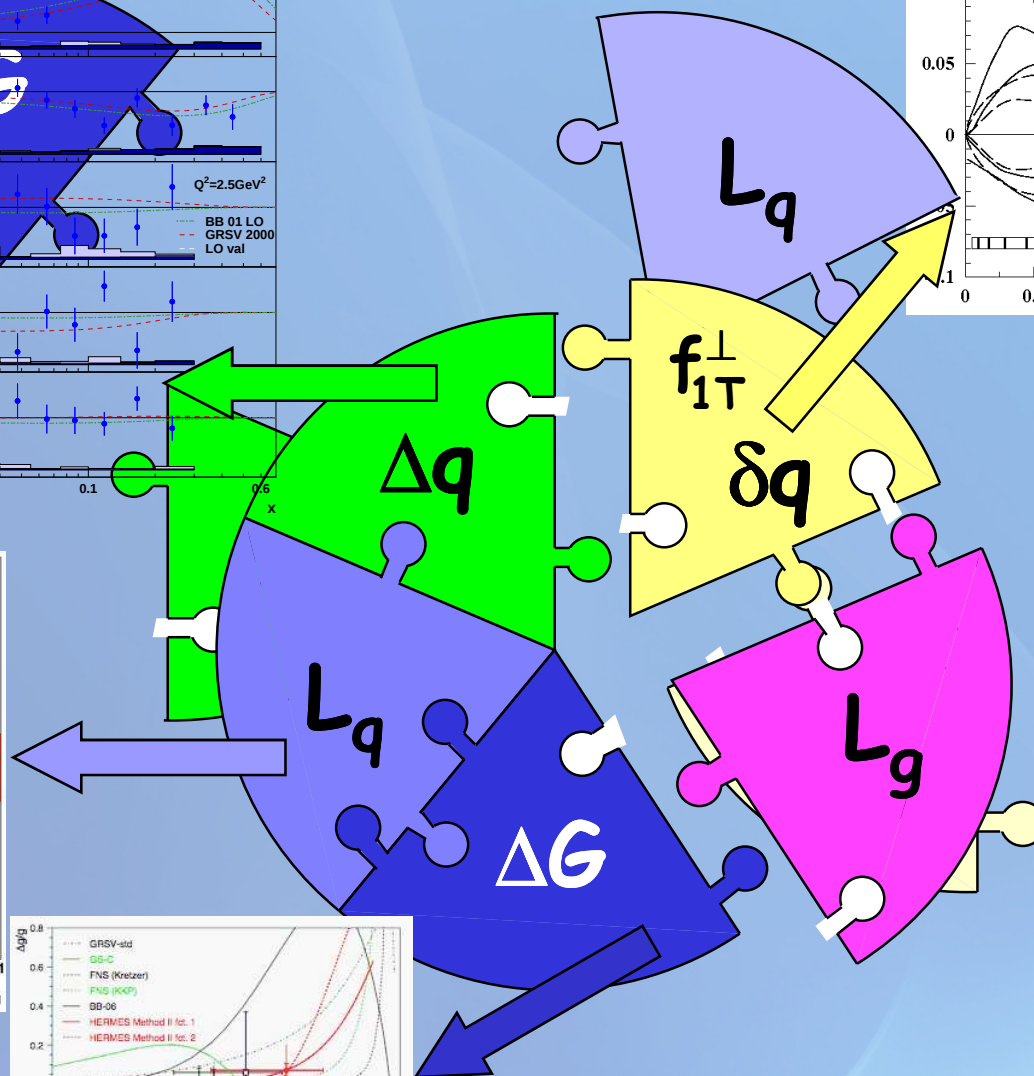
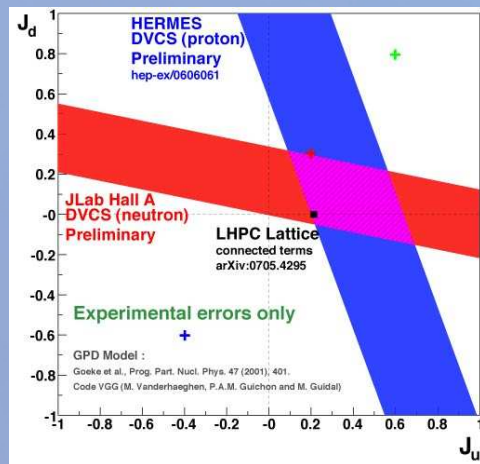
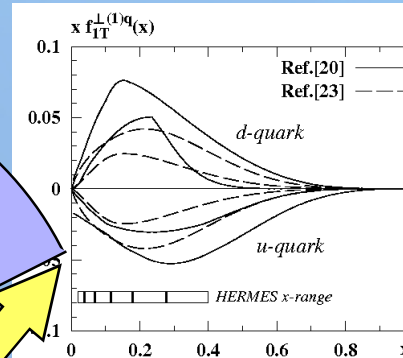
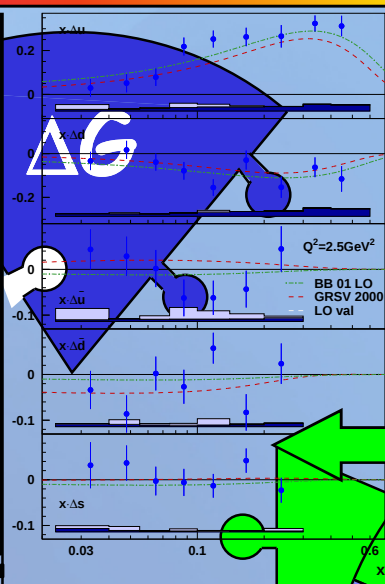
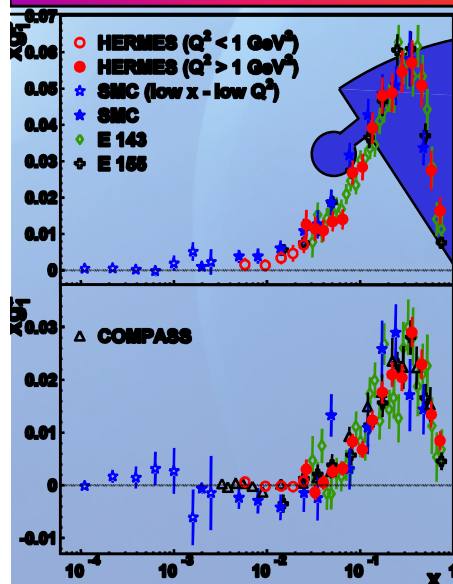
QCDSF/UKQCD Collab. (hep-ph/05110032)

Pheno. analysis from data:



- Anselmino et al., hep-ph/0511017
[20] Anselmino et al., PRD72 (05)
[21] Vogelsang, Yuan, PRD72 (05)
[23] Collins et al., hep-ph/0510342

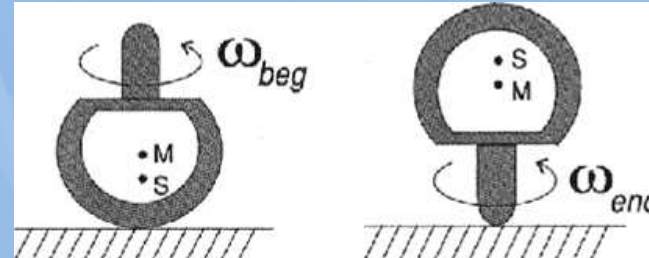
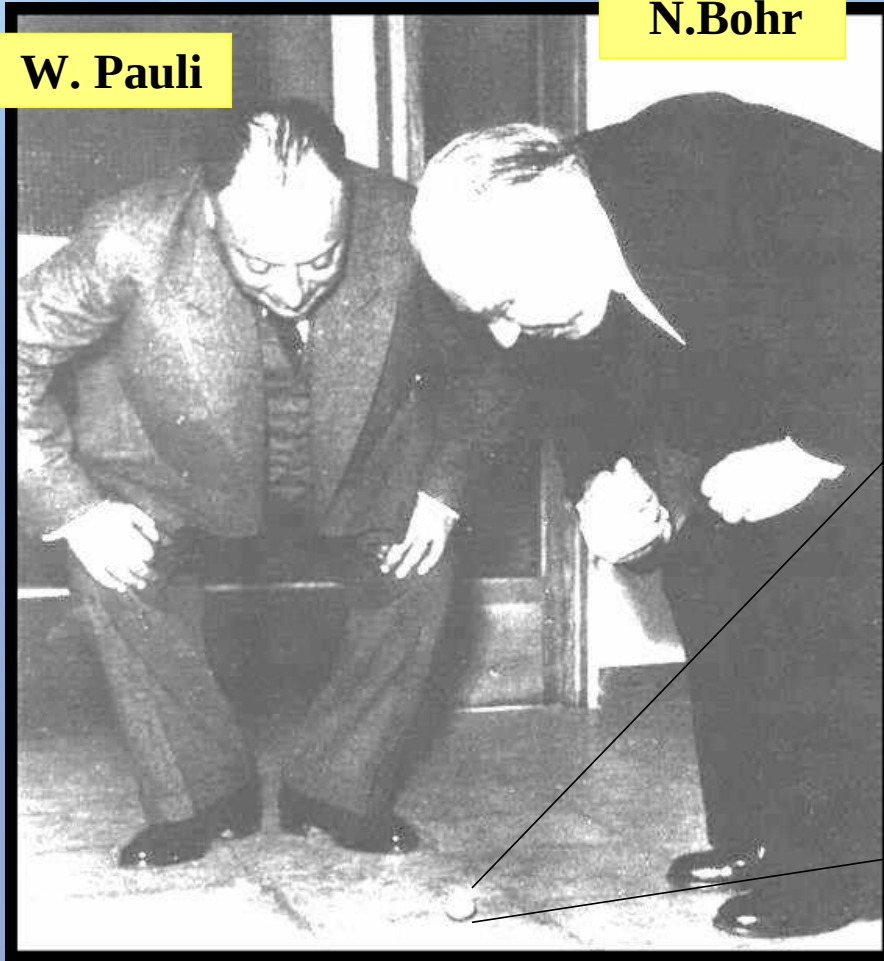
Summary



Spin is fascinating

W. Pauli

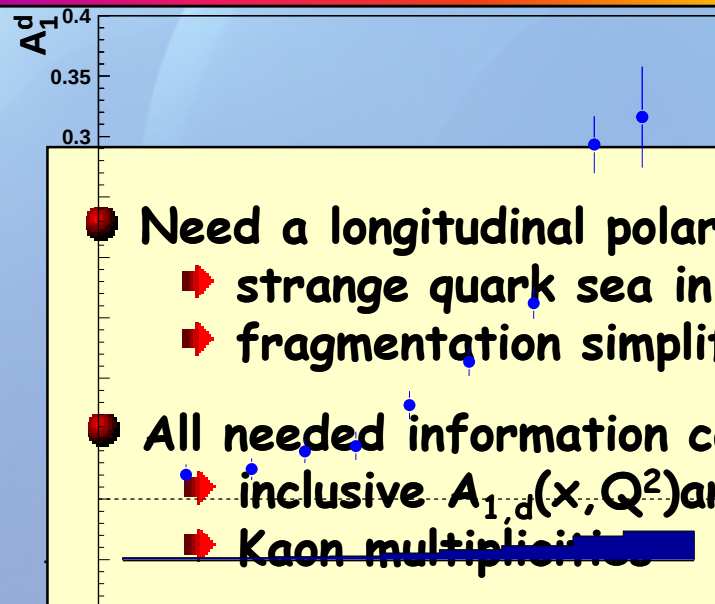
N. Bohr



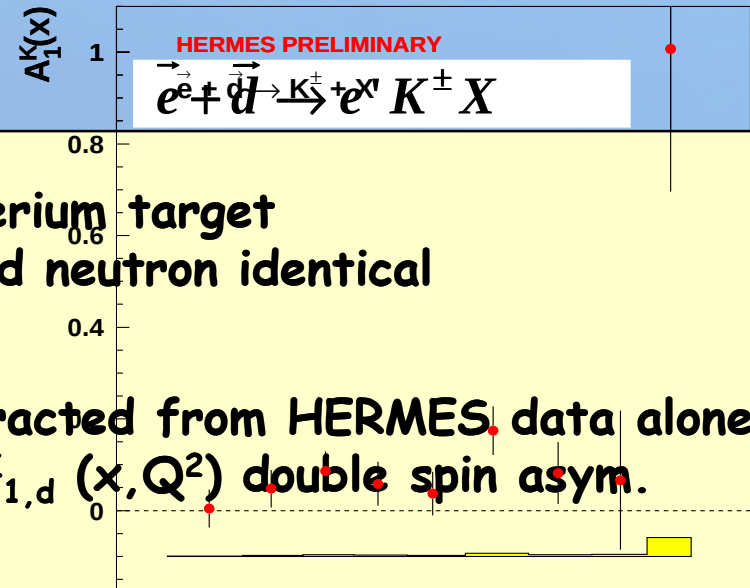
Thank you for your attention

BACKUP SLIDES

Results for ΔQ and ΔS



Inc

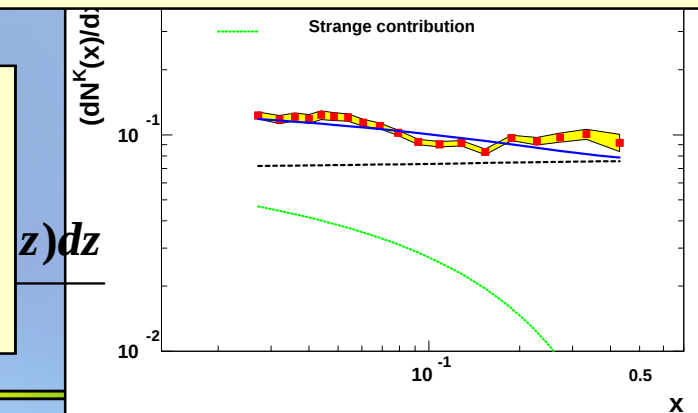


Kaon

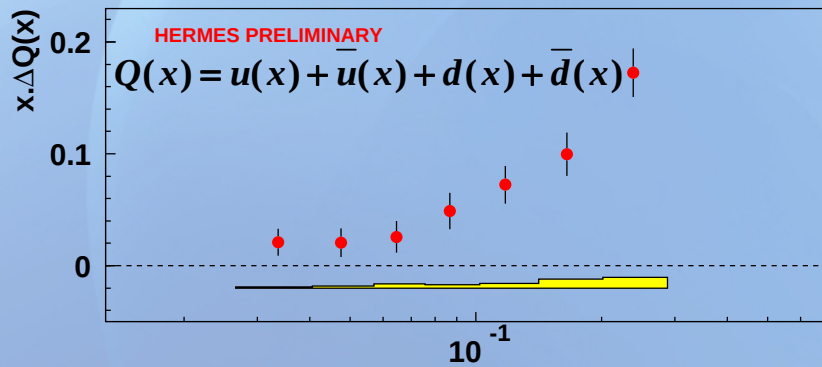
- Need a longitudinal polarized deuterium target
 - strange quark sea in proton and neutron identical
 - fragmentation simplifies
- All needed information can be extracted from HERMES data alone
 - inclusive $A_{1,d}(x, Q^2)$ and kaon $A_{1,d}^K(x, Q^2)$ double spin asym.
 - Kaon multiplicities
- Only assumptions used:
 - isospin symmetry between proton and neutron
 - charge-conjugation invariance in fragmentation

$$\int D_{\text{strange}}^K(z) dz = 2 \int D(z)_s dz$$

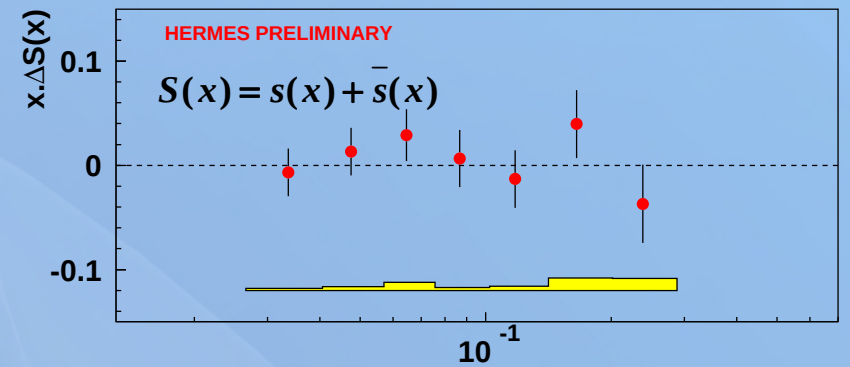
	This Work	Kretzer	KKP
$\int D_{\text{nstrg}}^K(z) dz$	0.41 ± 0.02	1.103	1.111
$\int D_{\text{strg}}^K(z) dz$	1.41 ± 0.29	0.783	0.296



Results for ΔQ and ΔS

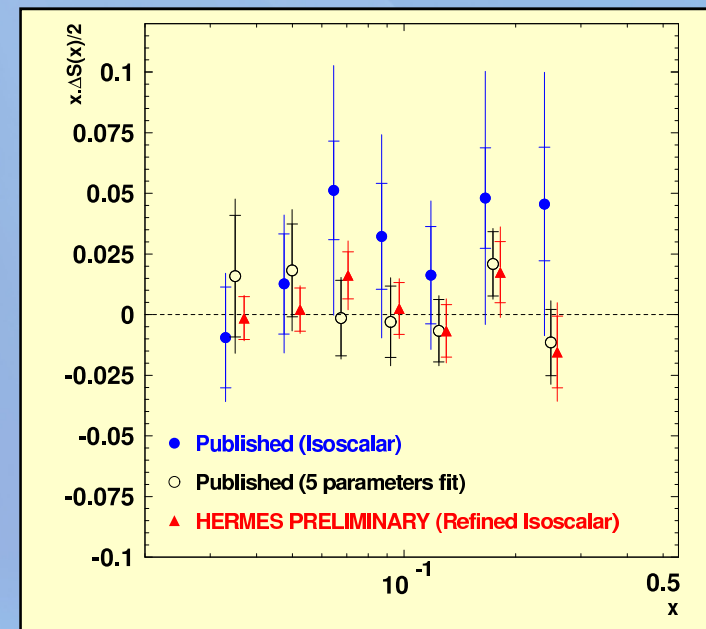


$$\int_{0.02}^1 \Delta Q = 0.286 \pm 0.026 \pm 0.011$$



$$\int_{0.02}^1 \Delta S = 0.006 \pm 0.029 \pm 0.007$$

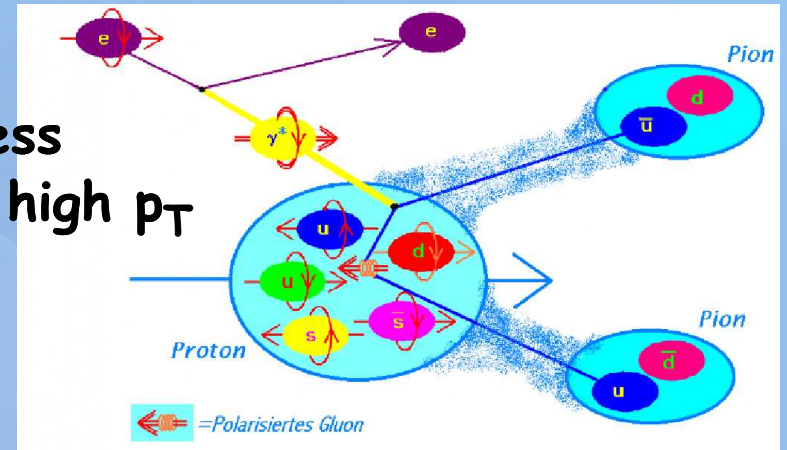
- Earlier HERMES conclusions of unpolarized strange sea confirmed
 - ➡ factor 2 smaller error bars
- Errors very sensitive to FF input



The golden channels

Idea: Direct measurement of ΔG

- ⇒ Isolate the photon gluon fusion process
- detection of hadronic final states with high p_T
 - ➡ high p_T pairs of hadrons
 - ➡ single high p_T hadrons



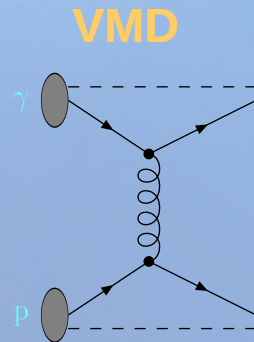
	<u>COMPASS</u>	<u>HERMES</u>	
$h^\pm h^\pm$	$\langle Q^2 \rangle > 1$	---	less sub-processes contributing ☺
	$\langle Q^2 \rangle < 1$	$\langle Q^2 \rangle < 0.1$	more sub-processes contributing ☹ higher statistics ☺
$h^\pm$	---	$\langle Q^2 \rangle > 0.1$	less sub-processes contributing ☺
	---	$\langle Q^2 \rangle < 0.1$	more sub-processes contributing ☹ higher statistics ☺

$h^\pm h^\pm$ vs. $h^\pm$:

$h^\pm$ more inclusive → pQCD NLO calculations (easier) possible

The Model in PYTHIA-6

- **Model:** Mix processes with different photon characteristics



a)
„Soft“ VMD:

Elastic, diffractive,

non-diffractive

(„minimum bias“)

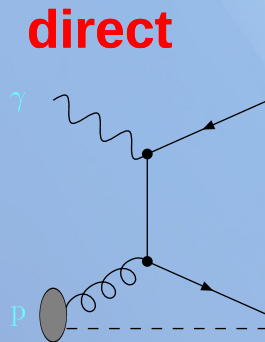
„low- p_T “

processes

Hard VMD:

„resolved“

processes



b)

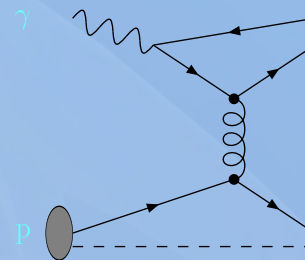
QCD-

Compton

DIS (= „resolved“)

processes

anomalous photon



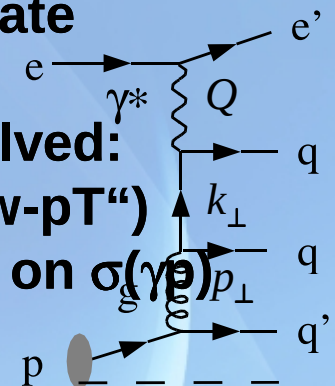
c)

Splitting of

$\gamma \rightarrow qq$

processes

- Small Q^2 : VMD + anomalous (= „resolved“)
- Large Q^2 : LO DIS dominates DIS processes
- Choice of hard process according to hardest scale involved:
- If all scales are too small, photon's VMD (diffractive or „low- p_T “)
- The „resolved“ part is modeled to match the world data on $\sigma(\gamma p)$



- Measured asymmetry is an incoherent superposition of different sub-process asymmetries:

$$A_{\parallel}^{meas}(p_t) = \sum_i f_i A_{\parallel}^i = f_{Bg} A_{\parallel}^{Bg} + f_{Sig} A_{\parallel}^{Sig}; \quad f_i = \frac{\sigma_i}{\sigma_{tot}}$$

- Signal: Gluon of the nucleon in the initial state

$$A_{\parallel}^{Sig}(p_t) = \left\langle \hat{a}_{Sig} \right\rangle \left\langle \frac{\Delta G}{G} \right\rangle; \quad \hat{a}_{Sig} : \text{Asymmetry of the hard sub-process}$$

- Background: all other sub-processes

- The gluon polarization is then:

$$\left\langle \frac{\Delta G}{G} \right\rangle = \frac{1}{f_{Sig} \langle \hat{a} \rangle} [A_{\parallel}^{meas} - f_{Bg} A_{\parallel}^{Bg}]$$

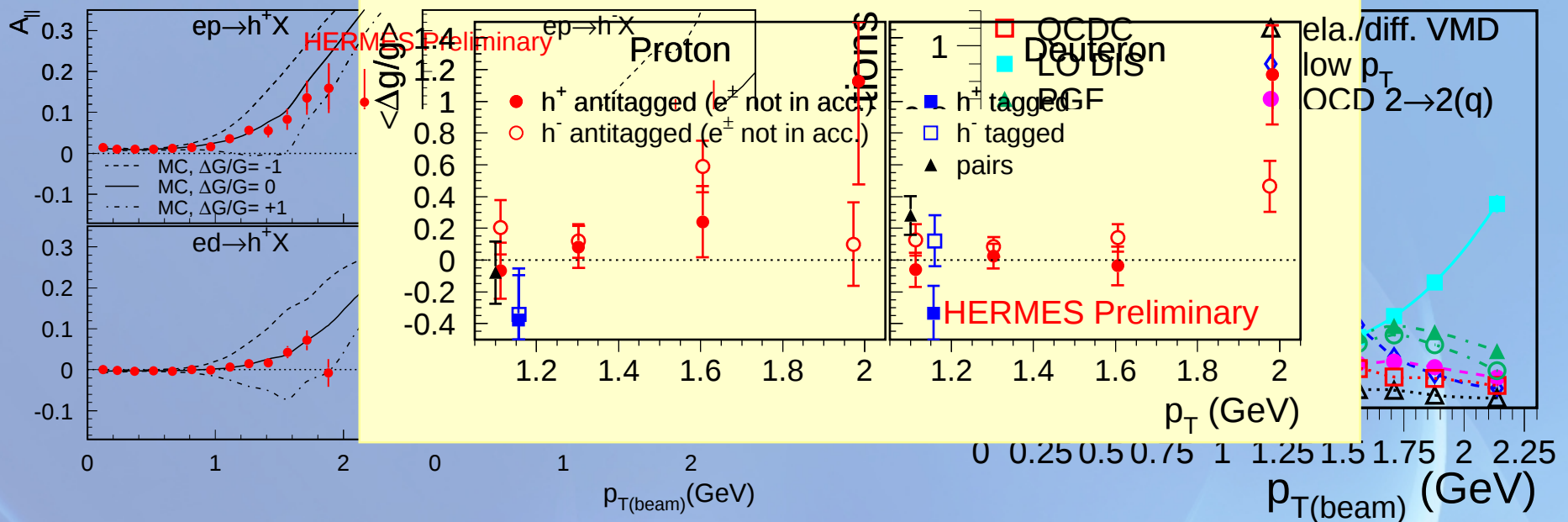
HERMES Results

Channels:

- ➡ $h^\pm$ with $Q^2 < 0.1 \text{ GeV}^2$; $p_T > 1 \text{ GeV}$
- ➡ $h^\pm$ with $Q^2 > 0.1 \text{ GeV}^2$; $p_T > 1 \text{ GeV}$
- ➡ $h^\pm h^\pm$ with $Q^2 < 0.1 \text{ GeV}^2$;
 $p_T^1, p_T^2 > 0.5 \text{ GeV}$ $p_{T,1}^2 + p_{T,2}^2 > 2.0 \text{ GeV}^2$

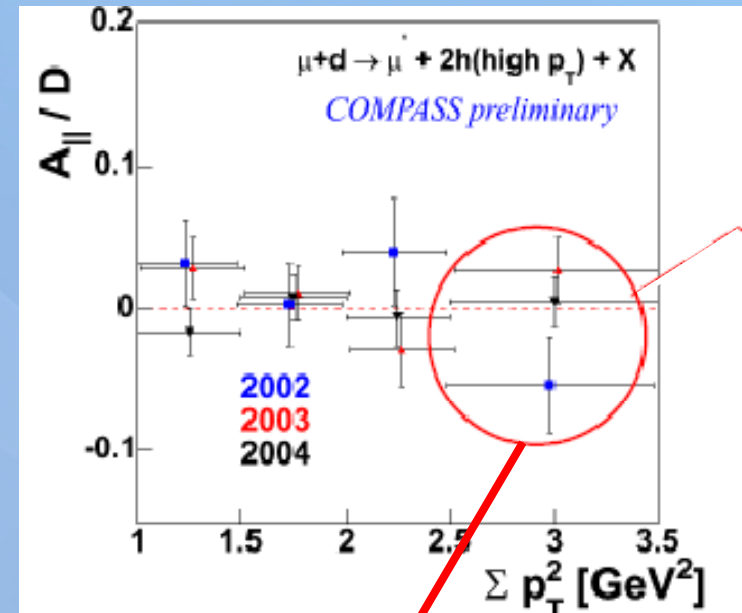
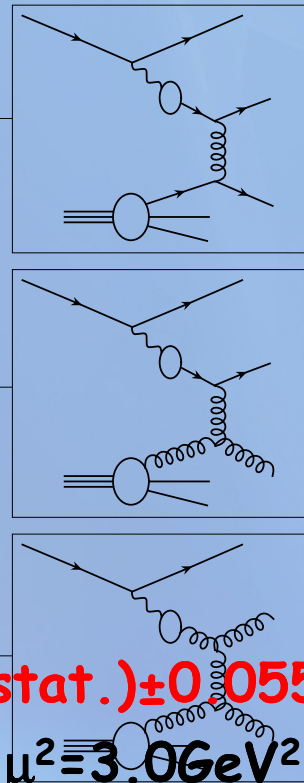
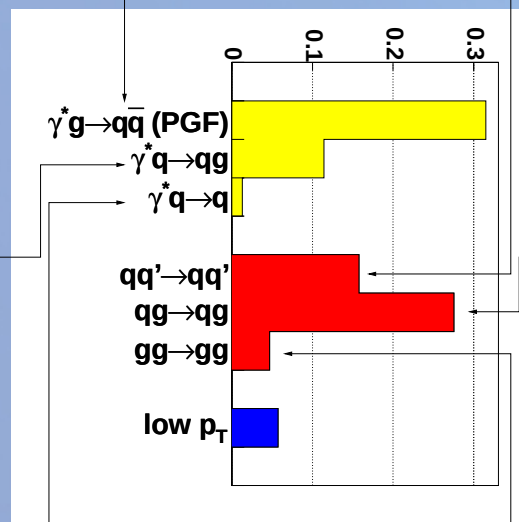
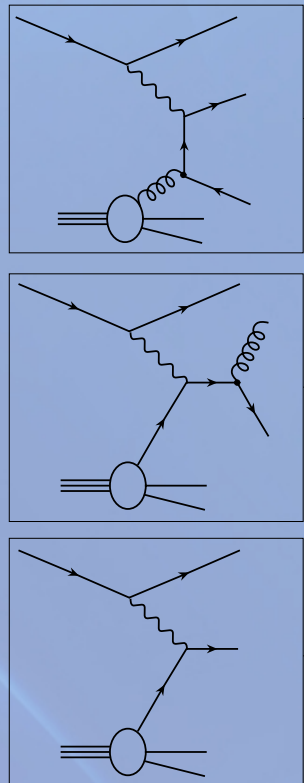
Considered sub-processes:

- ➡ LO-MC: PYTHIA 6.2



COMPASS Results

- Channel: $h^\pm h^\pm$ with $Q^2 < 1 \text{ GeV}^2$
- Cuts:
 $p_T^1, p_T^2 > 0.7 \text{ GeV}$ $p_{T,1}^2 + p_{T,2}^2 > 2.5 \text{ GeV}^2$
- Considered sub-processes:
 ➔ LO-MC: PYTHIA 6.2



used for $\Delta g/g$ extraction

$$\Delta g/g = 0.016 \pm 0.058(\text{stat.}) \pm 0.055(\text{syst.})$$

$$\langle x_g \rangle = 0.085^{+0.07}_{-0.035} \quad \mu^2 = 3.0 \text{ GeV}^2$$

Cross Section Data-MC

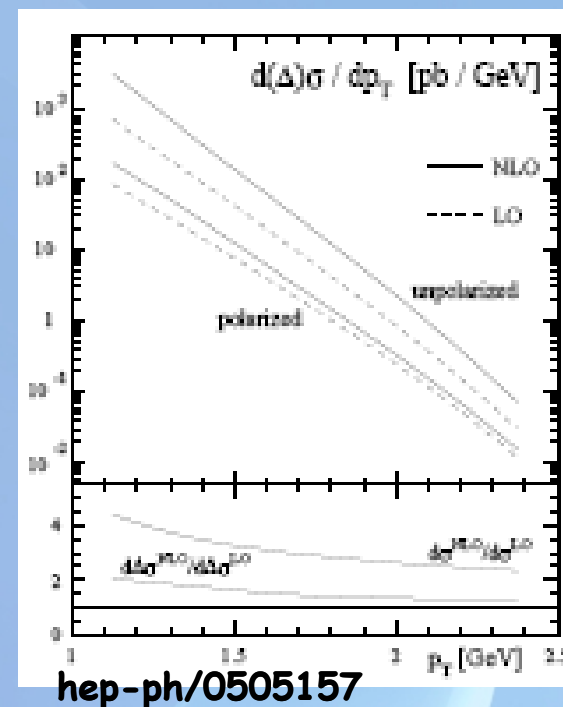
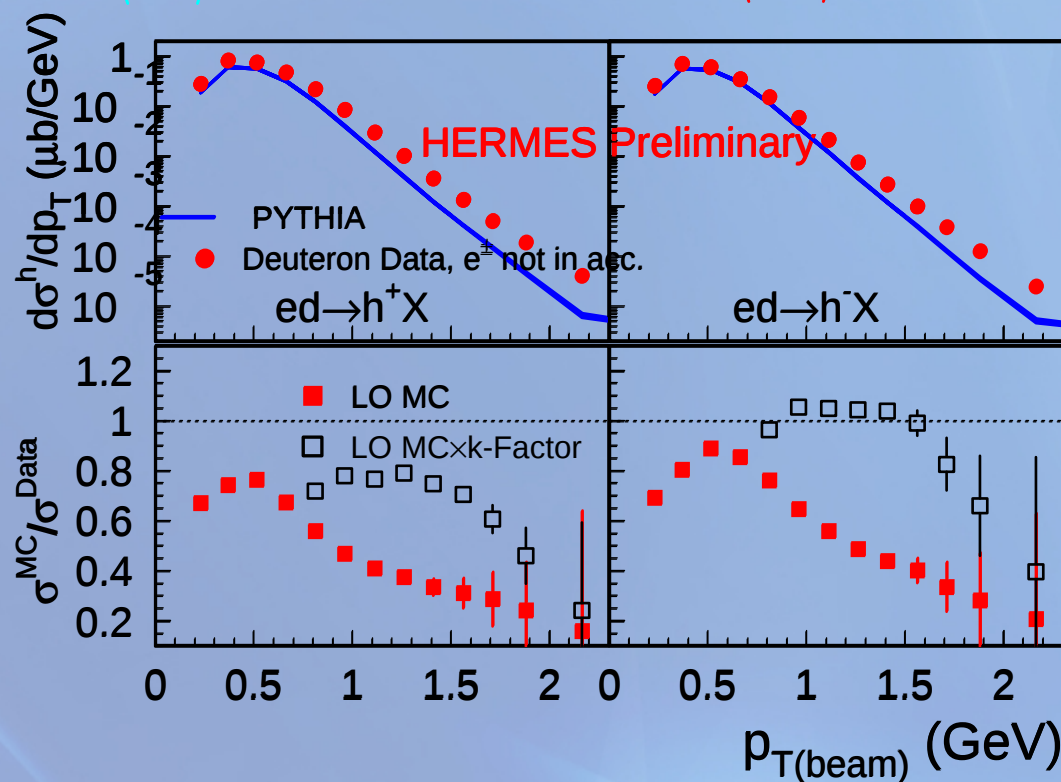
HERMES cross section:

k_T , p_T standard values:

$$\langle k_T^{p,\gamma} \rangle = 0.40 \text{ GeV} \quad \langle p_T^{\text{frag}} \rangle = 0.40 \text{ GeV}$$

➡ M. Anselmino et al. Phys. Rev. D71,074006

$$\langle k_t^2 \rangle = 0.25 \text{ GeV}^2 \pm 20\% \quad \langle p_t^2 \rangle = 0.20 \text{ GeV}^2 \pm 20\%$$

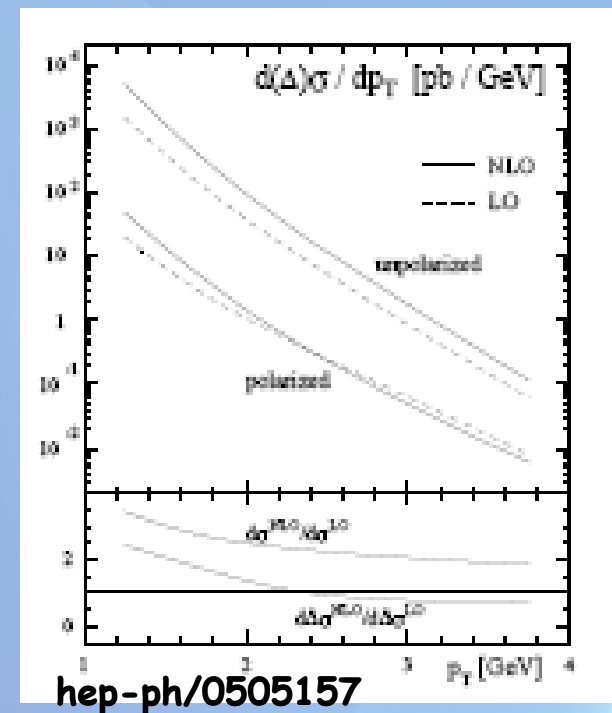
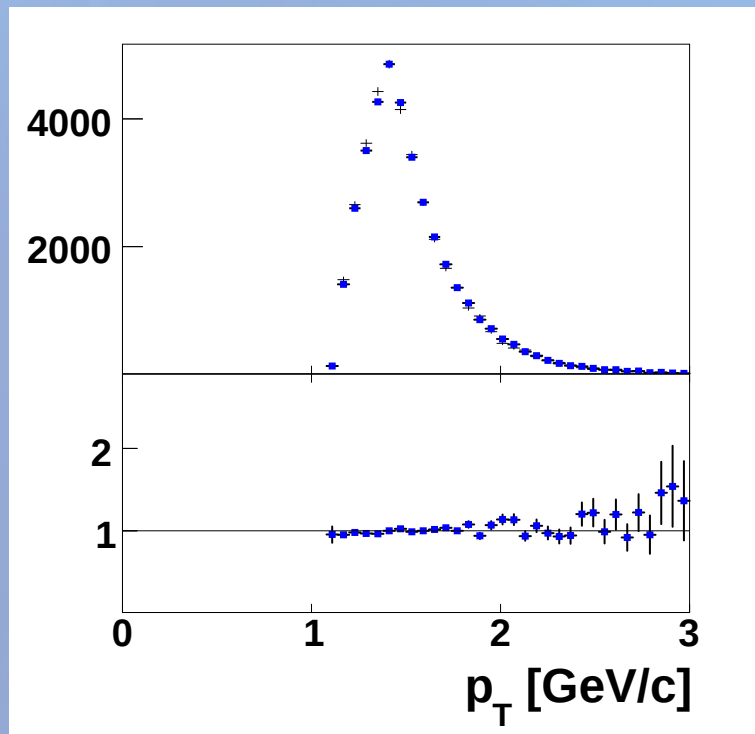


Cross Section Data - MC

■ COMPASS standard values:

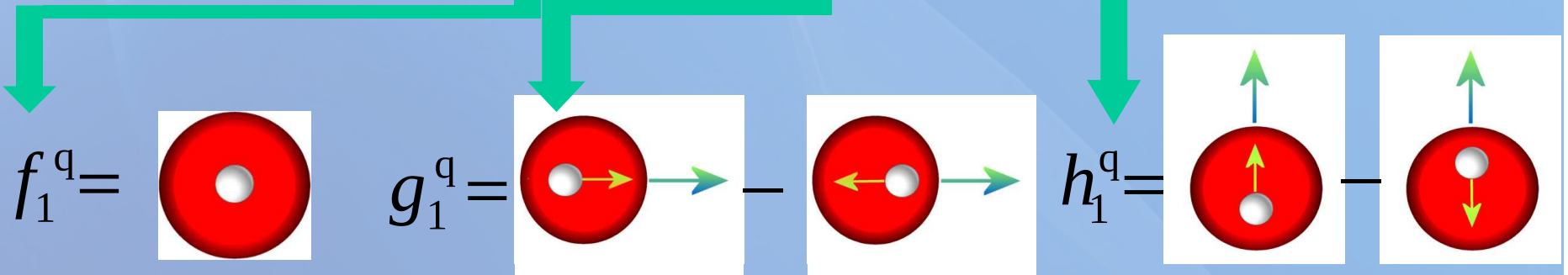
$$\langle k_T^\gamma \rangle = 0.50 \text{ GeV} \quad \langle k_T^p \rangle = 1.00 \text{ GeV}$$

$$\langle p_T^{\text{frag}} \rangle = 0.36 \text{ GeV}$$



The 3rd Twist-2 structure function

$$\Phi_{\text{Corr}}^{\text{Tw2}}(x) = \frac{1}{2} \left\{ f_1(x) + S_L g_1(x) \gamma_5 + h_1(x) \gamma_5 \gamma^1 S_T \right\} n^+$$



unpolarised quarks
and nucleons

longitudinally polarized
quarks and nucleons

transversely polarized
quarks and nucleons

$q(x)$: spin averaged

$\Delta q(x)$: helicity difference

$\delta q(x)$: *helicity flip*

→ vector charge

→ axial charge

→ tensor charge

well known

known

measuring!

Peculiarities of transversity

$$\Phi_{\text{Corr}}^{\text{Tw2}}(x) = \frac{1}{2} \left\{ f_1(x) + S_L g_1(x) \gamma_5 + h_1(x) \gamma_5 \gamma^1 S_T \right\} n^+$$

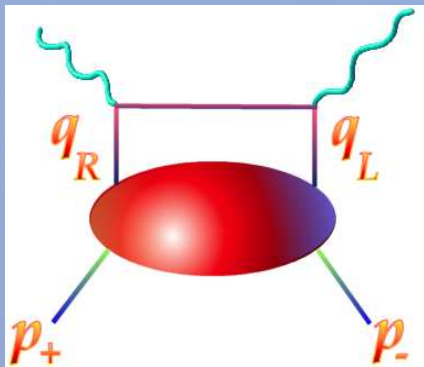
relativistic nature of quark:

in absence of relativistic effects $h_1(x) = g_1(x)$

Q^2 -evolution: unlike for $g_1^p(x)$,
the gluon doesn't mix with quark in $h_1^p(x)$
(no gluon analog for spin- $\frac{1}{2}$ nucleon)

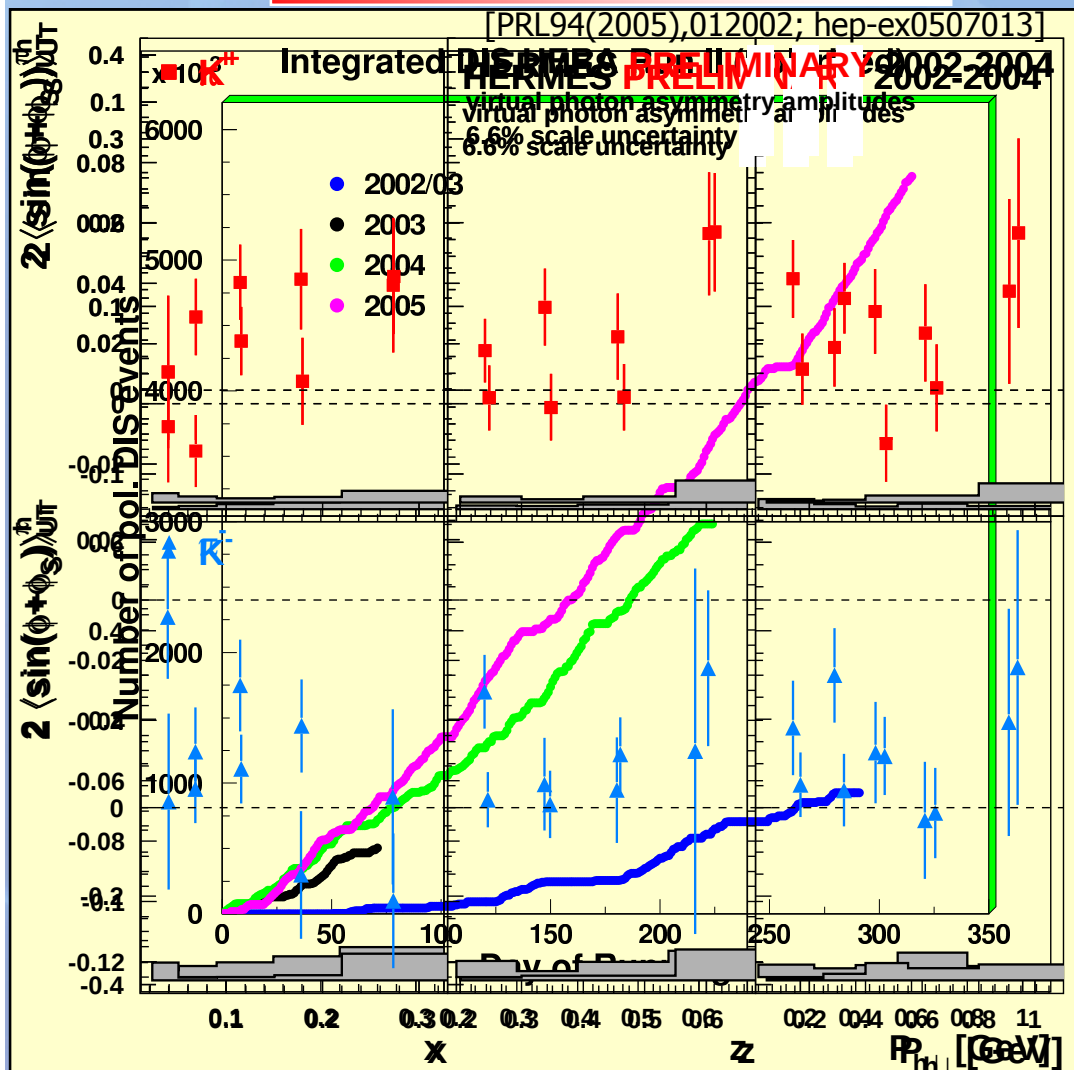
sensitive to the valence quark polarisation
 q and $\bar{q}$ have opposite sign.

tensor charge: first moment of h_1
(large from lattice QCD)



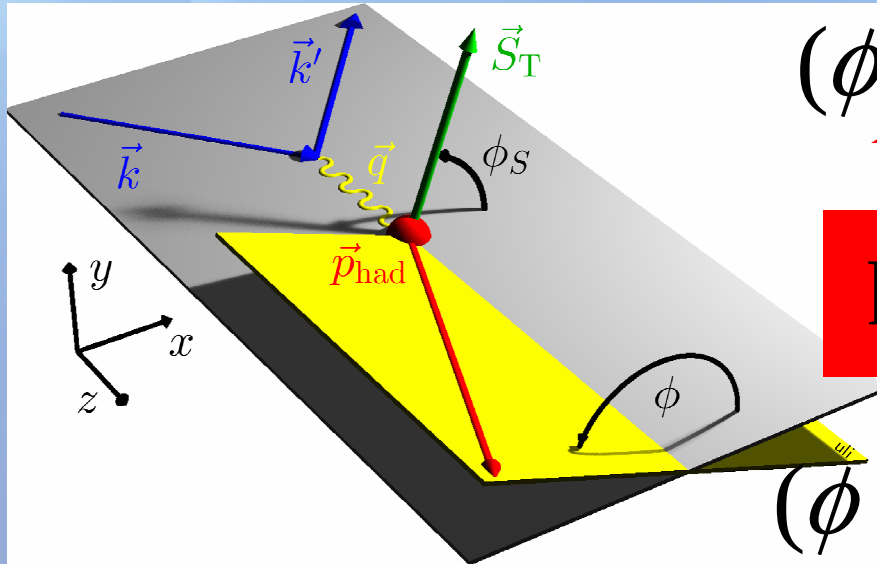
*single helicity
flip*

$$A_{\text{coll}} \propto h_1(x) H_1^\perp(z)$$



- Collins moment:
 $\pi^+ > 0$ $\pi^- < 0$
- π^- unexpected large
 $\Rightarrow$ role of unfavoured FF
- first data for Collins FF $H_1^\perp(z)$ available from Belle
 $\Rightarrow$ extraction of h_1 from Hermes asymmetries
- $K^+ > 0$ $K^- > 0$
 K^+ in agreement with π^+
- 2005 data:
 $\Rightarrow$ 2 dimensional binning disentangle x (PDF) and z (FF) dependence

Azimuthal angles and asymmetries



$(\phi + \phi_S)$ angle of hadron relative to *final* quark spin (Collins)

$h_1 \otimes H_1^\perp$ (Collins)

$(\phi - \phi_S)$ angle of hadron relative to *initial* quark spin (Sivers)

(Sivers)

$f_{1T}^\perp \otimes D_1$

peculiarity of $f_{1T}^\perp$
chiral-even naïve T-odd DF

related to parton orbital
momentum

violates naïve universality of PDF
⇒ different sign of $f_{1T}^\perp$ in DY